

Crystallography: A foray into reciprocal space

Chem 651B: Electron Microscopy

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May 19, 2015

Central Instrument Facility Summer school

State of the art instruments

Hands-on training for students

crystallography: learning goals

- Solids are often crystalline - *the language of waves*
 - Crystalline symmetry - *simplifying structures*
 - Describing structures - *using the Int. Tab. Crystallography*
 - Reciprocal space - *the Fourier transform of direct space*
 - Reciprocal lattice - *what diffraction shows us*
-
- Kinematic scattering - *deriving intensities*

Why do crystallography in the TEM?

take advantage of the 'local' probe!

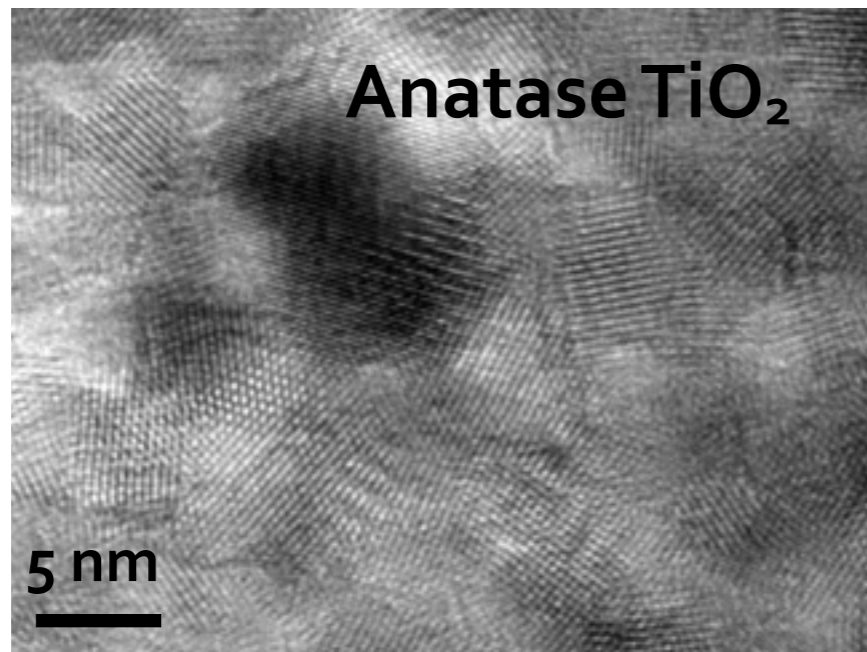
- Phase identification (lattice or point group)
- Orientation of crystals
- Finite sizes precludes macroscopic crystallography
 - *the nanostructure problem; challenges in crystal growth.*
- Supersymmetry: strong electron contrast
 - *modulated structures, charge density wave materials (e.g., TaS₂, LiFeBO₃)*
- Disorder: domains and local distortions
 - *e.g., metal-insulator transitions, ferroelectrics, superconductors*

what are crystals?

definition: crystallinity

The presence of three-dimensional order on the level of atomic dimensions. Crystallinity may be detected by diffraction techniques, heat-of-fusion measurements, etc. The amount of disorder within the crystalline region is not incompatible with this concept

- *IUPAC Compendium of Chemical Terminology - the Gold Book* (<http://goldbook.iupac.org>)



what are crystals?

Many condensed phases are crystals.

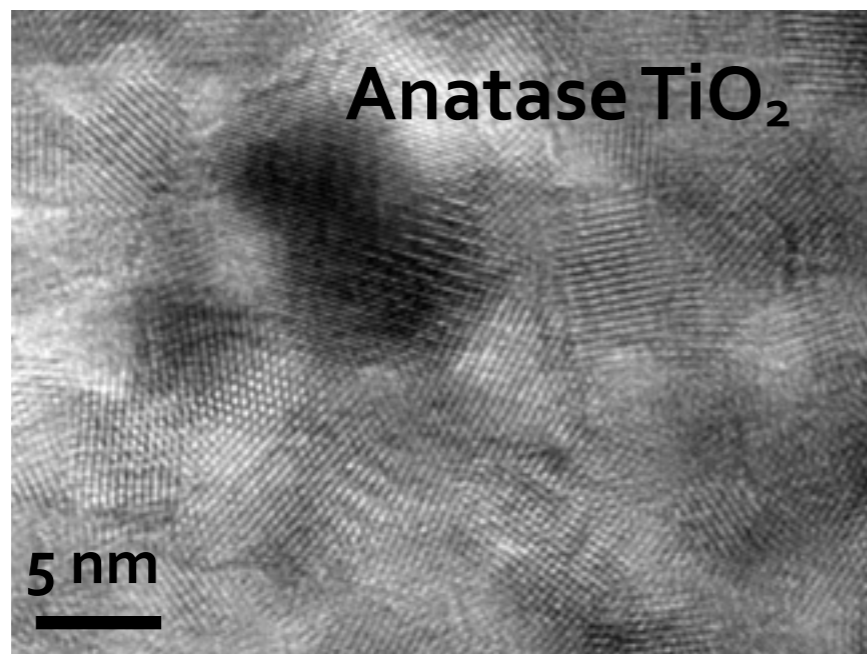
We treat them as being *infinitely periodic*.

(what's the difference between $\sim 10^{23}$ and ∞ ?)

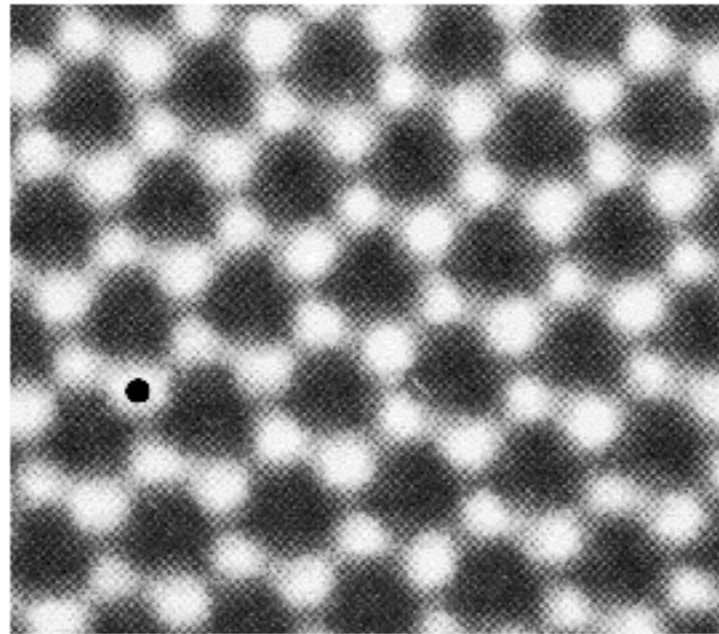
Why? Consider the packing of spheres (*i.e.*, oranges at the fruit stand).

The most dense packing for spheres (T. Hales; Kepler conjecture).

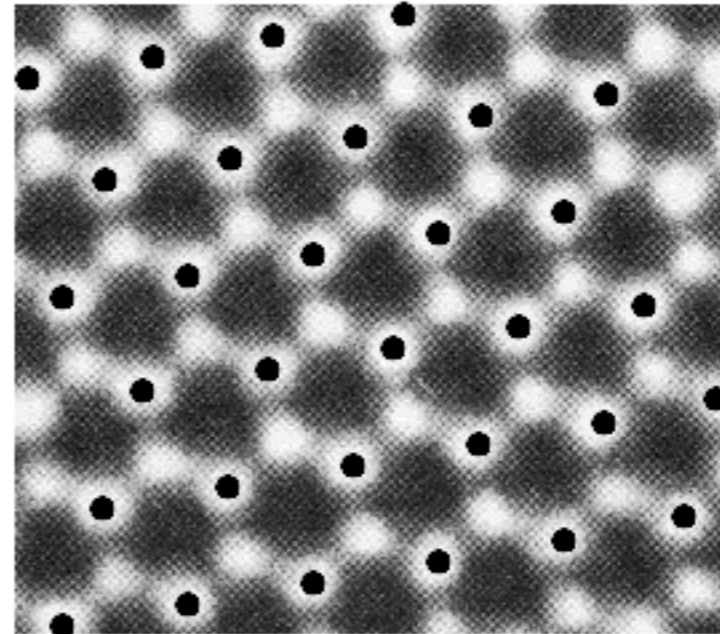
This hexagonal/cubic close packing: $\pi/(3\sqrt{2}) \sim 74\%$ (derive this for fun)



translational symmetry: the periodic lattice

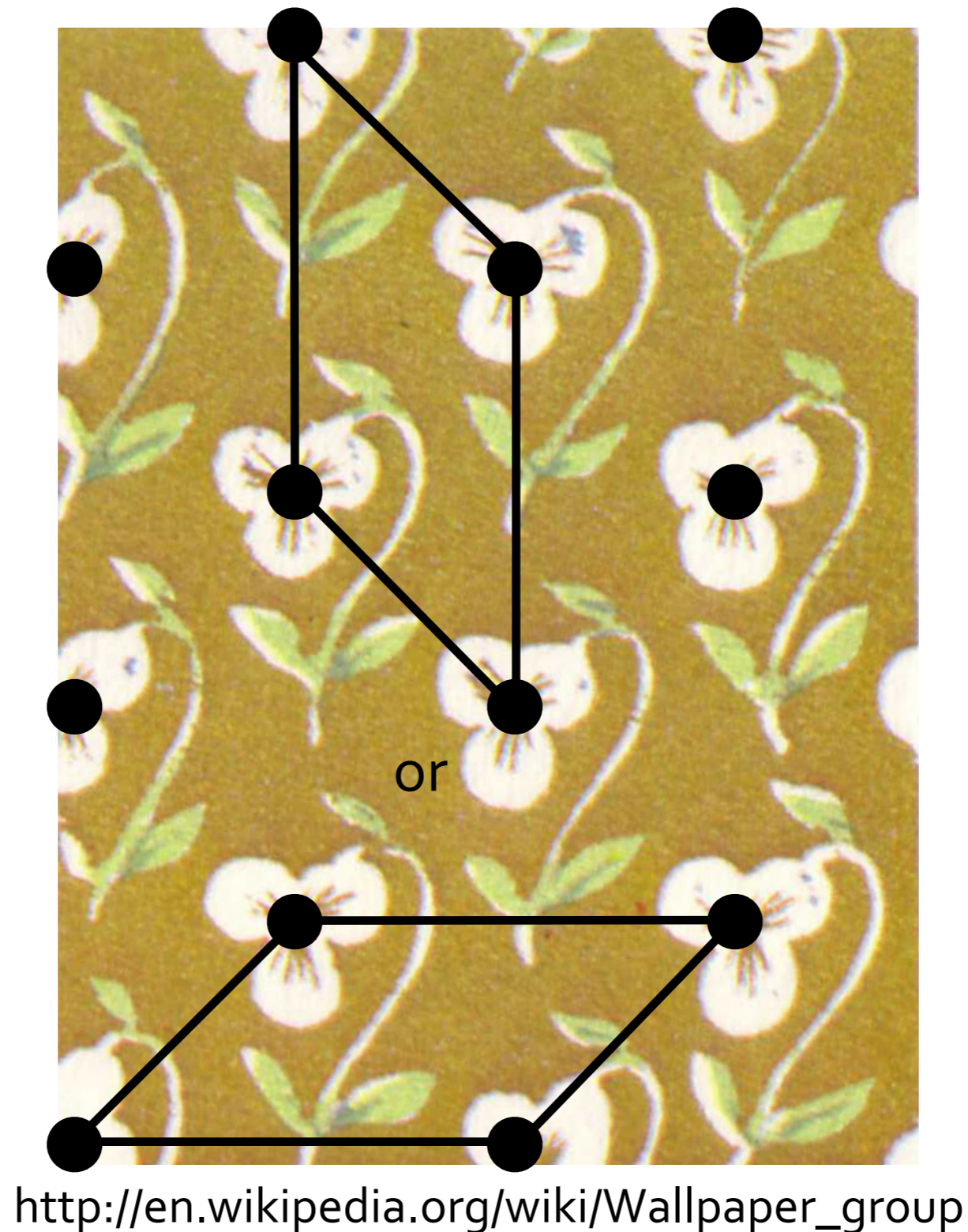


a lattice point



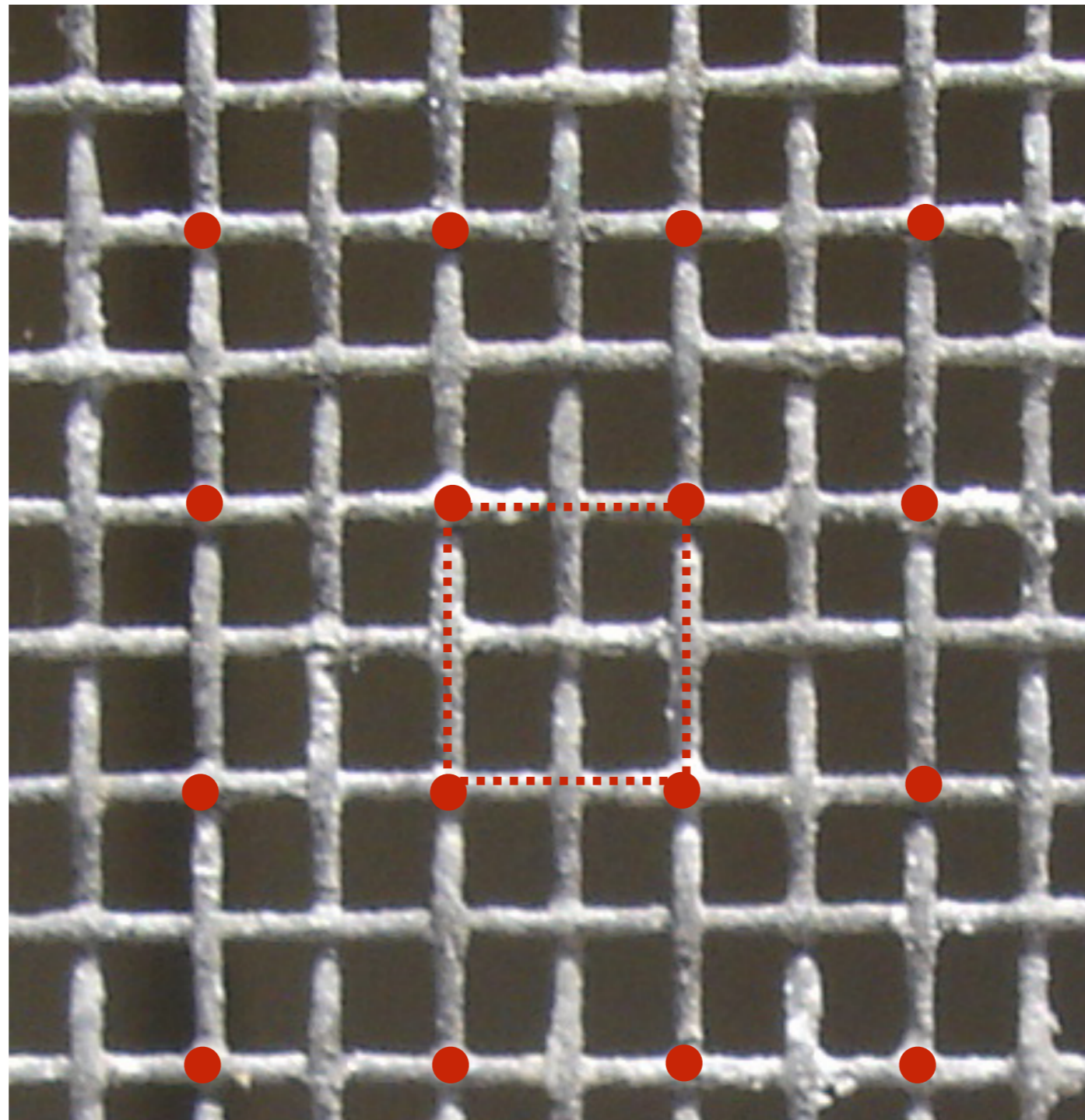
the periodic lattice

identify the lattice points



p1
Medieval Wall paper

identify the lattice points

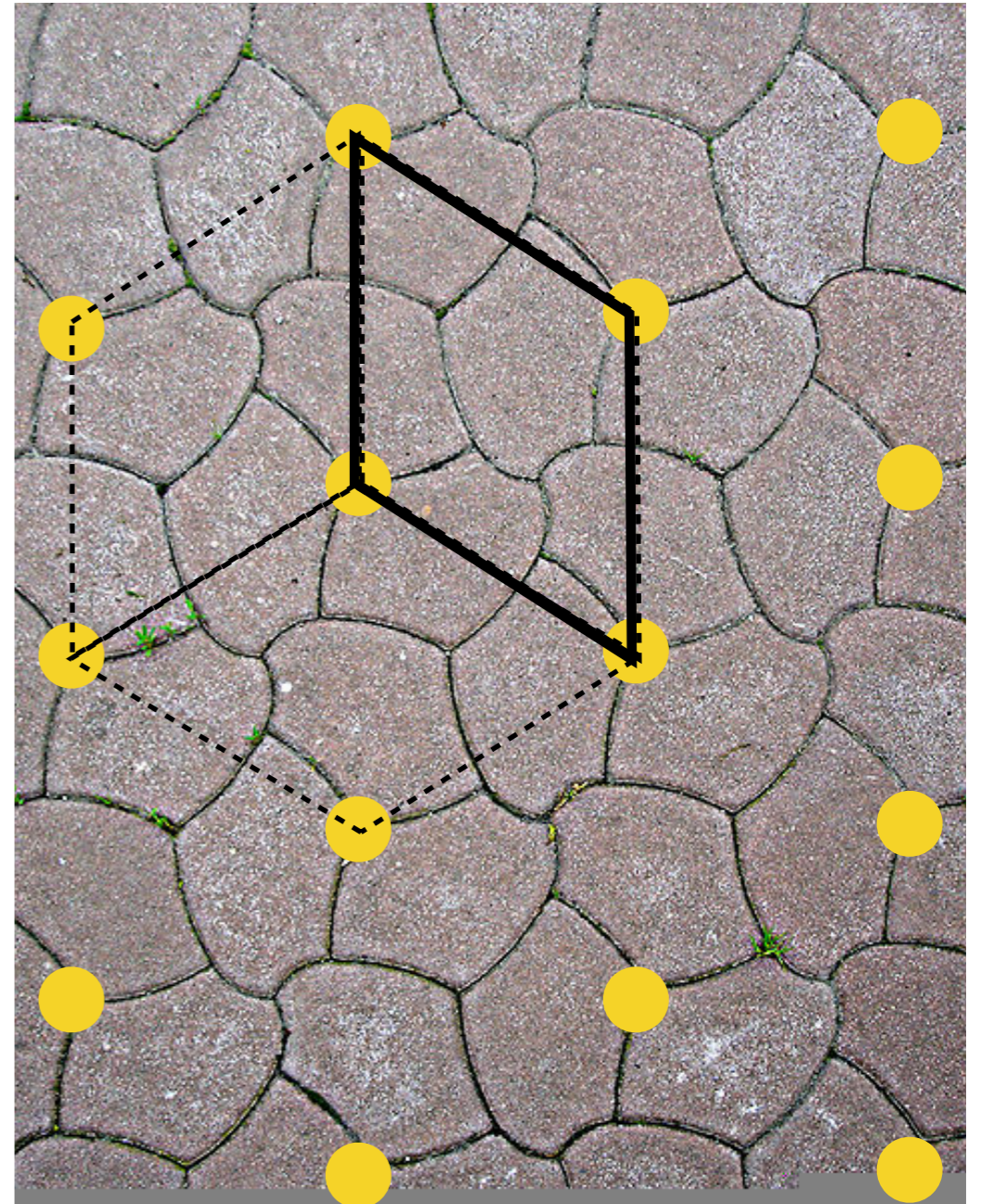


http://en.wikipedia.org/wiki/Wallpaper_group

p4g
Fly screen

the unit cell

- this is the object that tiles space
 - must tile completely, without gaps
 - the contents of the cell just also perfectly repeat
 - take the smallest repeating unit with the least # of parameters
- in 2D: parallelogram
- in 3D: parallelepiped



the geometry of the unit cell

- lattice parameters: $\mathbf{a}, \mathbf{b}, \mathbf{c}$ (use the right-hand rule)

- angles:

- α : angle between \mathbf{b} and \mathbf{c}

- β : angle between \mathbf{a} and \mathbf{c}

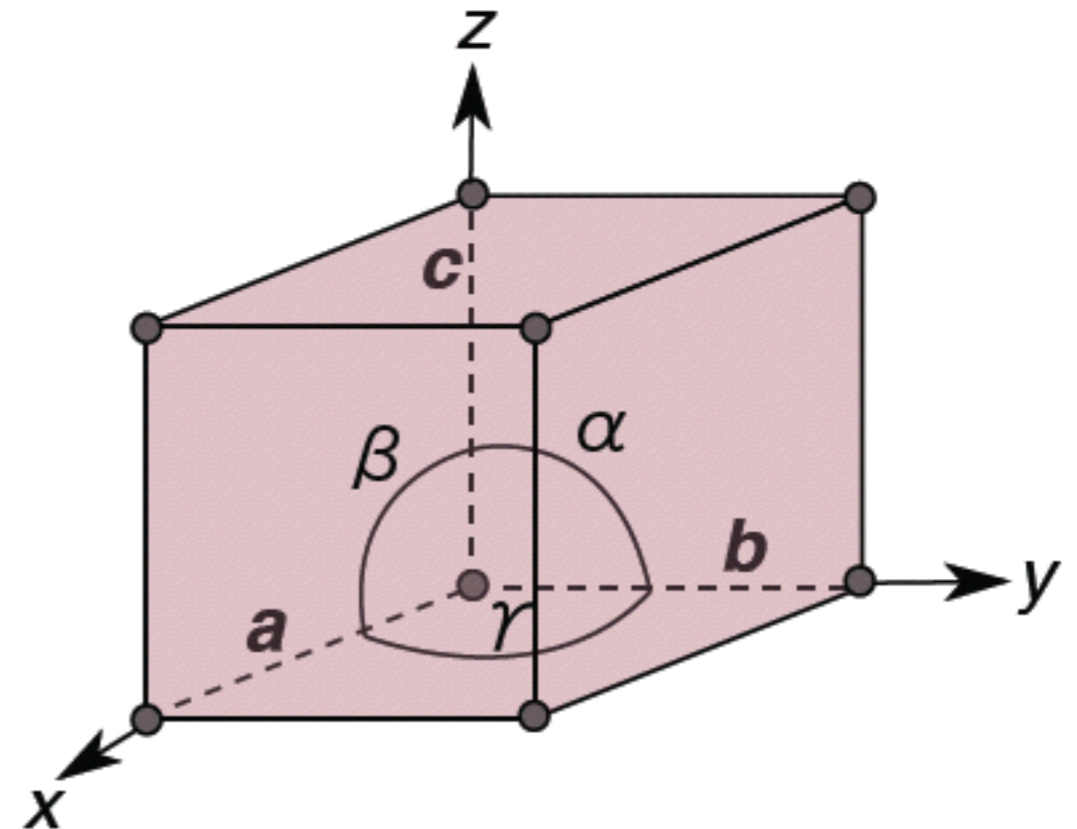
- γ : angle between \mathbf{a} and \mathbf{b}

- Lattice vectors:

- $\mathbf{t} = u\mathbf{a} + v\mathbf{b} + w\mathbf{c}$

- For shorthand, $\mathbf{t} = [u\ v\ w]$

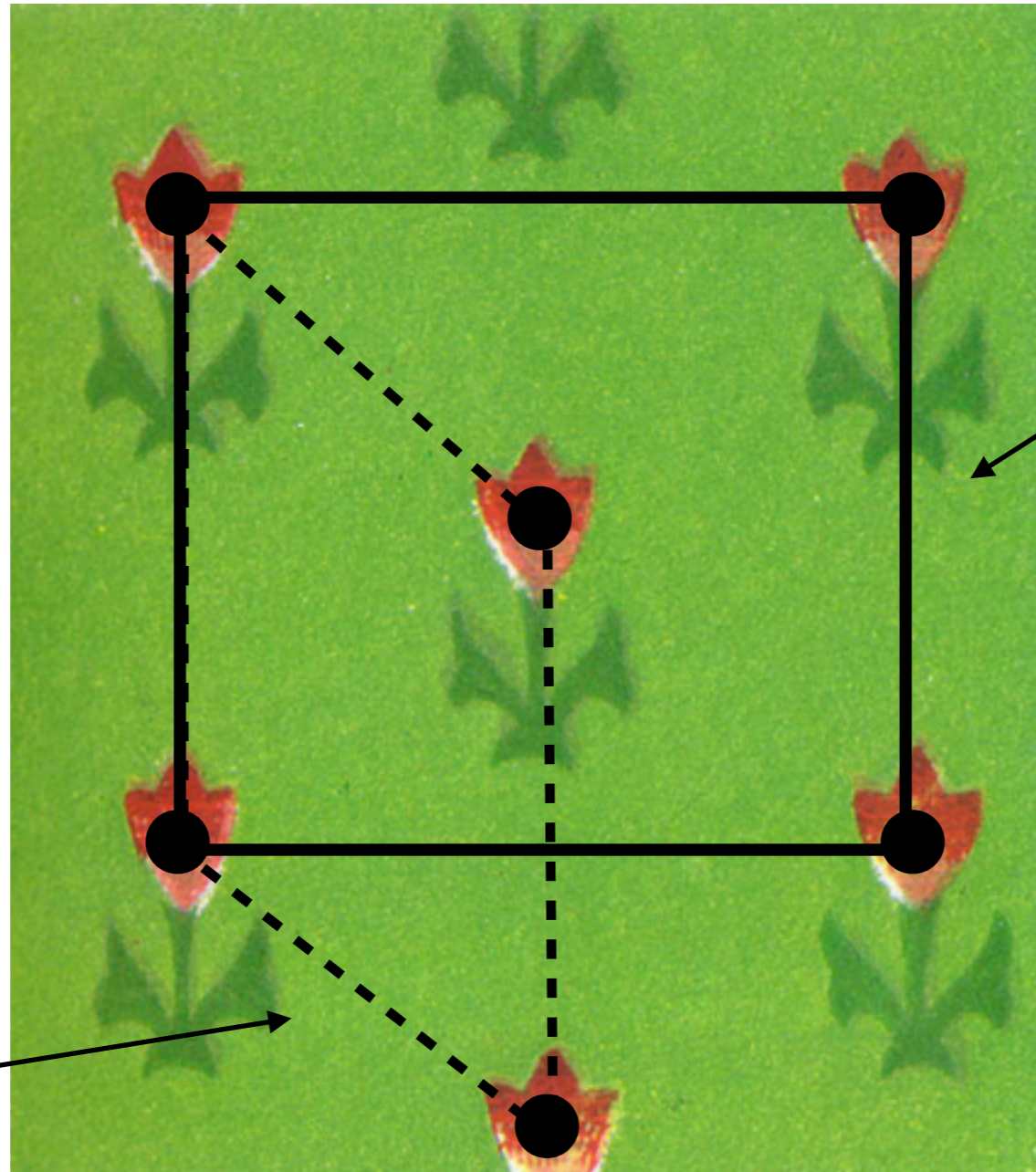
- for negative values, we write: $\mathbf{t} = [\bar{u}vw]$ (pronounced, "u-bar")



primitive vs. centered unit cells

pick the smallest cell that requires you to describe the least number of features

primitive cell
(2 edges + 1 angle)

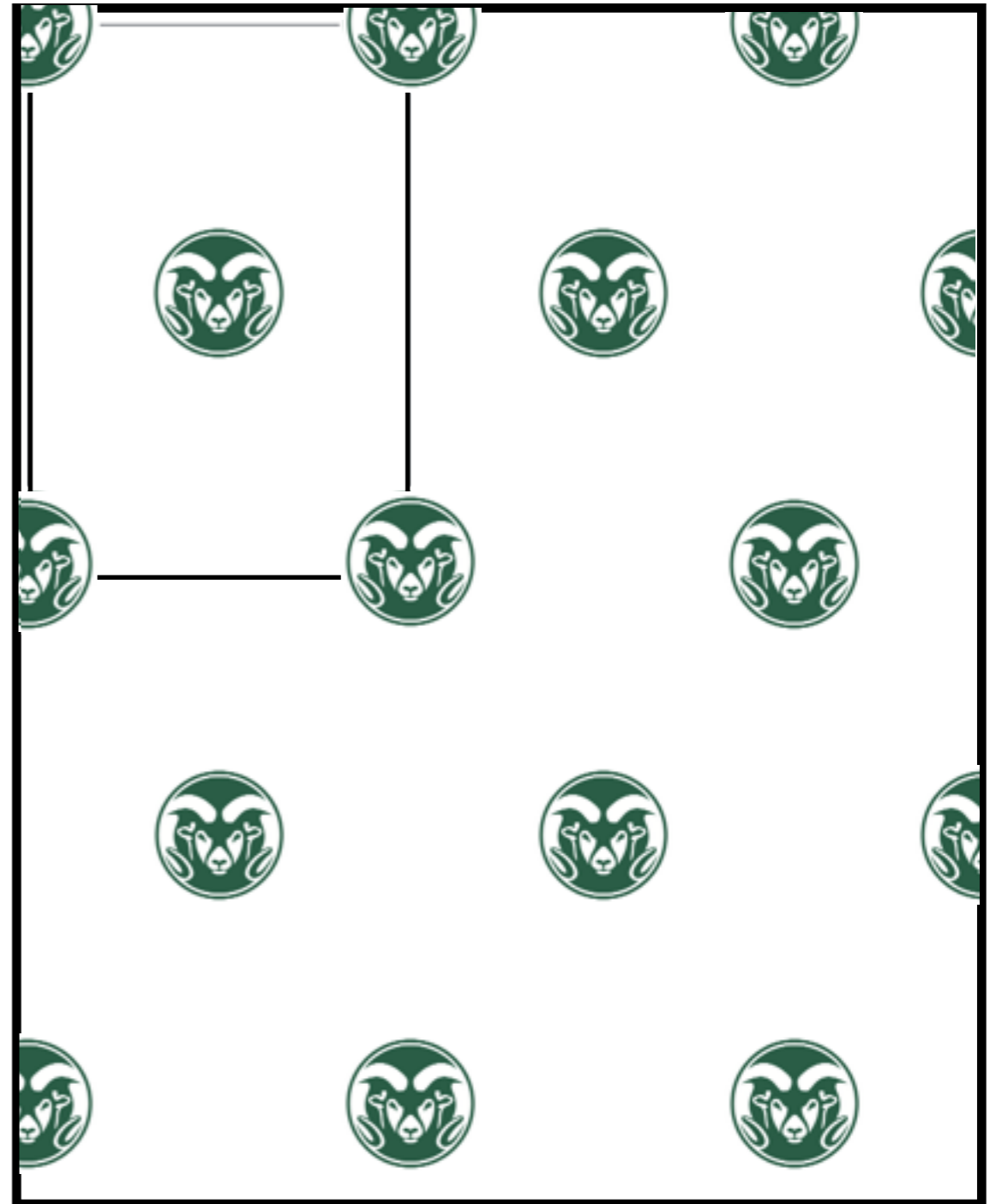


centered cell
(2 edges)

cm

building a crystal structure

- crystal system
- identify the lattice type (centering)
 - insert motifs into the unit cell
 - coordinates expressed as fractions of the lattice parameters
- tile with abandon

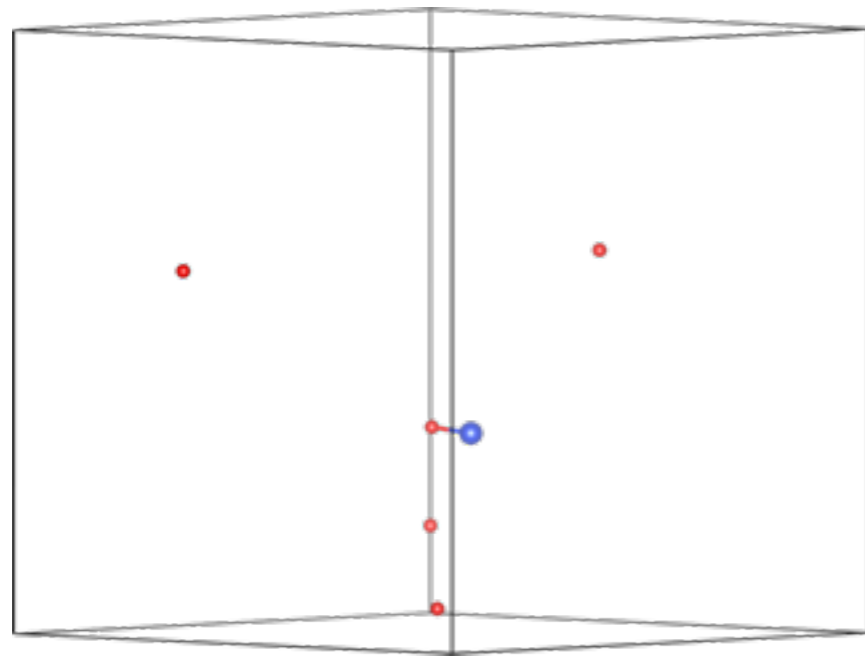


symmetry in the solid state

Why? Simplification of crystal structures

Consider the zeolite, "Faujasite" (SiO_2):

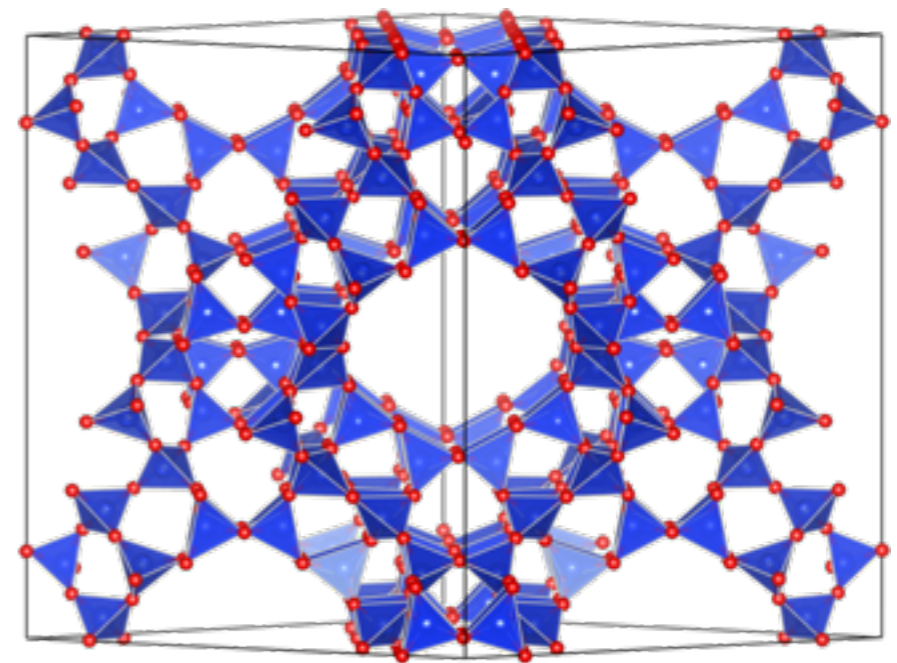
640 atoms tile space infinitely, yet only 5 symmetry-independent atoms.



*symmetry
operations*



$Fd\bar{3}m$



symmetry independent atoms
(5 atoms)

full unit cell
(640 atoms)

symmetry in the solid state

Why? Simplification of crystal structures **properties**

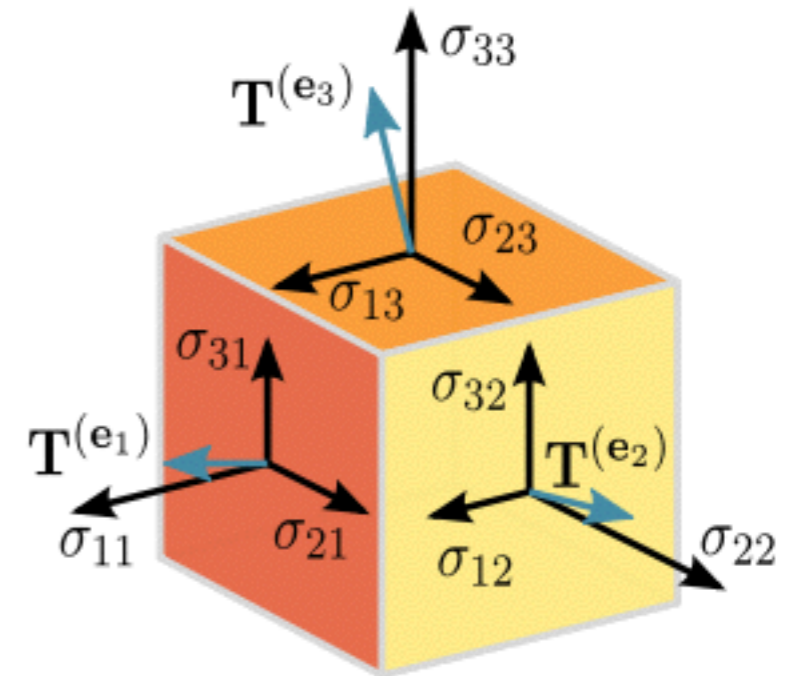
Example: the Cauchy stress tensor (2nd order)

No symmetry:

$$\sigma = \begin{bmatrix} \sigma_{11} & \sigma_{12} & \sigma_{13} \\ \sigma_{21} & \sigma_{22} & \sigma_{23} \\ \sigma_{31} & \sigma_{32} & \sigma_{33} \end{bmatrix}$$

Tetragonal symmetry:

$$\sigma = \begin{bmatrix} \sigma_{11} & 0 & 0 \\ 0 & \sigma_{22} = \sigma_{11} & 0 \\ 0 & 0 & \sigma_{33} \end{bmatrix}$$



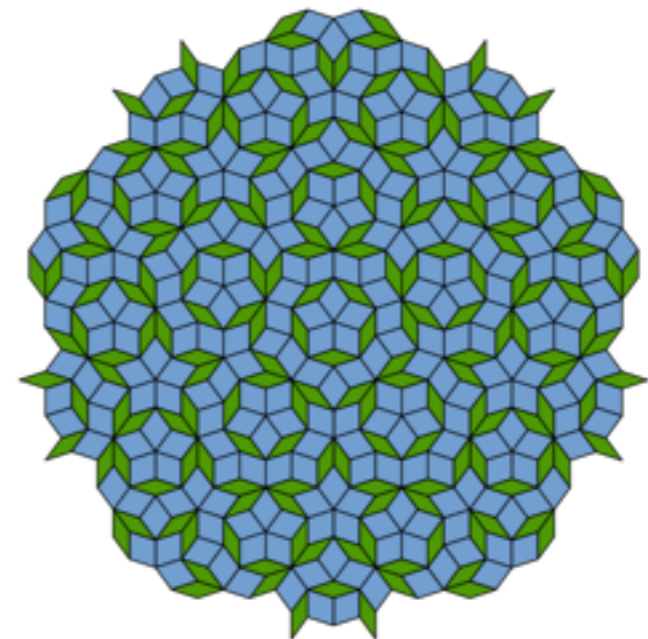
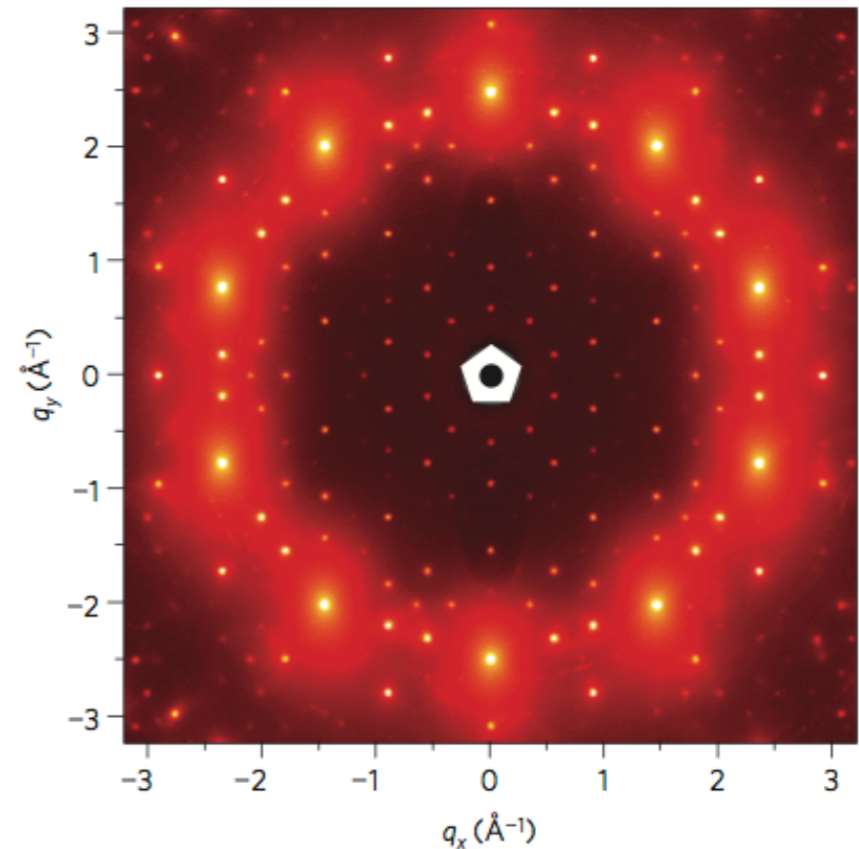
symmetry operations

- **32 point groups**
 - operations:
 - identity, rotations, mirrors, inversions, improper rotations
 - finite # of combinations (32)
- **in 3D, we also have translation operations**
 - introduces centering, screw axes and glide planes
 - centering: face centering (*F*), body centering (*I*), *c*-centering (*C*)
- **230 space groups (three-dimensional crystals belong to a closed group)**

I = Innenzentriert

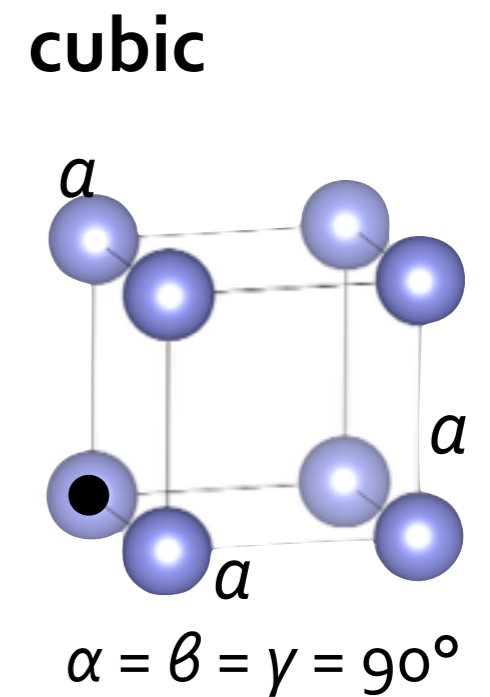
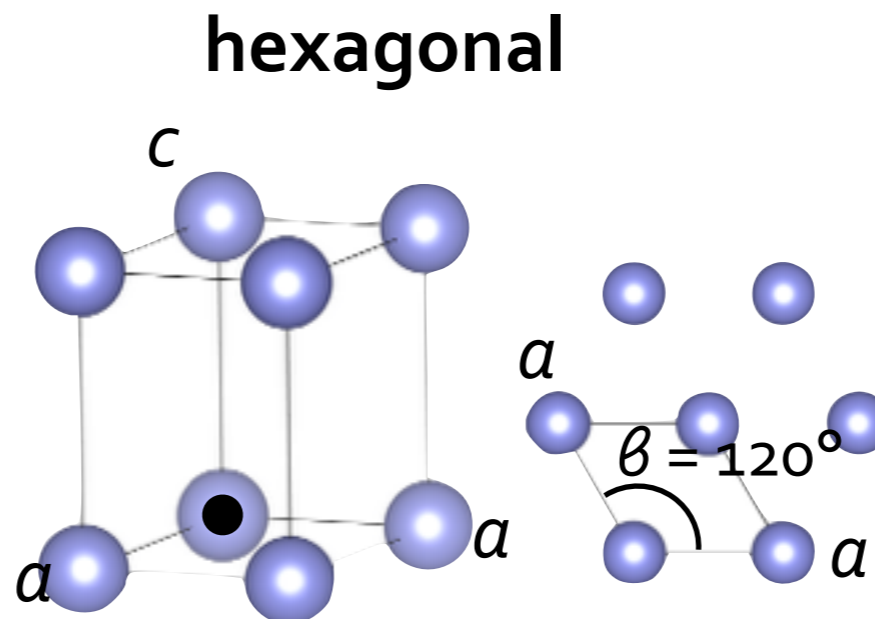
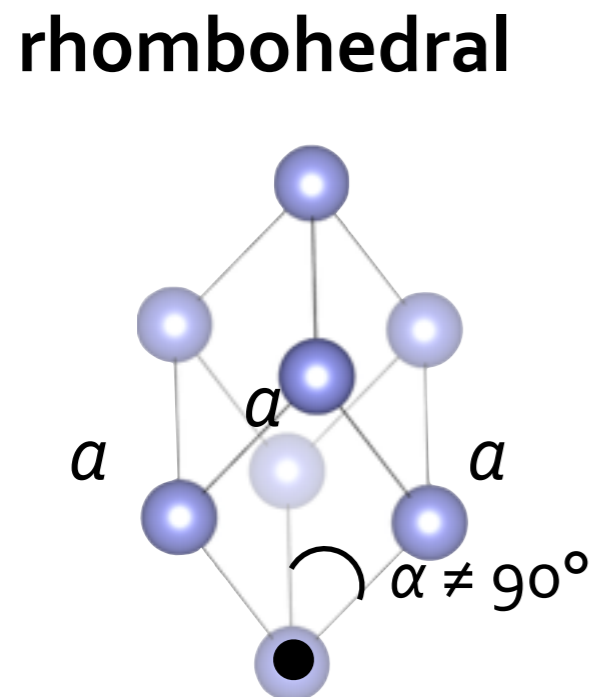
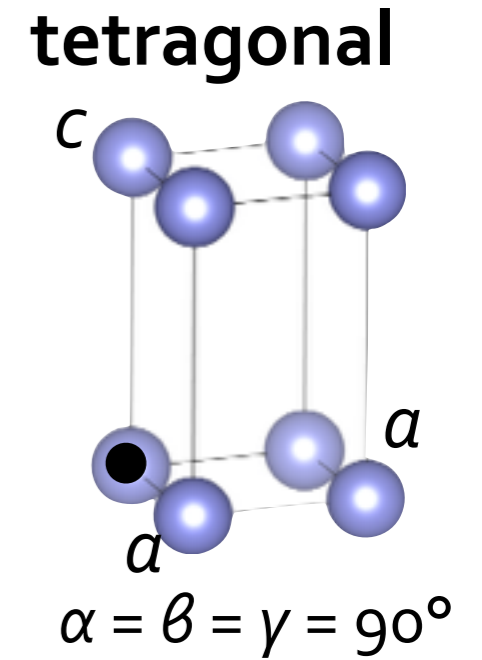
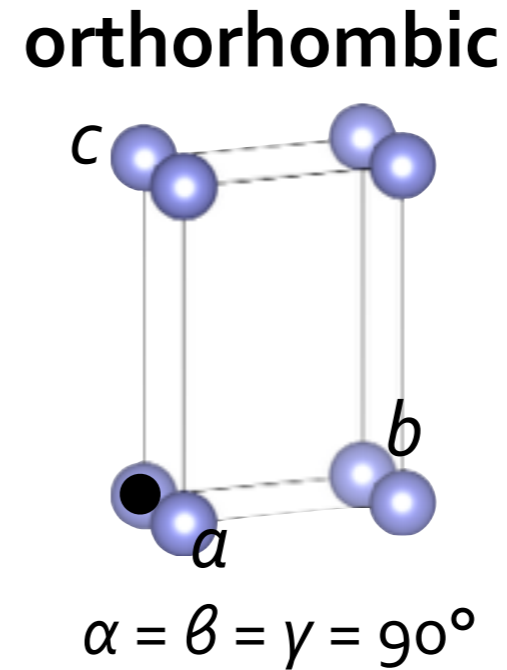
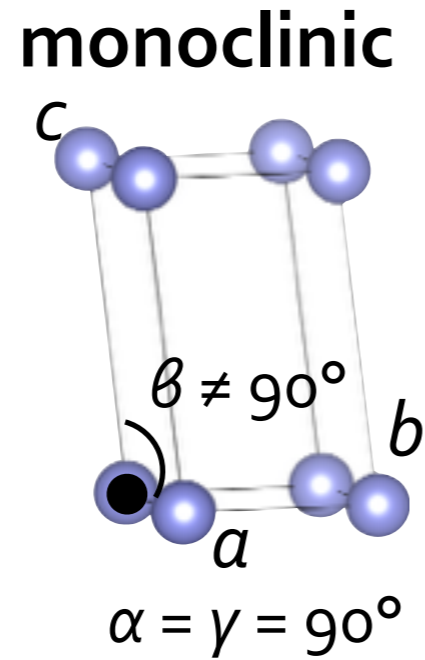
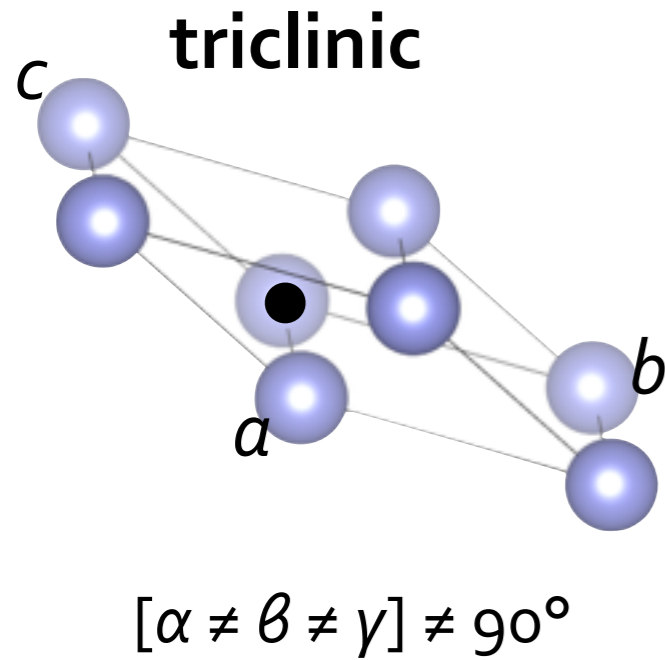
incompatible operations

- Only 2,3,4, and 6-fold rotation axes allowed in 3D.
- Yet, 5-fold rotation axes are observed in diffraction patterns!
 - **quasicrystals**
 - short range order; no long range *translational* order (in 3D). Akin to a Penrose tiling
 - possess perfect periodicity in higher dimensions (4- to 6D!!)
 - incommensurate, superspace crystallography required



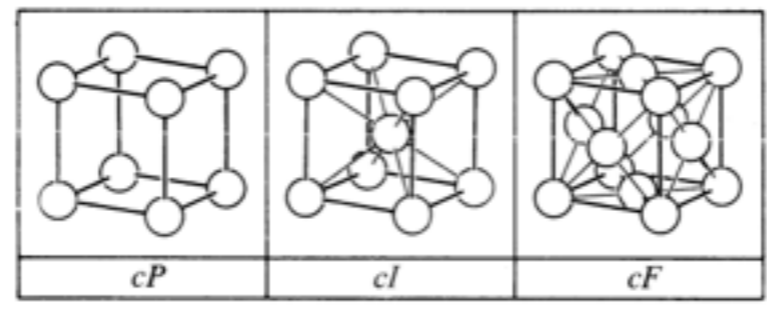
7 unique crystal systems

These are the possible parallelepipeds that tile space

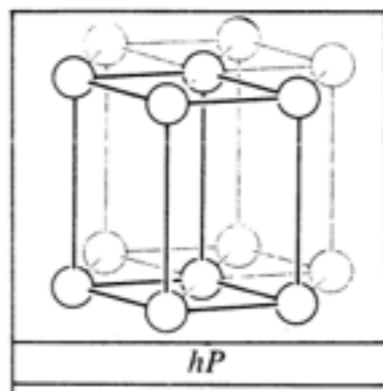


symmetry in three dimensions: crystal systems (7) and Bravais Lattices (14)

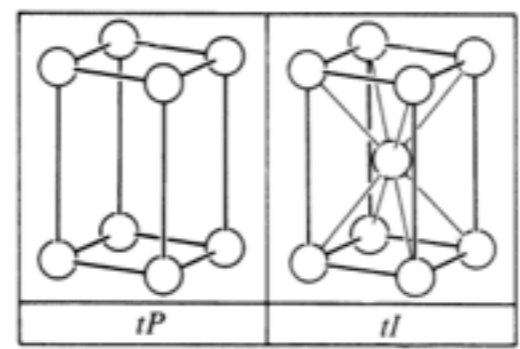
cubic



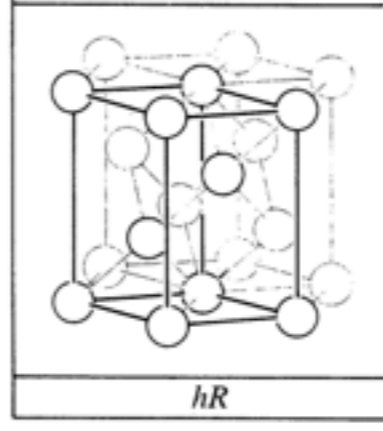
hexagonal



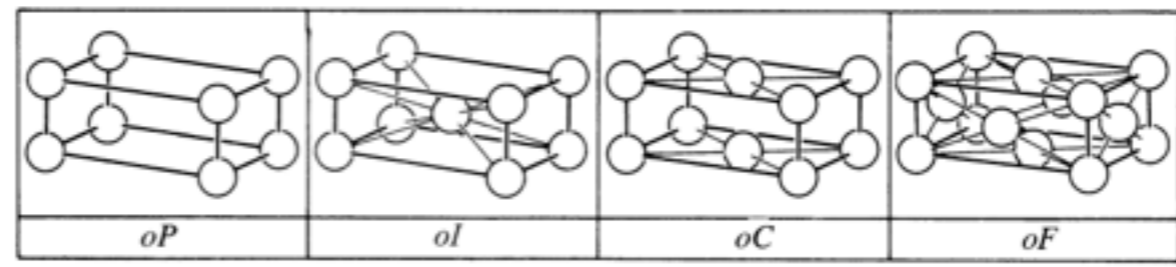
tetragonal



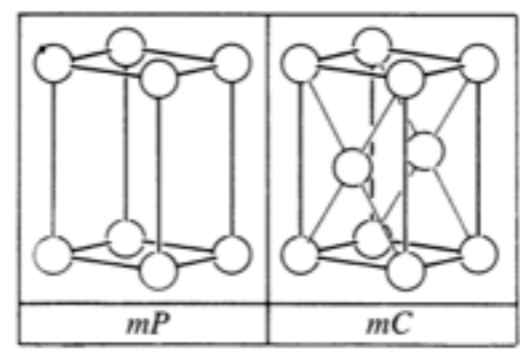
rhombohedral



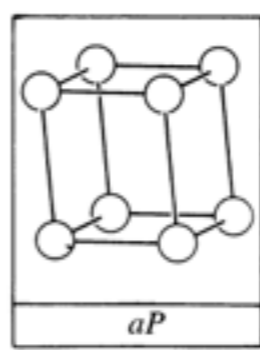
orthorhombic



monoclinic



triclinic



Bravais lattice: combine crystal systems and centering operations

symmetry in three dimensions: 7 crystal systems; 14 Bravais Lattices

Crystal System	Bravais Lattices	Axial lengths and angles
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cubic	P, I, F	$a = b = c; \alpha = \beta = \gamma = 90^\circ$
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tetragonal	P, I	$a = b \neq c; \alpha = \beta = \gamma = 90^\circ$
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orthorhombic	P, I, C, F	$a \neq b \neq c; \alpha = \beta = \gamma = 90^\circ$
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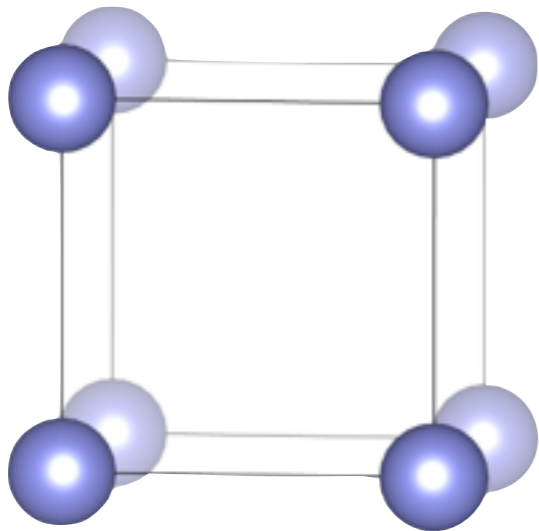
trigonal	P, R	$a = b = c; \alpha = \beta = \gamma \neq 90^\circ$
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hexagonal	P	$a = b \neq c; \alpha = \beta = 90^\circ, \gamma = 120^\circ$
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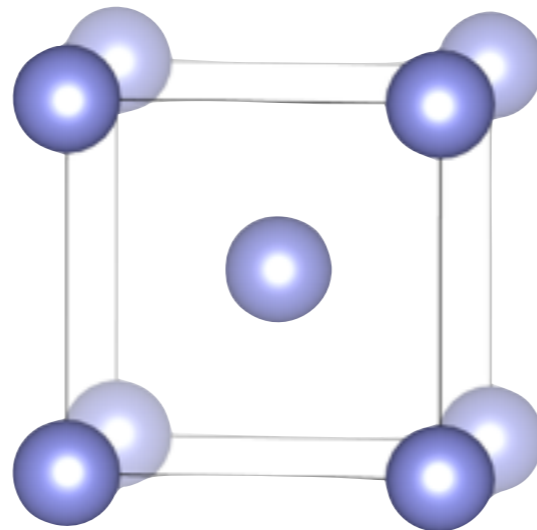
monoclinic	P, C	$a \neq b \neq c; \alpha = \gamma = 90^\circ, \beta \neq 90^\circ$
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triclinic	P	$a \neq b \neq c; [\alpha \neq \beta \neq \gamma] \neq 90^\circ$
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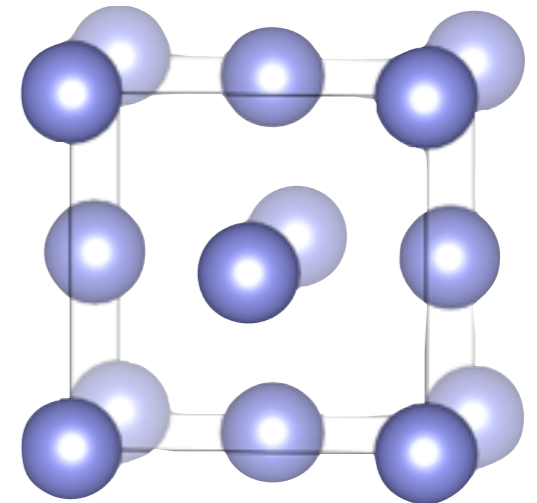
common structures of the elements



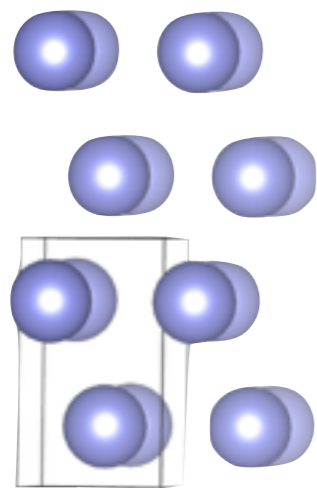
simple cubic: α -Po



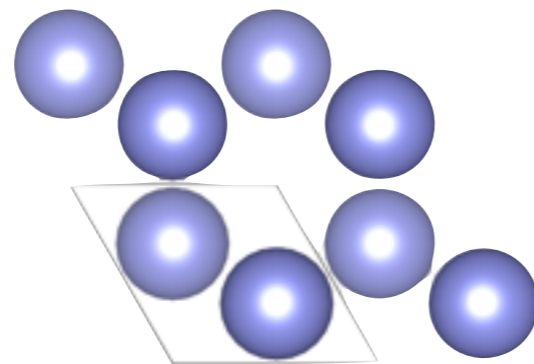
body-centered cubic: α -Fe



face-centered cubic: Cu

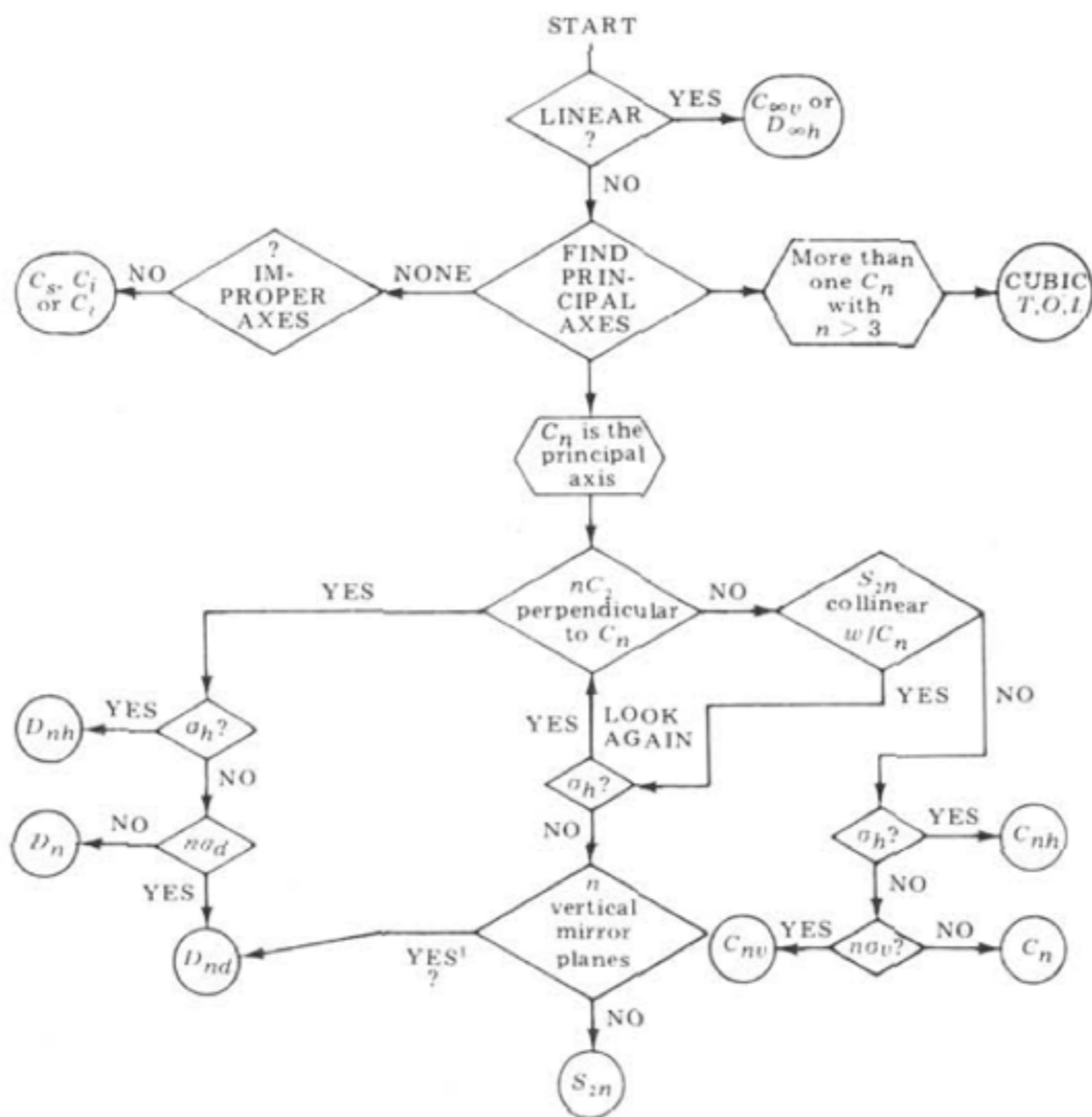


hexagonal-close-packed: Mg



diamond lattice: C, Si

point group identification



"Flow chart for identification of Schoenflies point groups." *J. Chem. Ed.*, 53, 3 (1976): 190.

plane group identification

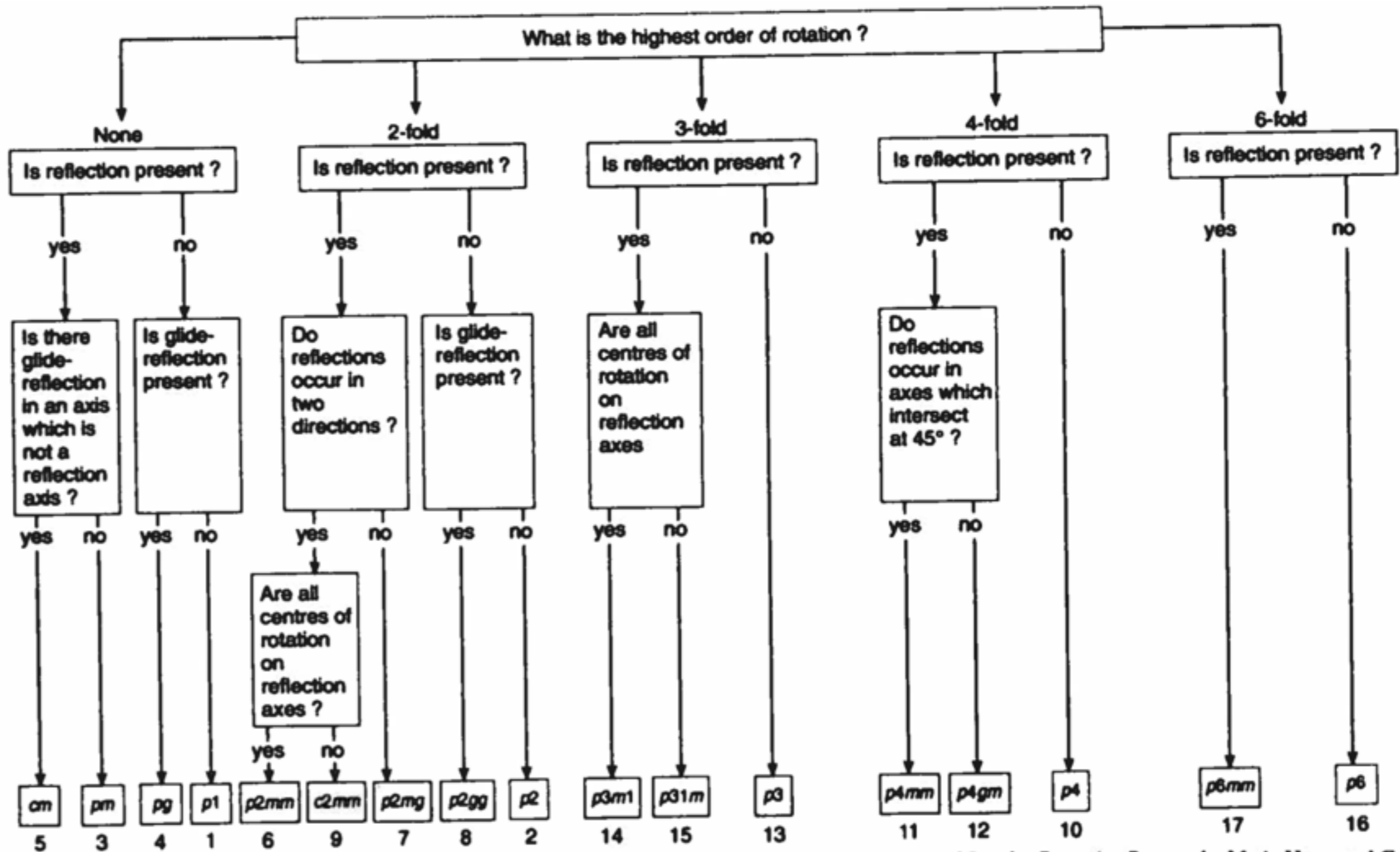


Fig. 2.8. Flow diagram for identifying one of the seventeen plane patterns (redrawn from *The Geometry of Regular Repeating Patterns* by M. A. Hann and G. M. Thomson, the Textile Institute, Manchester, 1992). The numbering is that which is arbitrarily assigned in the International Tables (see Fig 2.6).

space group identification

FINDSYM SUITE HELP

FINDSYM

Version 4.1, May 2013
Harold T. Stokes, Branton J. Campbell, and Dorian M. Hatch, Department of Physics and Astronomy, Brigham Young University, Provo, Utah, 84602, USA, stokesh@byu.edu

Description: Identify the space group of a crystal, given the positions of the atoms in a unit cell.

[Help](#)
[Version History](#)

How to cite FINDSYM: ISOTROPY Software Suite, iso.byu.edu.
Also H. T. Stokes and D. M. Hatch, "FINDSYM: Program for Identifying the Space Group Symmetry of a Crystal." *J. Appl. Cryst.* **38**, 237-238 (2005).

NEW VERSION: Includes magnetic space groups.

Title (copied to output)

Tolerance

Lattice parameters that define a "conventional" unit cell. (alpha is the angle between b and c; beta is the angle between a and c; gamma is the angle between a and b.) Units for a,b,c are arbitrary. Units for alpha,beta,gamma are degrees.

a: b: c:

alpha: beta: gamma:

(use a computer)

<http://stokes.byu.edu/iso/findsym.php>

Let's build a structure

face-centered cubic: Cu

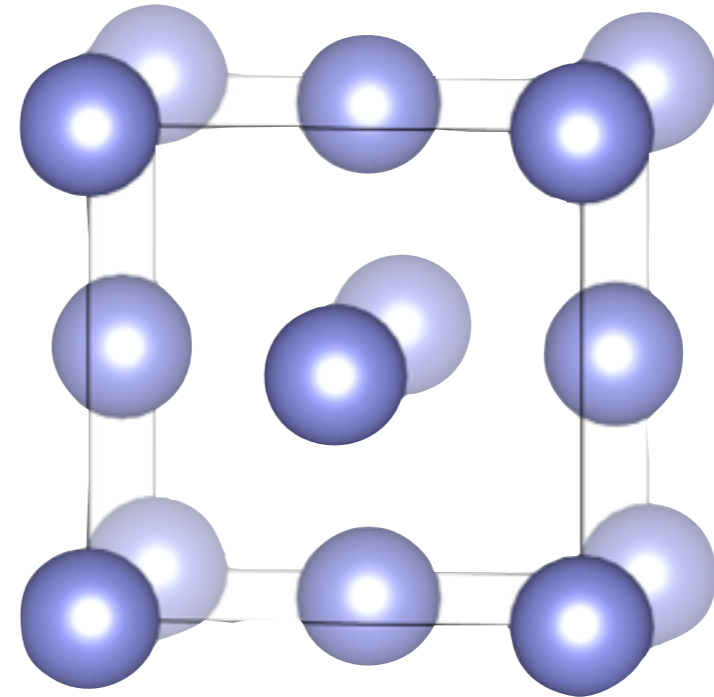
Cubic, $a = 3.60 \text{ \AA}$

Space-group: $225 \quad Fm\bar{3}m$

Atom	x	y	z
------	-----	-----	-----

Cu	0	0	0
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(fractional coordinates)



We only describe the symmetry inequivalent atoms;
let symmetry do the rest.

**How many atoms are in a face-centered cubic,
close-packed unit cell?**

Consulting the Int'l Tables of Crystall.

International Tables for Crystallography (2006). Vol. A, Space group 225, pp. 688–691.

$Fm\bar{3}m$

O_h^5

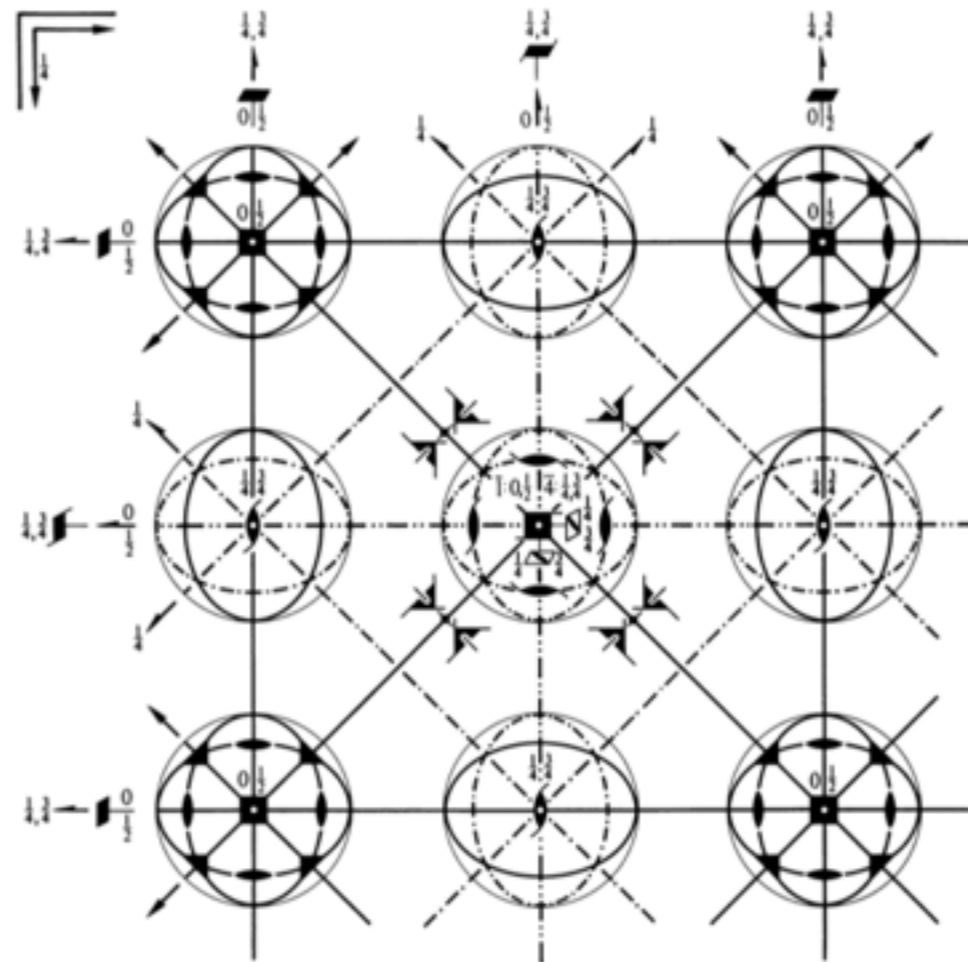
$m\bar{3}m$

Cubic

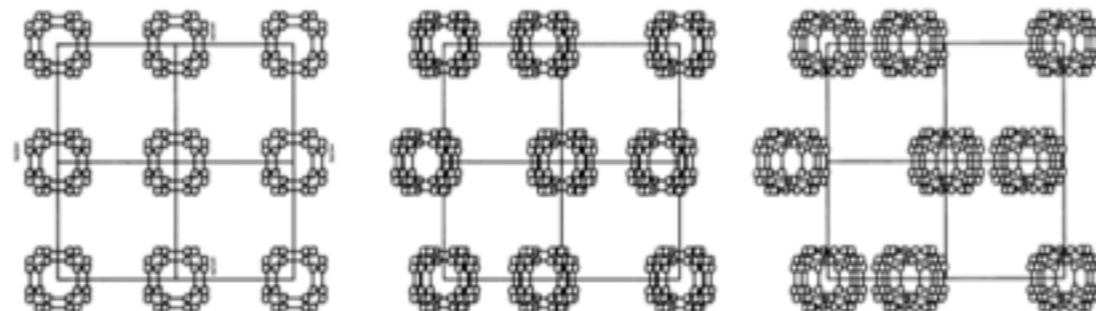
No. 225

$F 4/m \bar{3} 2/m$

Patterson symmetry $Fm\bar{3}m$



Upper left quadrant only



Positions

Multiplicity,
Wyckoff letter,
Site symmetry

Coordinates

$(0,0,0)+$ $(0,\frac{1}{2},\frac{1}{2})+$ $(\frac{1}{2},0,\frac{1}{2})+$ $(\frac{1}{2},\frac{1}{2},0)+$

face
centering
operations

Reflection conditions

h, k, l permutable

General:

$hkl : h+k, h+l, k+l = 2n$

$0kl : k, l = 2n$

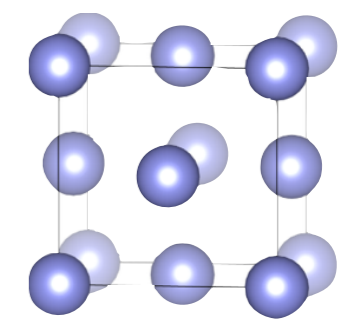
$hhl : h+l = 2n$

$h00 : h = 2n$

192	I	1	(1) x, y, z	(2) \bar{x}, \bar{y}, z	(3) \bar{x}, y, \bar{z}	(4) x, \bar{y}, \bar{z}
			(5) z, x, y	(6) z, \bar{x}, \bar{y}	(7) \bar{z}, \bar{x}, y	(8) \bar{z}, x, \bar{y}
			(9) y, z, x	(10) \bar{y}, z, \bar{x}	(11) y, \bar{z}, \bar{x}	(12) \bar{y}, \bar{z}, x
			(13) y, x, \bar{z}	(14) $\bar{y}, \bar{x}, \bar{z}$	(15) y, \bar{x}, z	(16) \bar{y}, x, z
			(17) x, z, \bar{y}	(18) \bar{x}, z, y	(19) $\bar{x}, \bar{z}, \bar{y}$	(20) x, \bar{z}, y
			(21) z, y, \bar{x}	(22) z, \bar{y}, x	(23) \bar{z}, y, x	(24) $\bar{z}, \bar{y}, \bar{x}$
			(25) $\bar{x}, \bar{y}, \bar{z}$	(26) x, y, \bar{z}	(27) x, \bar{y}, z	(28) \bar{x}, y, z
			(29) $\bar{z}, \bar{x}, \bar{y}$	(30) \bar{z}, x, y	(31) z, x, \bar{y}	(32) z, \bar{x}, y
			(33) $\bar{y}, \bar{z}, \bar{x}$	(34) y, \bar{z}, x	(35) \bar{y}, z, x	(36) y, z, \bar{x}
			(37) \bar{y}, \bar{x}, z	(38) y, x, z	(39) \bar{y}, x, \bar{z}	(40) y, \bar{x}, \bar{z}
			(41) \bar{x}, \bar{z}, y	(42) x, \bar{z}, \bar{y}	(43) x, z, y	(44) \bar{x}, z, \bar{y}
			(45) $\bar{z}, \bar{y}, \bar{x}$	(46) \bar{z}, y, \bar{x}	(47) z, \bar{y}, \bar{x}	(48) z, y, x

96	k	$. . m$	x, x, z	\bar{x}, \bar{x}, z	\bar{x}, x, \bar{z}	x, \bar{x}, \bar{z}	z, x, x	z, \bar{x}, \bar{x}
			\bar{z}, \bar{x}, x	\bar{z}, x, \bar{x}	x, z, x	\bar{x}, z, \bar{x}	x, \bar{z}, \bar{x}	\bar{x}, \bar{z}, x
			x, x, \bar{z}	$\bar{x}, \bar{x}, \bar{z}$	x, \bar{x}, z	\bar{x}, x, z	x, z, \bar{x}	\bar{x}, z, x
			$\bar{x}, \bar{z}, \bar{x}$	x, \bar{z}, x	z, x, \bar{x}	z, \bar{x}, x	\bar{z}, x, x	$\bar{z}, \bar{x}, \bar{x}$
96	j	$m . .$	$0, y, z$	$0, \bar{y}, z$	$0, y, \bar{z}$	$0, \bar{y}, \bar{z}$	$z, 0, y$	$z, 0, \bar{y}$
			$\bar{z}, 0, y$	$\bar{z}, 0, \bar{y}$	$y, z, 0$	$\bar{y}, z, 0$	$y, \bar{z}, 0$	$\bar{y}, \bar{z}, 0$
			$y, 0, \bar{z}$	$\bar{y}, 0, \bar{z}$	$y, 0, z$	$\bar{y}, 0, z$	$0, z, \bar{y}$	$0, z, y$
			$0, \bar{z}, \bar{y}$	$0, \bar{z}, y$	$z, y, 0$	$z, \bar{y}, 0$	$\bar{z}, y, 0$	$\bar{z}, \bar{y}, 0$
48	i	$m . m 2$	$\frac{1}{2}, y, y$	$\frac{1}{2}, \bar{y}, y$	$\frac{1}{2}, y, \bar{y}$	$\frac{1}{2}, \bar{y}, \bar{y}$	$y, \frac{1}{2}, y$	$y, \frac{1}{2}, \bar{y}$
			$\bar{y}, \frac{1}{2}, y$	$\bar{y}, \frac{1}{2}, \bar{y}$	$y, y, \frac{1}{2}$	$\bar{y}, y, \frac{1}{2}$	$y, \bar{y}, \frac{1}{2}$	$\bar{y}, \bar{y}, \frac{1}{2}$
48	h	$m . m 2$	$0, y, y$	$0, \bar{y}, y$	$0, y, \bar{y}$	$0, \bar{y}, \bar{y}$	$y, 0, y$	$y, 0, \bar{y}$
			$\bar{y}, 0, y$	$\bar{y}, 0, \bar{y}$	$y, y, 0$	$\bar{y}, y, 0$	$y, \bar{y}, 0$	$\bar{y}, \bar{y}, 0$
48	g	$2 . mm$	$x, \frac{1}{2}, \frac{1}{2}$	$\bar{x}, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, x, \frac{1}{2}$	$\frac{1}{2}, \bar{x}, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, x$	$\frac{1}{2}, \frac{1}{2}, \bar{x}$
			$\frac{1}{2}, x, \frac{1}{2}$	$\frac{1}{2}, \bar{x}, \frac{1}{2}$	$x, \frac{1}{2}, \frac{1}{2}$	$\bar{x}, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, \bar{x}$	$\frac{1}{2}, \frac{1}{2}, x$
32	f	$. 3 m$	x, x, x	\bar{x}, \bar{x}, x	\bar{x}, x, \bar{x}	x, \bar{x}, \bar{x}		
			x, x, \bar{x}	$\bar{x}, \bar{x}, \bar{x}$	x, \bar{x}, x	\bar{x}, x, \bar{x}		
24	e	$4 m . m$	$x, 0, 0$	$\bar{x}, 0, 0$	$0, x, 0$	$0, \bar{x}, 0$	$0, 0, x$	$0, 0, \bar{x}$
24	d	$m . mm$	$0, \frac{1}{4}, \frac{1}{4}$	$0, \frac{3}{4}, \frac{1}{4}$	$\frac{1}{4}, 0, \frac{1}{4}$	$\frac{1}{4}, 0, \frac{3}{4}$	$\frac{1}{4}, \frac{1}{4}, 0$	$\frac{3}{4}, \frac{1}{4}, 0$
8	c	$\bar{4} 3 m$	$\frac{1}{4}, \frac{1}{4}, \frac{1}{4}$	$\frac{1}{4}, \frac{1}{4}, \frac{3}{4}$				
4	b	$m \bar{3} m$	$\frac{1}{2}, \frac{1}{2}, \frac{1}{2}$					
4	a	$m \bar{3} m$	$0, 0, 0$					

our 1 symm. inequiv. atom



Symmetry
inequivalent atoms:

Atom	x	y	z
Cu	0	0	0

Multiplicity: 4
(face centering)

All atoms:

Atom	x	y	z
Cu1	0	0	0
Cu2	0	0.5	0.5
Cu3	0.5	0	0.5
Cu4	0.5	0.5	0

Now we have
all 4 atoms

Let's build another structure: a perovskite!

General formula: ABX_3

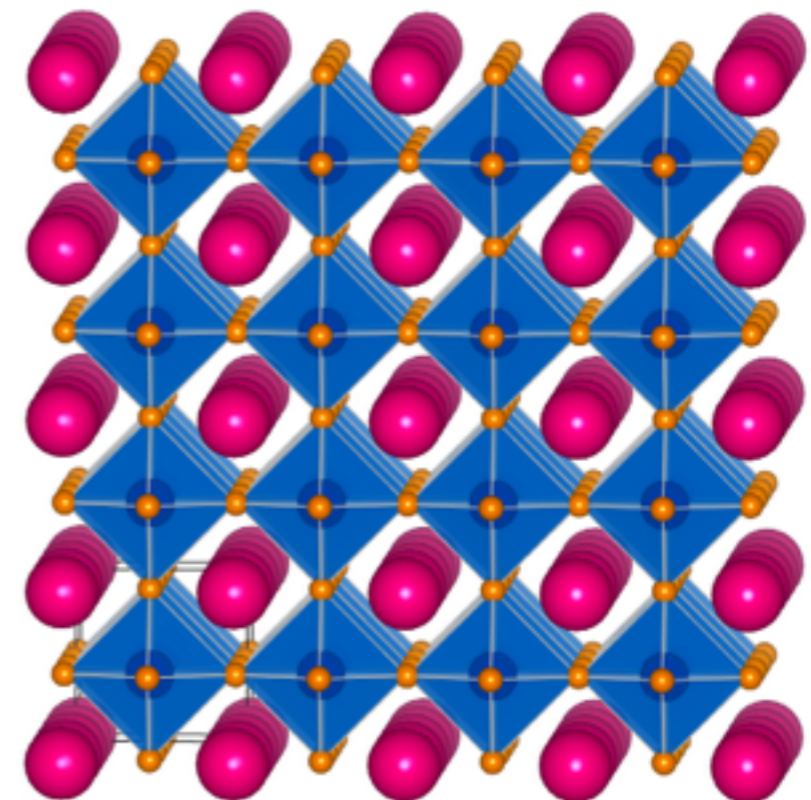
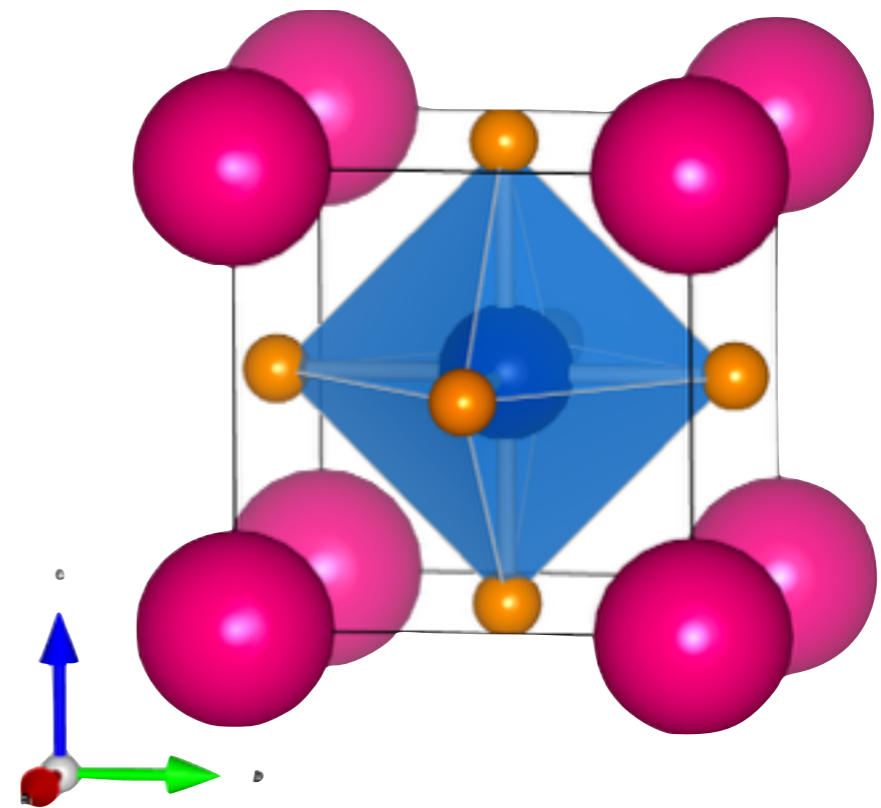
Primitive cubic: SrTiO_3

Cubic, $a = 4.00 \text{ \AA}$

Space-group: 221 $Pm\bar{3}m$

Atom	x	y	z
Sr	0	0	0
Ti	0.5	0.5	0.5
O	0.5	0.5	0

Corner-linked TiO_6 octahedra



Consulting the Int'l Tables of Crystall.

International Tables for Crystallography (2006). Vol. A, Space group 221, pp. 672–674.

$Pm\bar{3}m$

O_h^1

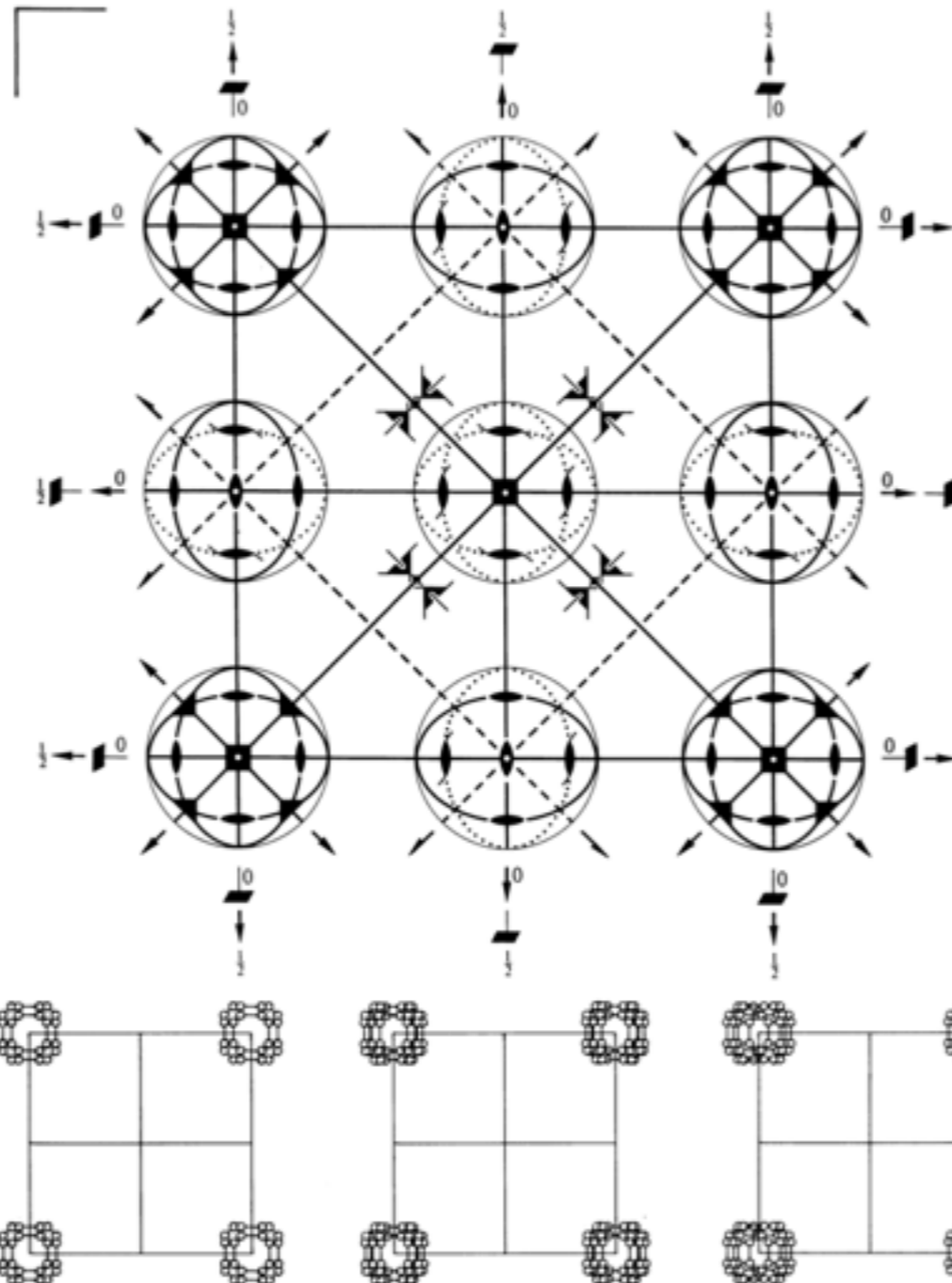
$m\bar{3}m$

Cubic

No. 221

$P 4/m \bar{3} 2/m$

Patterson symmetry $Pm\bar{3}m$

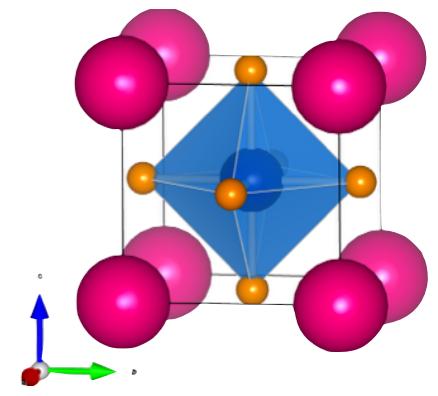


48	<i>n</i>	1	(1) x, y, z (5) z, x, y (9) y, z, x (13) y, x, \bar{z} (17) x, z, \bar{y} (21) z, y, \bar{x} (25) $\bar{x}, \bar{y}, \bar{z}$ (29) $\bar{z}, \bar{x}, \bar{y}$ (33) $\bar{y}, \bar{z}, \bar{x}$ (37) \bar{y}, \bar{x}, z (41) \bar{x}, \bar{z}, y (45) \bar{z}, \bar{y}, x	(2) \bar{x}, \bar{y}, z (6) z, \bar{x}, \bar{y} (10) \bar{y}, z, \bar{x} (14) $\bar{y}, \bar{x}, \bar{z}$ (18) \bar{x}, z, y (22) z, \bar{y}, x (26) x, y, \bar{z} (30) \bar{z}, x, y (34) y, \bar{z}, x (38) y, x, z (42) x, \bar{z}, \bar{y} (46) \bar{z}, y, \bar{x}	(3) \bar{x}, y, \bar{z} (7) \bar{z}, \bar{x}, y (11) y, \bar{z}, \bar{x} (15) y, \bar{x}, z (19) $\bar{x}, \bar{z}, \bar{y}$ (23) \bar{z}, y, x (27) x, \bar{y}, z (31) z, x, \bar{y} (35) \bar{y}, z, x (39) \bar{y}, x, \bar{z} (43) x, z, y (47) z, \bar{y}, \bar{x}	(4) x, \bar{y}, \bar{z} (8) \bar{z}, x, \bar{y} (12) \bar{y}, \bar{z}, x (16) \bar{y}, x, z (20) x, \bar{z}, y (24) $\bar{z}, \bar{y}, \bar{x}$ (28) \bar{x}, y, z (32) z, \bar{x}, y (36) y, z, \bar{x} (40) y, \bar{x}, \bar{z} (44) \bar{x}, z, \bar{y} (48) z, y, x	no conditions
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Special: no extra conditions

24	<i>m</i>	$. . m$	x, x, z \bar{z}, \bar{x}, x x, x, \bar{z} $\bar{x}, \bar{z}, \bar{x}$	\bar{x}, \bar{x}, z \bar{z}, x, \bar{x} $\bar{x}, \bar{x}, \bar{z}$ x, \bar{z}, x	\bar{x}, x, \bar{z} x, z, x x, \bar{x}, z z, x, \bar{x}	x, \bar{x}, \bar{z} \bar{x}, z, \bar{x} \bar{x}, x, z z, \bar{x}, x	z, x, x x, \bar{z}, \bar{x} x, z, \bar{x} \bar{z}, x, x	z, \bar{x}, \bar{x} \bar{x}, \bar{z}, x \bar{x}, z, x $\bar{z}, \bar{x}, \bar{x}$
24	<i>l</i>	$m . .$	$\frac{1}{2}, y, z$ $\bar{z}, \frac{1}{2}, y$ $y, \frac{1}{2}, \bar{z}$ $\frac{1}{2}, \bar{z}, \bar{y}$	$\frac{1}{2}, \bar{y}, z$ $\bar{z}, \frac{1}{2}, \bar{y}$ $\bar{y}, \frac{1}{2}, \bar{z}$ $\frac{1}{2}, \bar{z}, y$	$\frac{1}{2}, y, \bar{z}$ $y, z, \frac{1}{2}$ $y, \frac{1}{2}, z$ $z, y, \frac{1}{2}$	$\frac{1}{2}, \bar{y}, \bar{z}$ $\bar{y}, z, \frac{1}{2}$ $\bar{y}, \frac{1}{2}, z$ $z, \bar{y}, \frac{1}{2}$	$z, \frac{1}{2}, y$ $y, \bar{z}, \frac{1}{2}$ $\frac{1}{2}, z, \bar{y}$ $\bar{z}, y, \frac{1}{2}$	$z, \frac{1}{2}, \bar{y}$ $\bar{y}, \bar{z}, \frac{1}{2}$ $\frac{1}{2}, z, y$ $\bar{z}, \bar{y}, \frac{1}{2}$
24	<i>k</i>	$m . .$	$0, y, z$ $\bar{z}, 0, y$ $y, 0, \bar{z}$ $0, \bar{z}, \bar{y}$	$0, \bar{y}, z$ $\bar{z}, 0, \bar{y}$ $\bar{y}, 0, \bar{z}$ $0, \bar{z}, y$	$0, y, \bar{z}$ $y, z, 0$ $y, 0, z$ $z, y, 0$	$0, \bar{y}, \bar{z}$ $\bar{y}, z, 0$ $\bar{y}, 0, z$ $z, \bar{y}, 0$	$z, 0, y$ $y, \bar{z}, 0$ $0, z, \bar{y}$ $\bar{z}, y, 0$	$z, 0, \bar{y}$ $\bar{y}, \bar{z}, 0$ $0, z, y$ $\bar{z}, \bar{y}, 0$
12	<i>j</i>	$m . m 2$	$\frac{1}{2}, y, y$ $\bar{y}, \frac{1}{2}, y$	$\frac{1}{2}, \bar{y}, y$ $\bar{y}, \frac{1}{2}, \bar{y}$	$\frac{1}{2}, y, \bar{y}$ $y, y, \frac{1}{2}$	$\frac{1}{2}, \bar{y}, \bar{y}$ $\bar{y}, y, \frac{1}{2}$	$y, \frac{1}{2}, y$ $y, \bar{y}, \frac{1}{2}$	$y, \frac{1}{2}, \bar{y}$ $\bar{y}, \bar{y}, \frac{1}{2}$
12	<i>i</i>	$m . m 2$	$0, y, y$ $\bar{y}, 0, y$	$0, \bar{y}, y$ $\bar{y}, 0, \bar{y}$	$0, y, \bar{y}$ $y, y, 0$	$0, \bar{y}, \bar{y}$ $\bar{y}, y, 0$	$y, 0, y$ $y, \bar{y}, 0$	$y, 0, \bar{y}$ $\bar{y}, \bar{y}, 0$
12	<i>h</i>	$m m 2 . .$	$x, \frac{1}{2}, 0$ $\frac{1}{2}, x, 0$	$\bar{x}, \frac{1}{2}, 0$ $\frac{1}{2}, \bar{x}, 0$	$0, x, \frac{1}{2}$ $x, 0, \frac{1}{2}$	$0, \bar{x}, \frac{1}{2}$ $\bar{x}, 0, \frac{1}{2}$	$\frac{1}{2}, 0, x$ $0, \frac{1}{2}, \bar{x}$	$\frac{1}{2}, 0, \bar{x}$ $0, \frac{1}{2}, x$
8	<i>g</i>	$. 3 m$	x, x, x x, x, \bar{x}	\bar{x}, \bar{x}, x $\bar{x}, \bar{x}, \bar{x}$	\bar{x}, x, \bar{x} x, \bar{x}, x	x, \bar{x}, \bar{x} \bar{x}, x, x		
6	<i>f</i>	$4 m . m$	$x, \frac{1}{2}, \frac{1}{2}$	$\bar{x}, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, x, \frac{1}{2}$	$\frac{1}{2}, \bar{x}, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, x$	$\frac{1}{2}, \frac{1}{2}, \bar{x}$
6	<i>e</i>	$4 m . m$	$x, 0, 0$	$\bar{x}, 0, 0$	$0, x, 0$	$0, \bar{x}, 0$	$0, 0, x$	$0, 0, \bar{x}$
3	<i>d</i>	$4/m m . m$	$\frac{1}{2}, 0, 0$	$0, \frac{1}{2}, 0$	$0, 0, \frac{1}{2}$			
3	<i>c</i>	$4/m m . m$	$0, \frac{1}{2}, \frac{1}{2}$	$\frac{1}{2}, 0, \frac{1}{2}$	$\frac{1}{2}, \frac{1}{2}, 0$			
1	<i>b</i>	$m \bar{3} m$	$\frac{1}{2}, \frac{1}{2}, \frac{1}{2}$					Ti(1)
1	<i>a</i>	$m \bar{3} m$	$0, 0, 0$					Sr(1)

O(1), O(2), O(3)



Symmetry
inequivalent atoms:

Atom	x	y	z
Sr	0	0	0
Ti	0.5	0.5	0.5
O	0.5	0.5	0

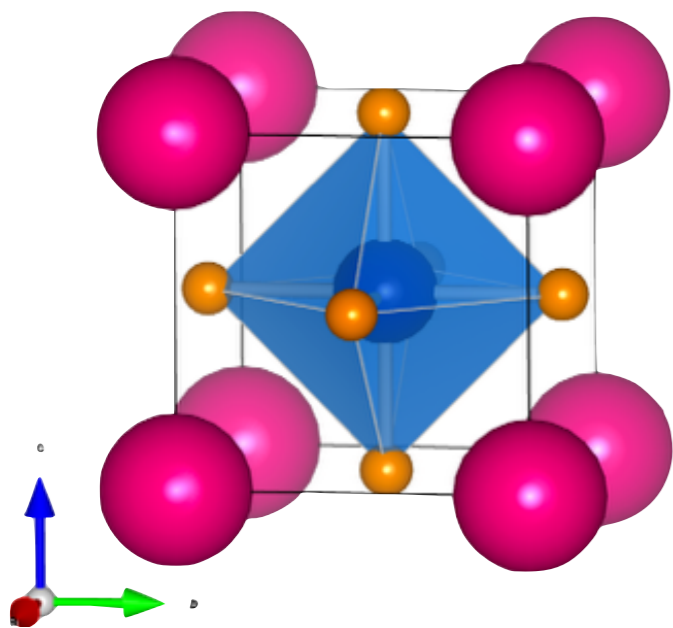
No centering operations

All atoms:

Atom	x	y	z
Sr1	0	0	0
Ti1	0.5	0.5	0.5
O1	0.5	0.5	0
O2	0.5	0	0.5
O3	0.5	0.5	0

tiling with abandon...

Atom	x	y	z
Sr1	0	0	0
<i>Ti1</i>	<i>0.5</i>	<i>0.5</i>	<i>0.5</i>
O1	0.5	0.5	0
<i>O2</i>	<i>0.5</i>	<i>0</i>	<i>0.5</i>
O3	0.5	0.5	0



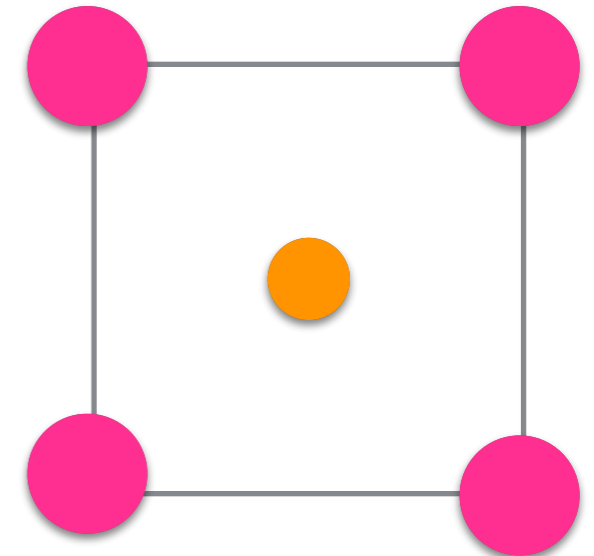
sanity check,
coordination
numbers:

O: 2

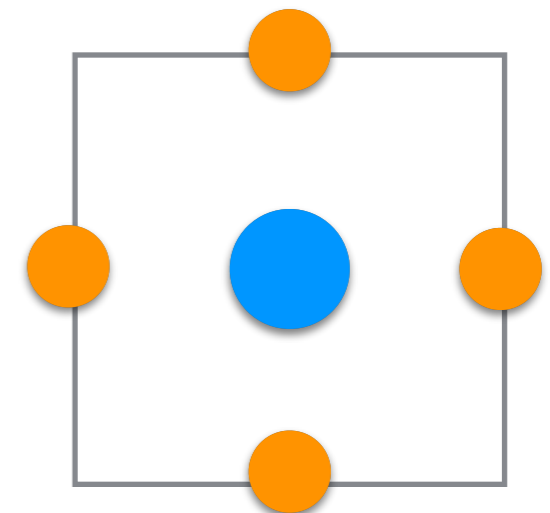
Ti: 6

Sr: 12

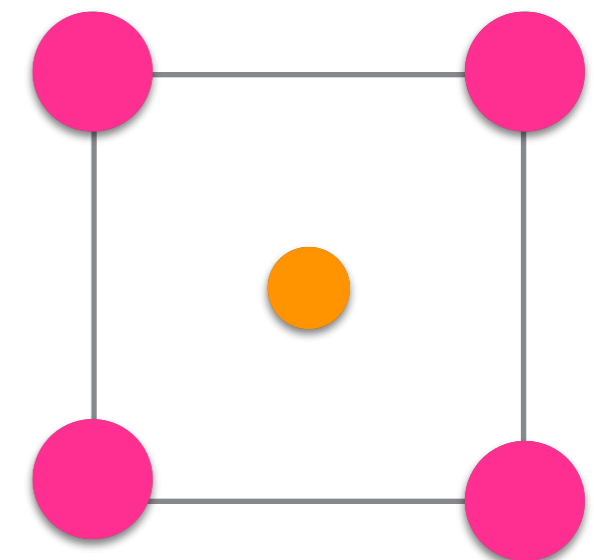
$z = 0$



$z = 0.5$



$z = 1$



entering into reciprocal space

Fourier transforms

mathematical operations that identify repeating patterns

Key:



f

Original function
(direct space)

f[^]

Fourier transform
(reciprocal space)

Fourier transforms

mathematical operations that identify repeating patterns

Key:

$$f(\mathbf{r}) = \int_{-\infty}^{\infty} \hat{f}(\mathbf{q}) e^{2\pi i \mathbf{q} \cdot \mathbf{r}} d\mathbf{q}$$

Inverse transform
(direct space)

$$\hat{f}(\mathbf{q}) = \int_{-\infty}^{\infty} f(\mathbf{r}) e^{-2\pi i \mathbf{q} \cdot \mathbf{r}} d\mathbf{r}$$

Fourier transform
(reciprocal space)



Crystals have repeating patterns

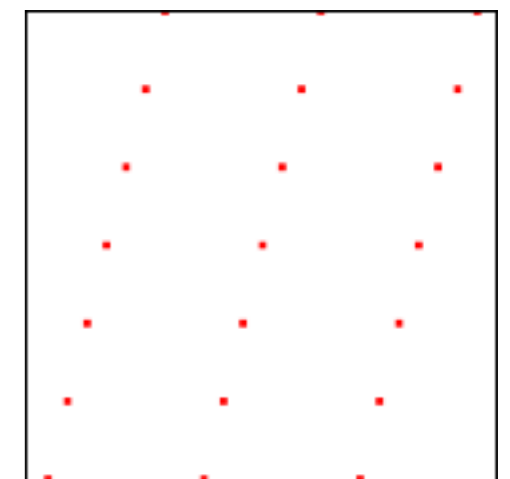
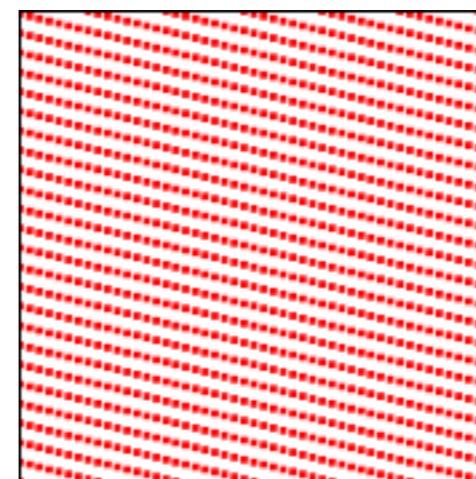
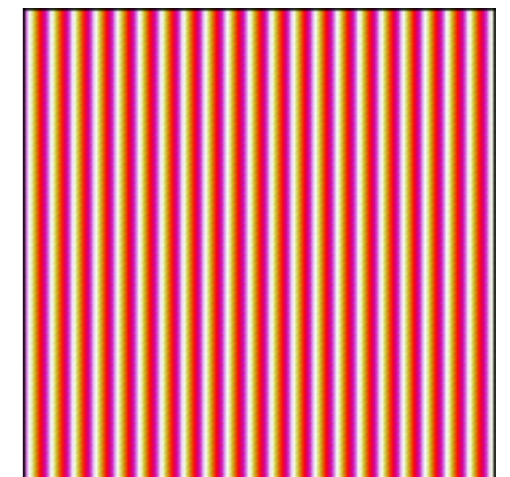
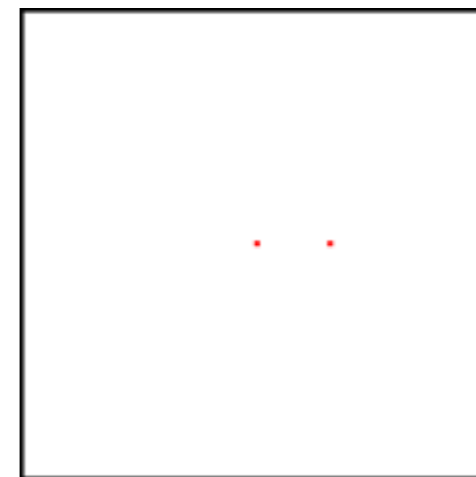
- For each plane of atoms, you have a characteristic periodicity (or frequency).
- What are the units? *inverse length*
- **Reciprocal space**

Fourier transforms of periodic structures: diffraction

We can do this with two-dimensional objects. Two slits will produce a grating interference pattern.

A grating will produce a lattice of points.

object \longrightarrow Fourier transform



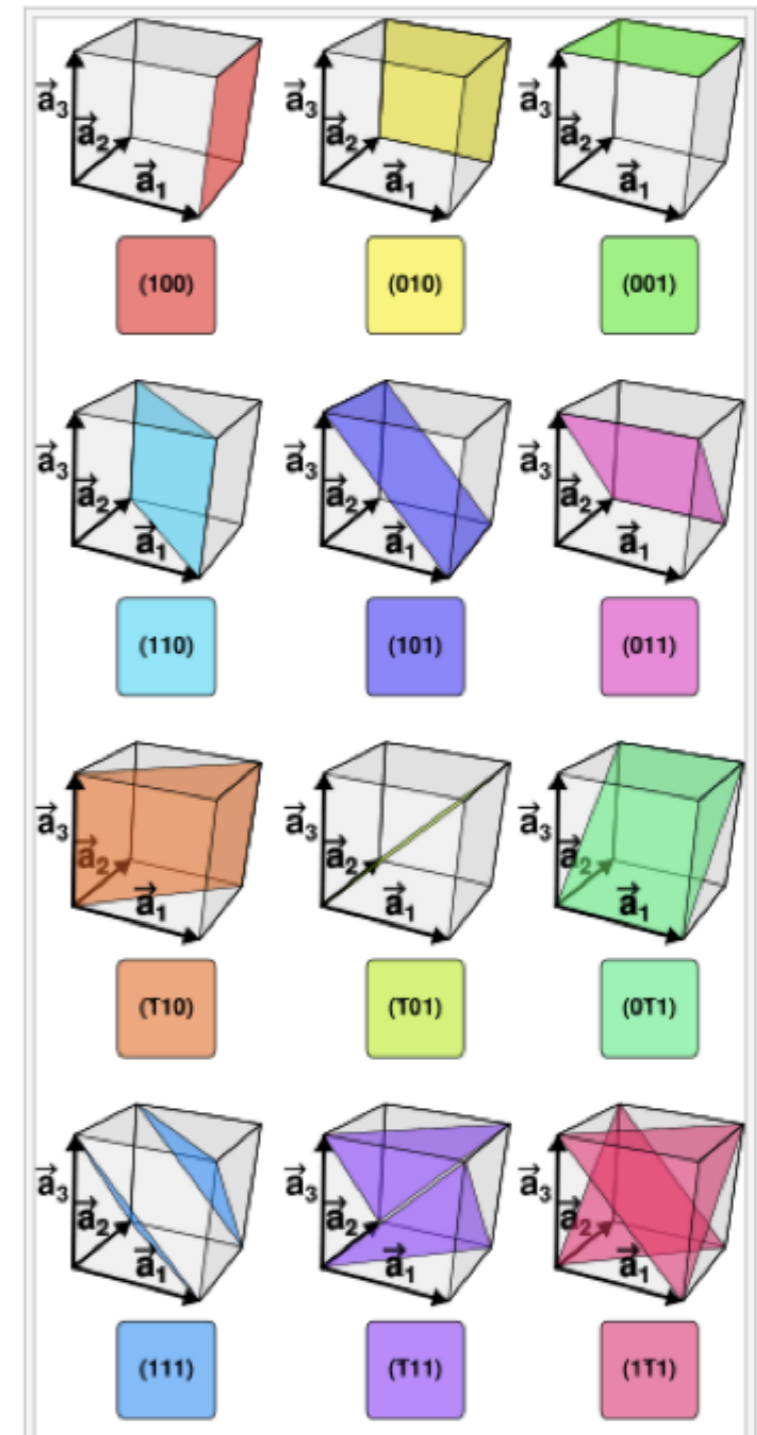
Images taken from Kevin Cowtan's Picture Book of Fourier Transforms:

<http://www.yesbl.york.ac.uk/~cowtan/fourier/fourier.html>

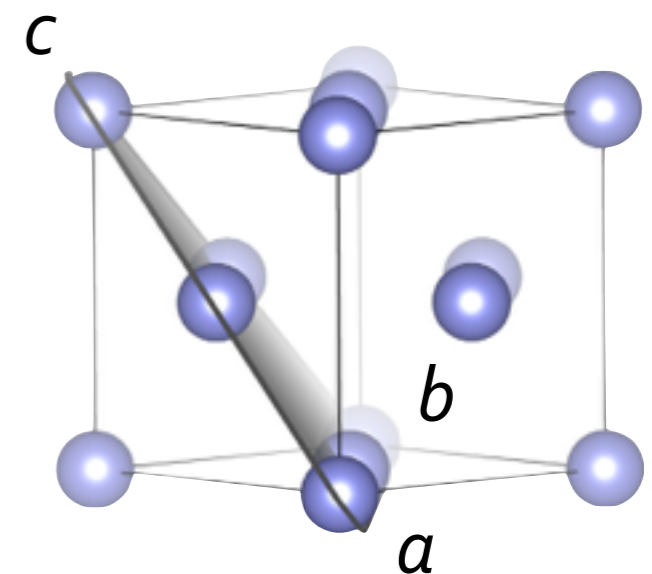
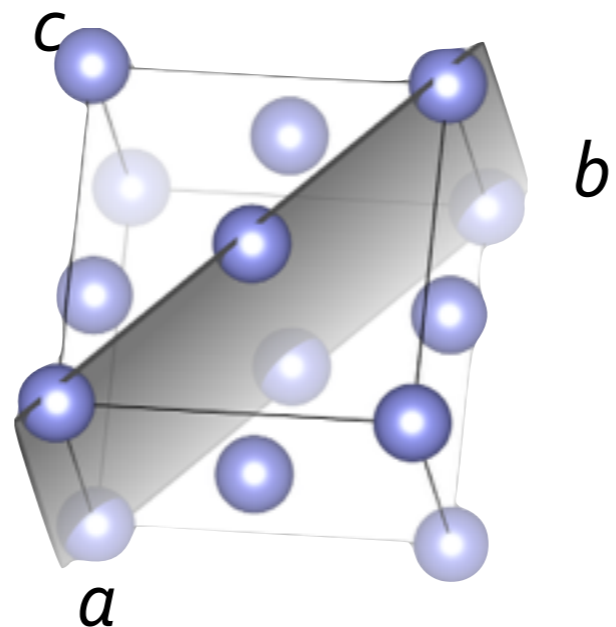
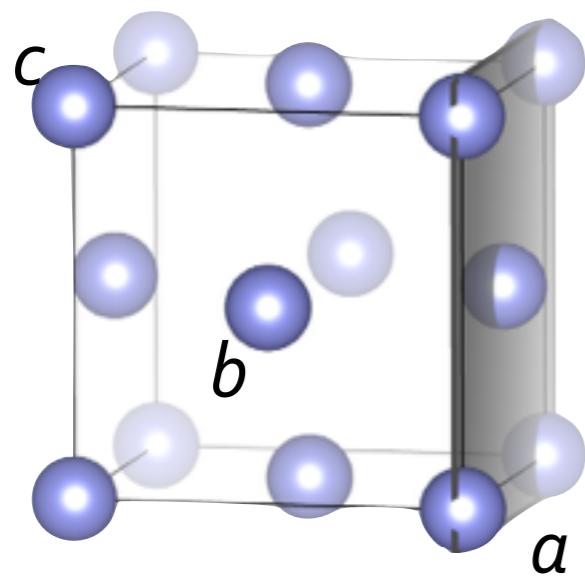
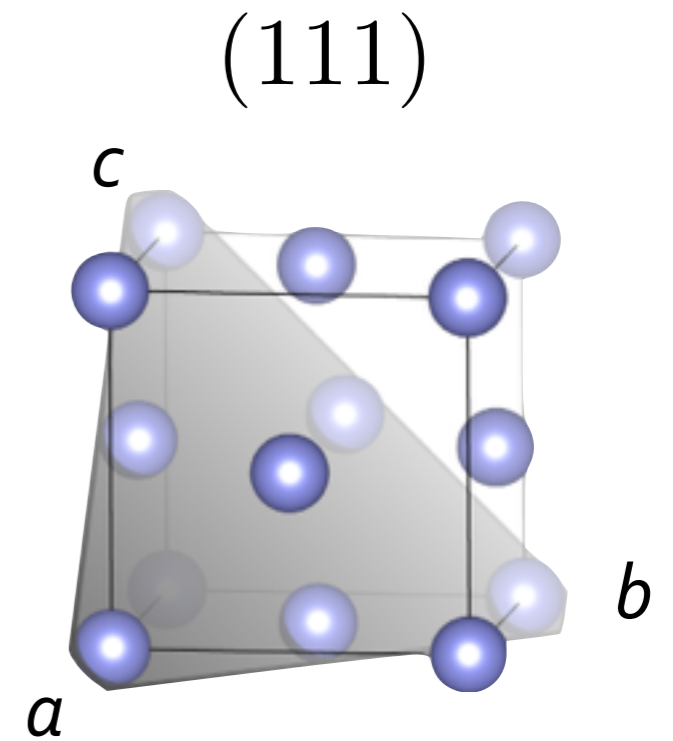
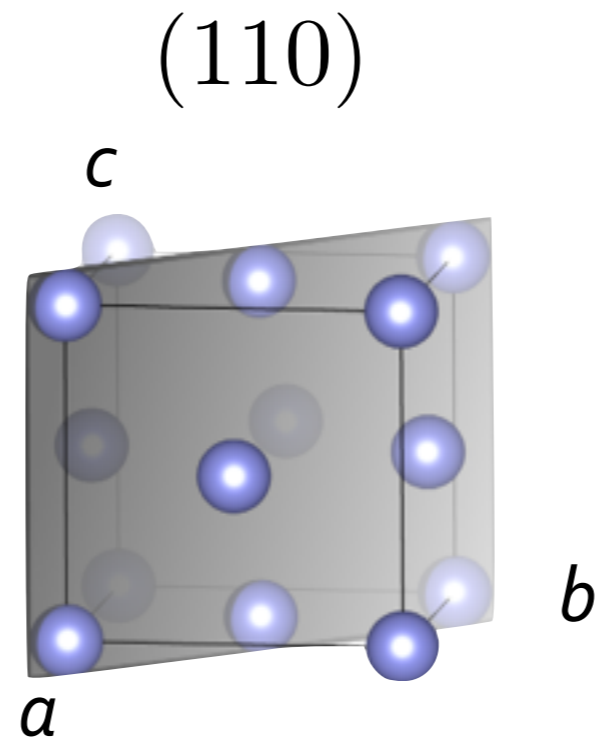
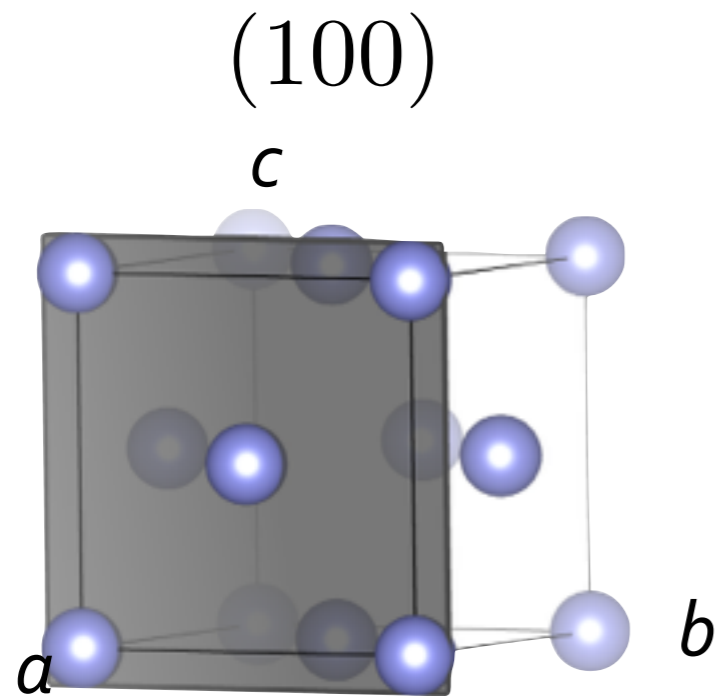
Visit for more (and interactive) examples.

Miller indices: labels for sets of planes

- the indices: (hkl)
 - h : $(\text{intercept with } a)^{-1}$
 - k : $(\text{intercept with } b)^{-1}$
 - l : $(\text{intercept with } c)^{-1}$
- Convention dictates that we use integers; therefore, multiply if needed.



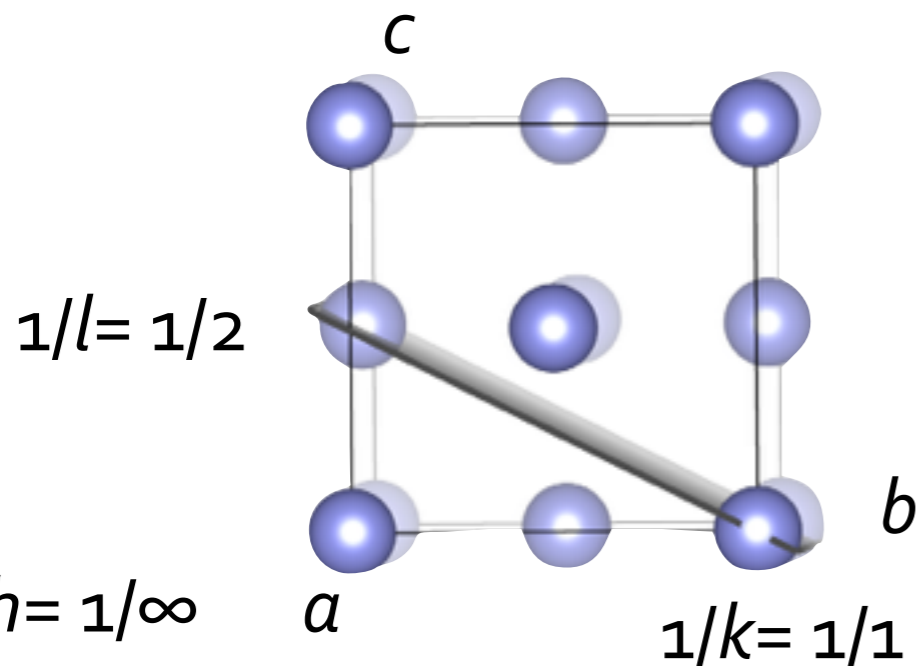
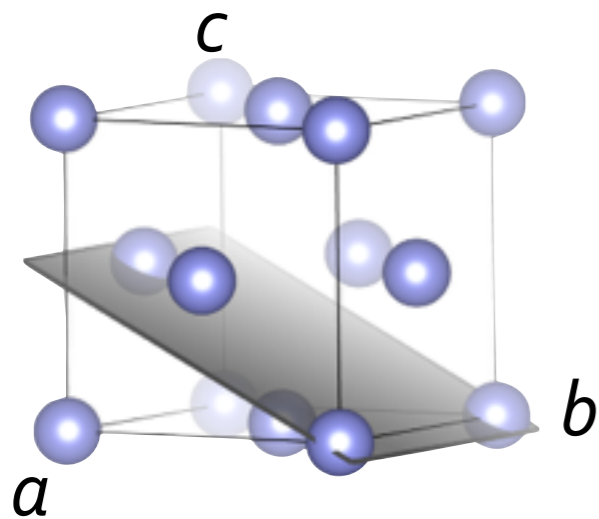
Examples of Miller Indices



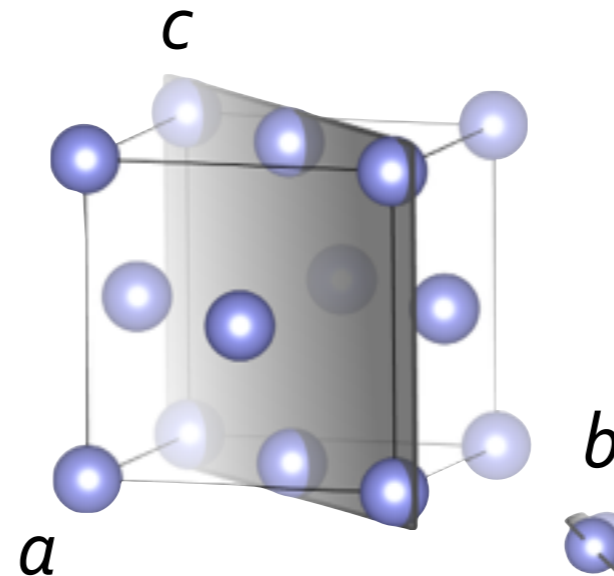
for no intercept: we take the intercept as ∞ ; the index is then zero

More examples of Miller Indices

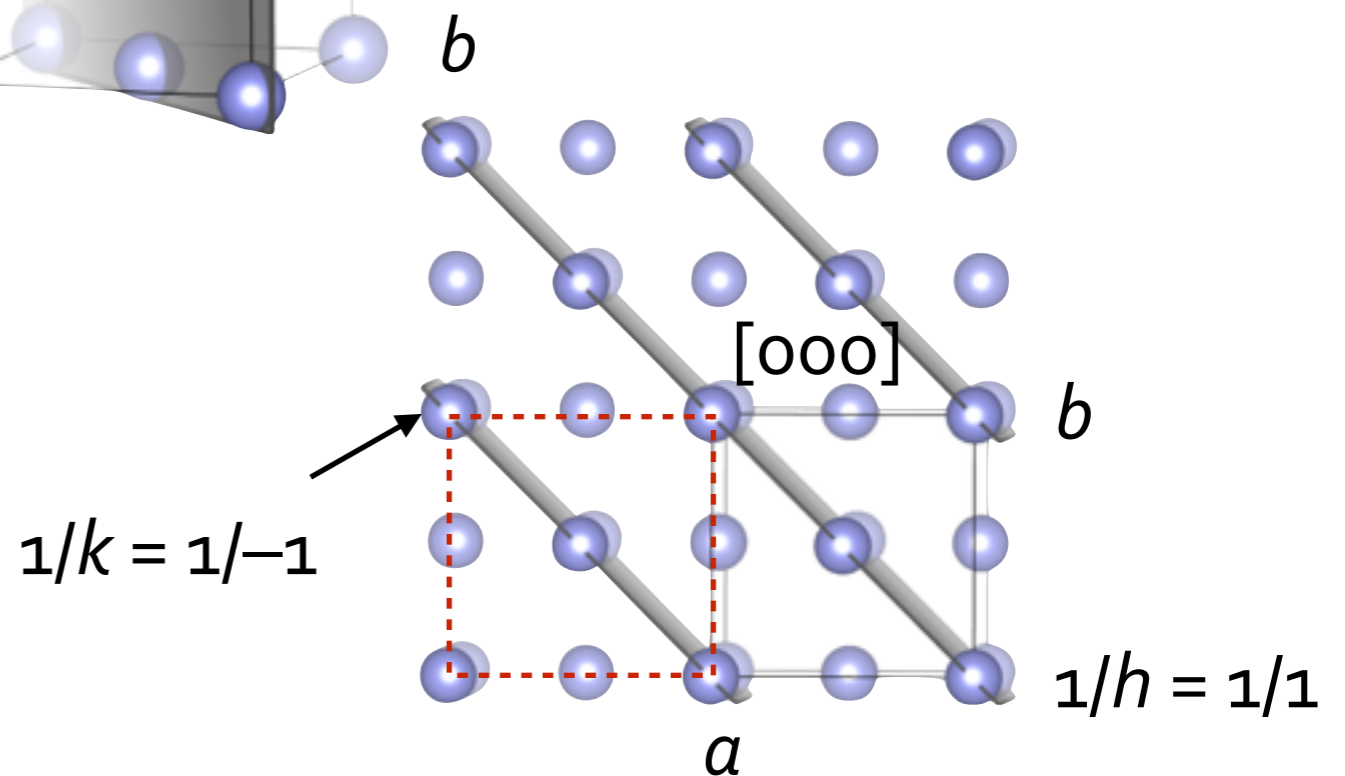
(012)



(1 $\bar{1}$ 0)



negative index?
tile with abandon



For animations, see:

http://www.doitpoms.ac.uk/tlplib/miller_indices/lattice_index.php

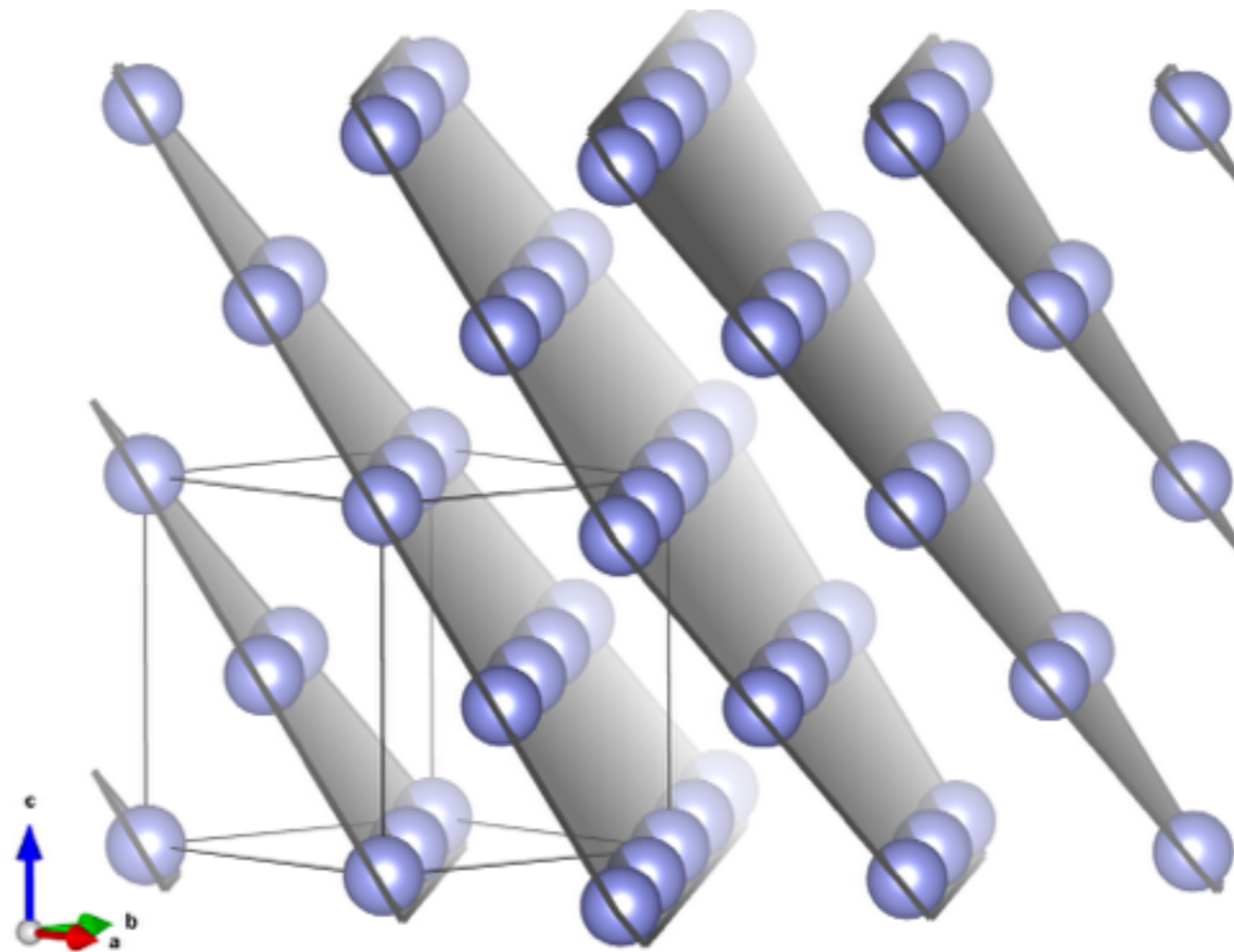
index notation (for completeness)

- (hkl)
 - indexes a *specific* lattice plane
- $\{hkl\}$
 - indexes all *symmetry-related* lattice planes
 - In a cubic system,
 $\{100\} \ni (100), (010), (001),$
 $(\bar{1}00), (0\bar{1}0), (00\bar{1})$

- $[uvw]$
 - indexes a *specific* lattice vector
- $\langle uvw \rangle$
 - indexes all *symmetry-related* lattice vectors
 - In a cubic system,
 $\langle 110 \rangle \ni [110], [011], [101],$
 $[\bar{1}10], [1\bar{1}0], [\bar{1}\bar{1}0],$
 $[0\bar{1}1], [01\bar{1}], [0\bar{1}\bar{1}],$
 $[\bar{1}01], [10\bar{1}], [\bar{1}0\bar{1}]$

Miller indices index all sets of planes

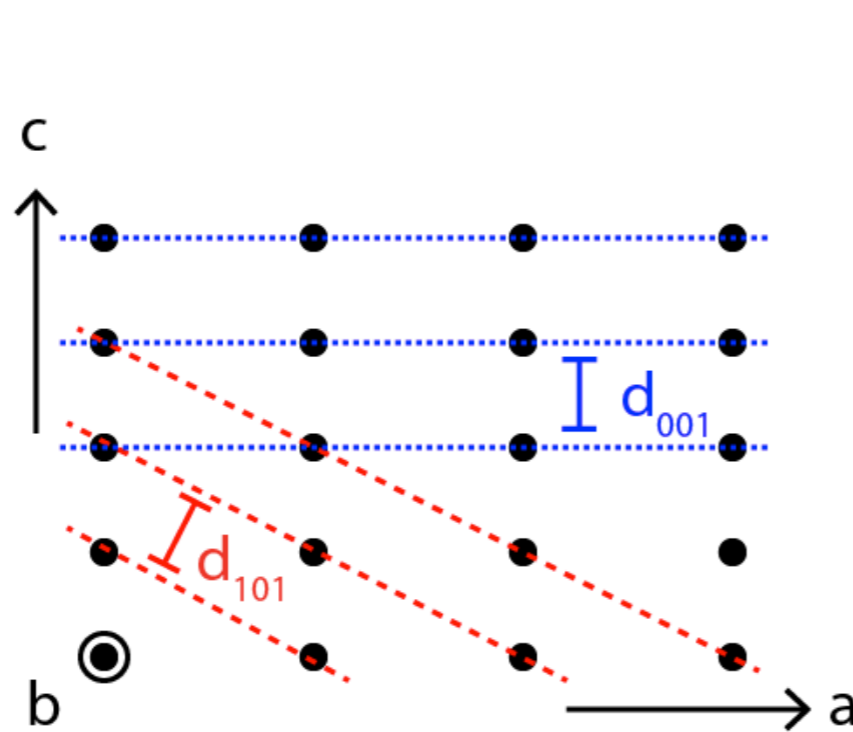
(111) = all of the planes drawn



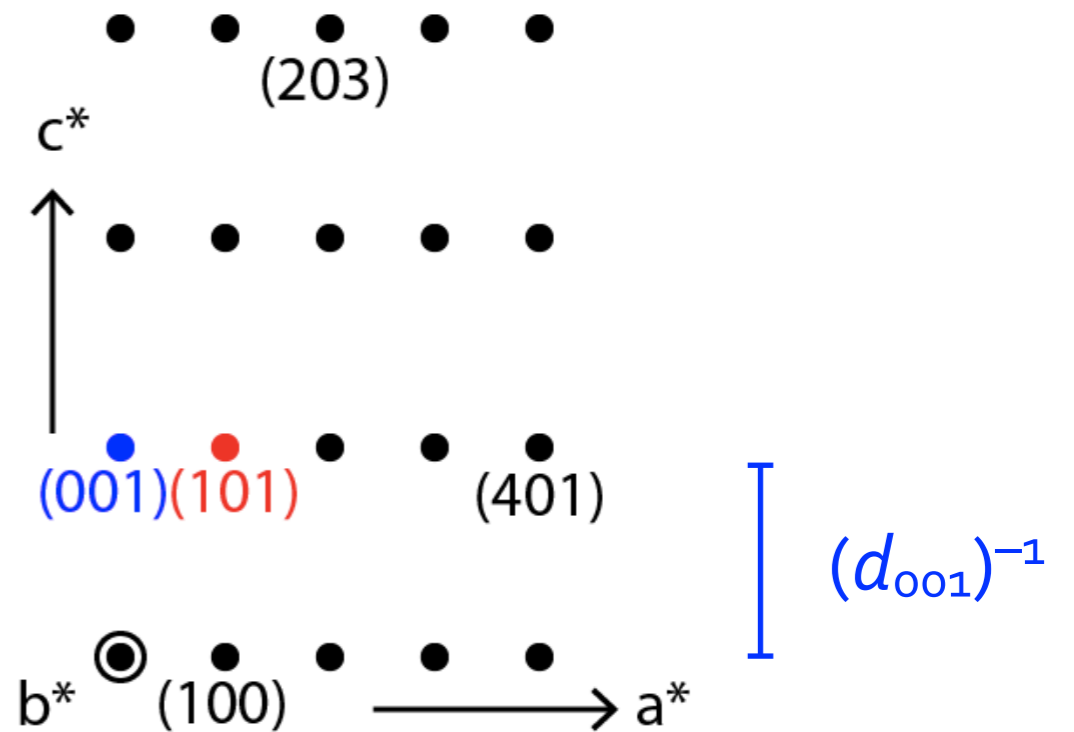
maybe there is an easier way to describe ALL of the planes of atoms?

reciprocal space and the reciprocal lattice

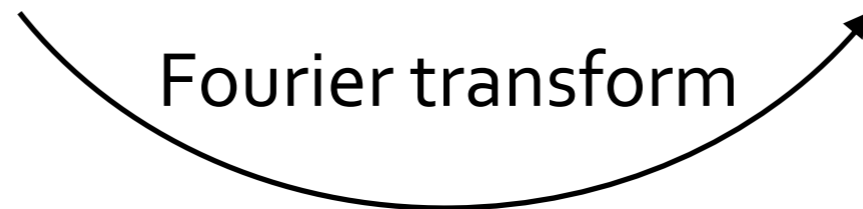
d_{hkl} -spacing:
distance
between
planes



atoms:
direct space (\AA)



recip. lattice points:
reciprocal space (\AA^{-1})



sets of planes in direct space become points in reciprocal space;
they form the basis of the reciprocal lattice.
*note the inverse relationship: big a , small a^**

Reciprocal space

$$V = a \cdot (b \times c) \quad \text{volume of the unit cell}$$

Reciprocal lattice vectors, \mathbf{a}^* , \mathbf{b}^* , and \mathbf{c}^* :

$$\mathbf{a}^* = \frac{\mathbf{b} \times \mathbf{c}}{a \cdot (\mathbf{b} \times \mathbf{c})}; \quad \mathbf{b}^* = \frac{\mathbf{c} \times \mathbf{a}}{a \cdot (\mathbf{b} \times \mathbf{c})}; \quad \mathbf{c}^* = \frac{\mathbf{a} \times \mathbf{b}}{a \cdot (\mathbf{b} \times \mathbf{c})};$$

Reciprocal lattice comprised of:

$$\mathbf{G} = h\mathbf{a}^* + k\mathbf{b}^* + l\mathbf{c}^*$$

where h , k , and l are the Miller indices (integers).

what can you learn about **crystals** from **diffraction?** *electron*

- Lattice constants
 - **locations** of points in *reciprocal space*
- Symmetry
 - **arrangement** and **presence** of points in *reciprocal space*
- Atom positions
 - **intensity** of points in *reciprocal space* [*caveat for electron diffraction*]
- Orientation
 - **intensity** and **arrangement** of points in reciprocal space

To learn about orientation, you need to know the underlying lattice structure

BaTiO₃ <001>

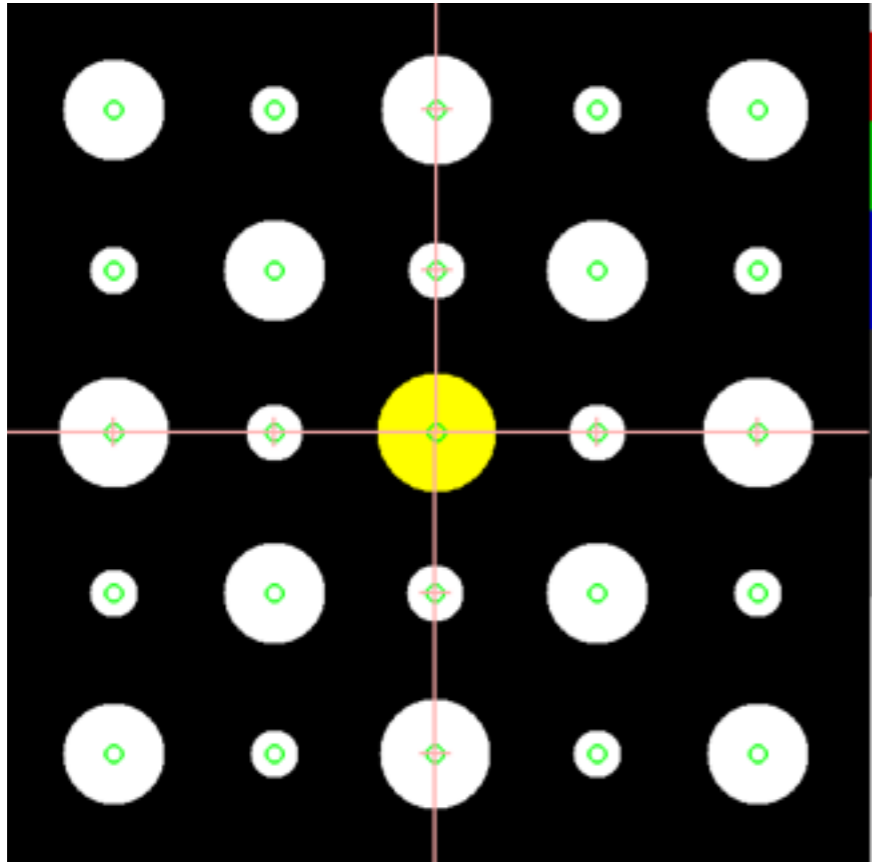
cubic - paraelectric

what changed?

tetragonal - ferroelectric

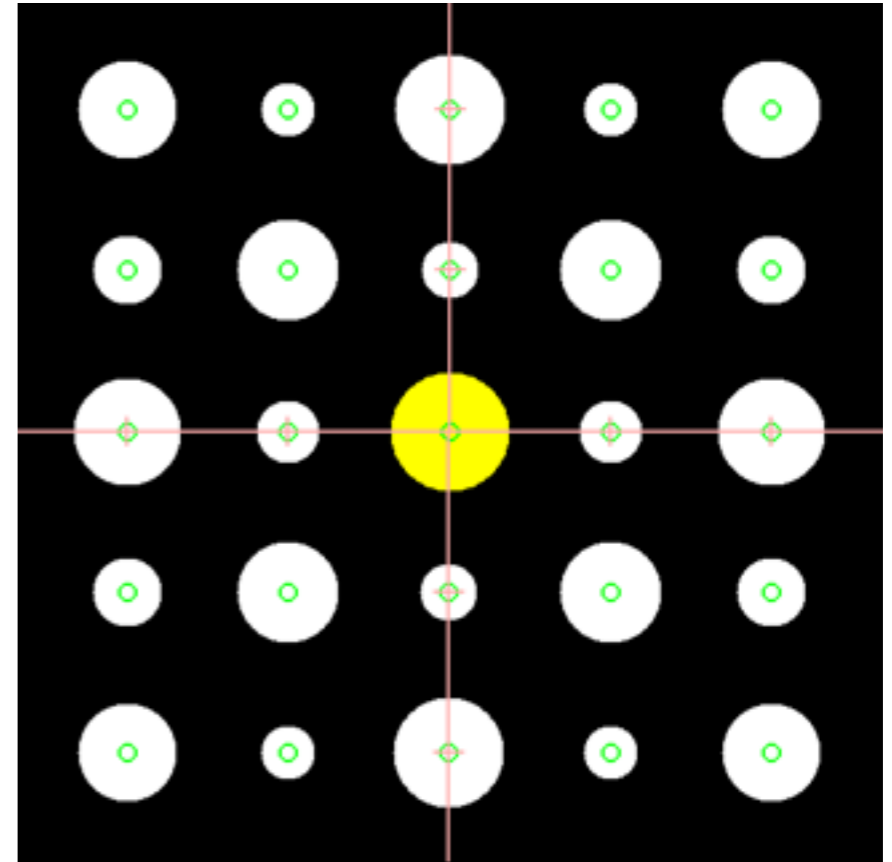
b^*

b^*

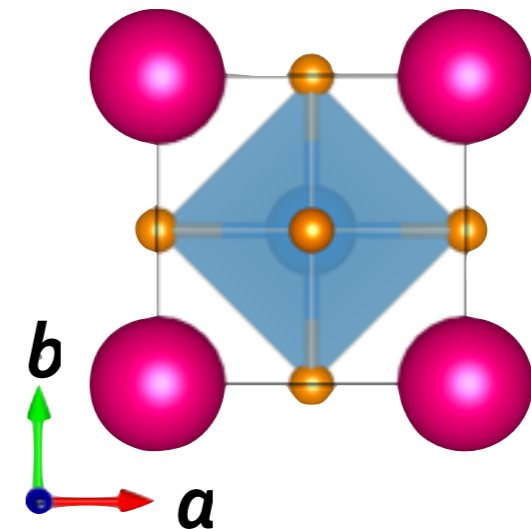
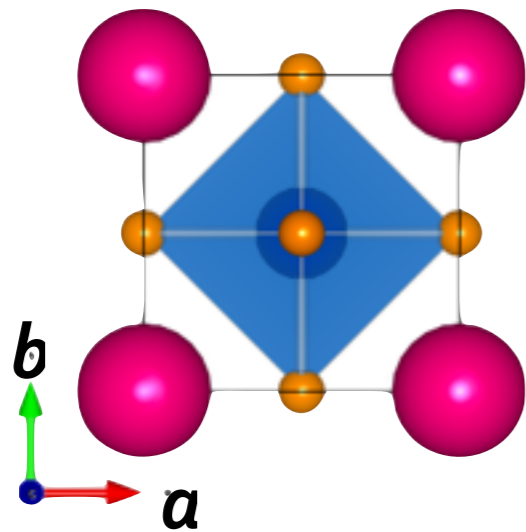


(nothing perceptible here)

a^*



a^*



simulated with ISODISTORT and VESTA

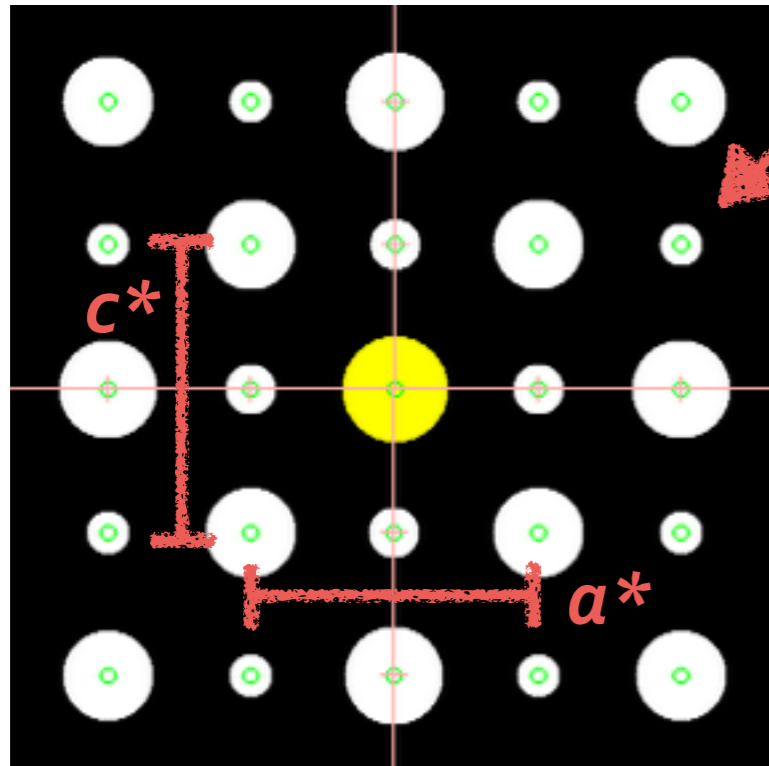
BaTiO₃ <010>

cubic - paraelectric

what changed?

tetragonal - ferroelectric

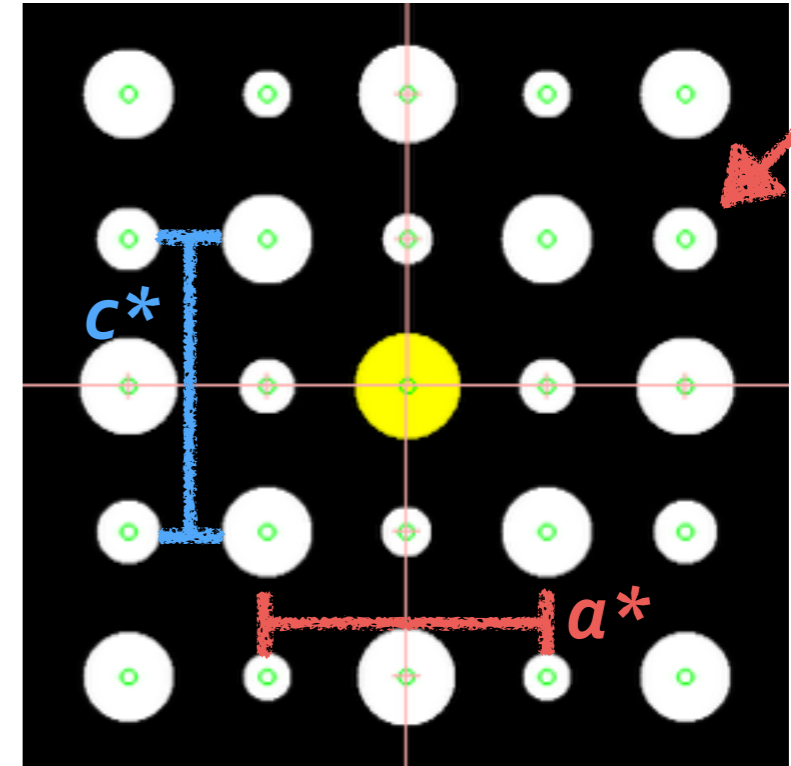
c^*



a^*

a^*

c^*



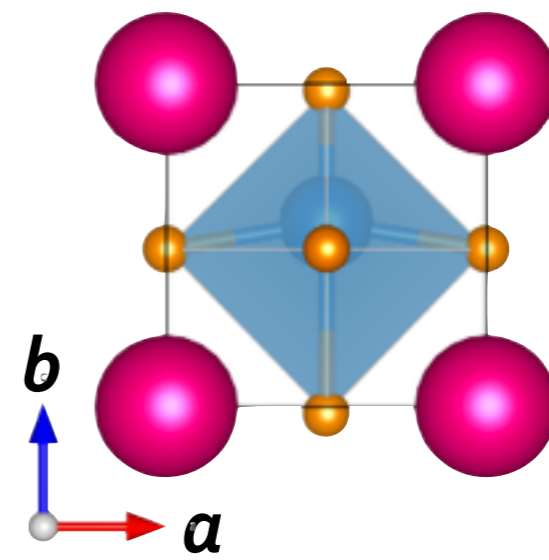
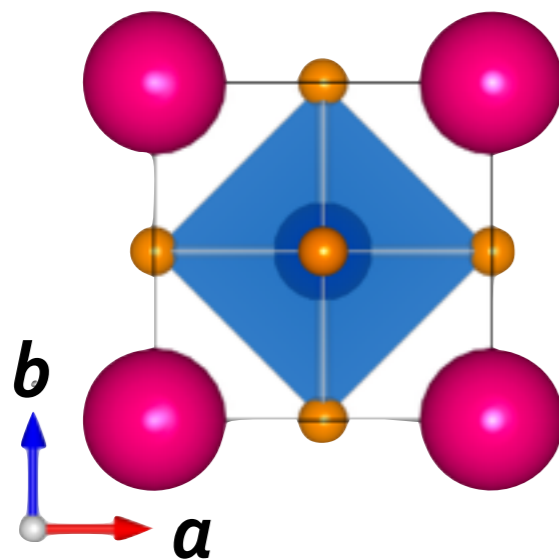
a^*

a^*

$a^* = c^*$

cell dimensions

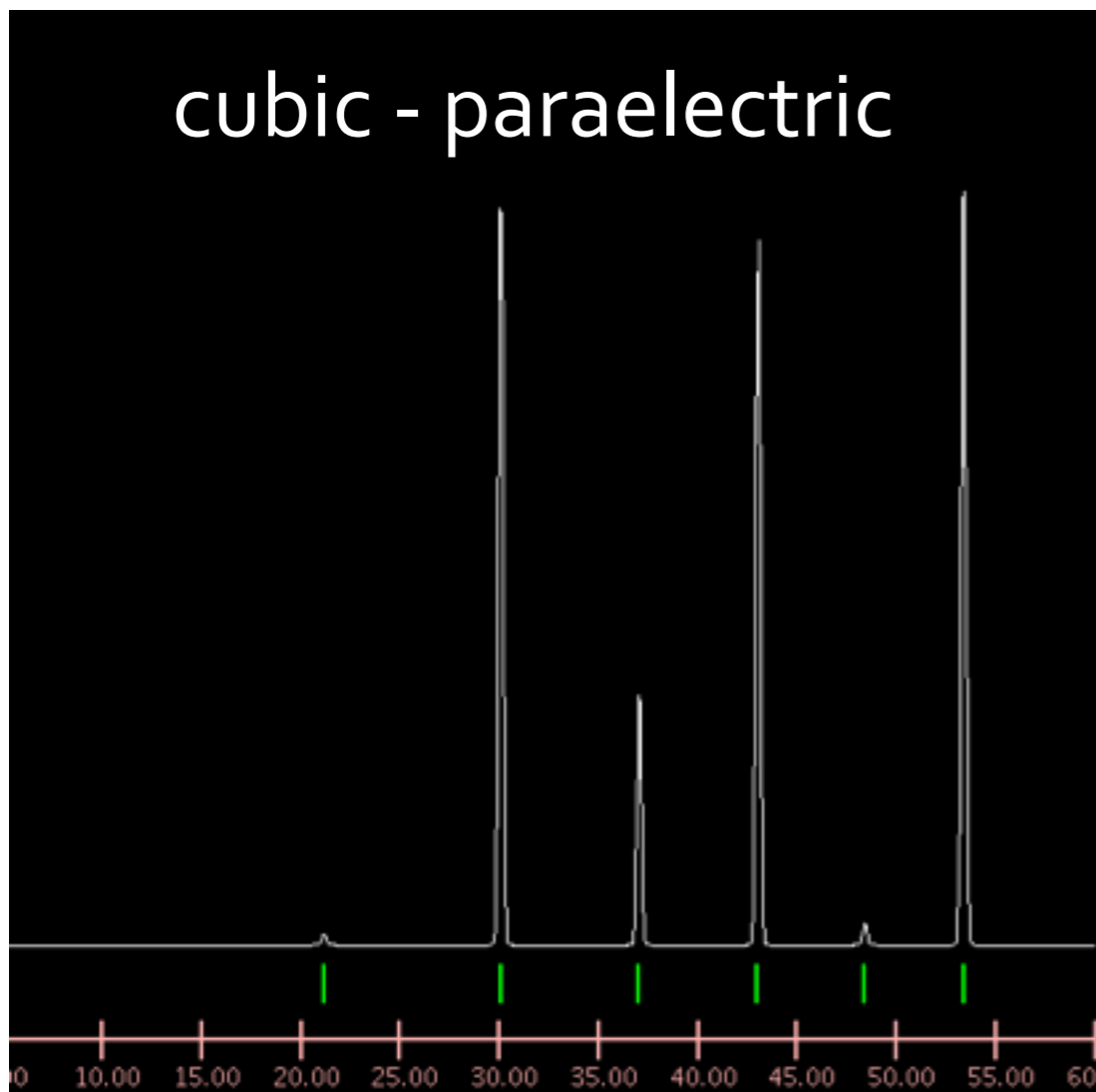
$a^* \neq c^*$



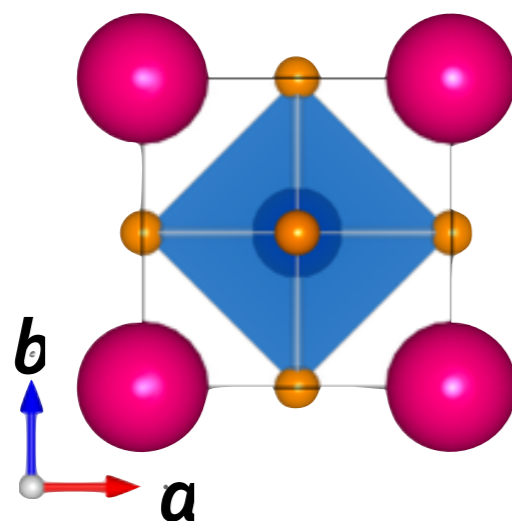
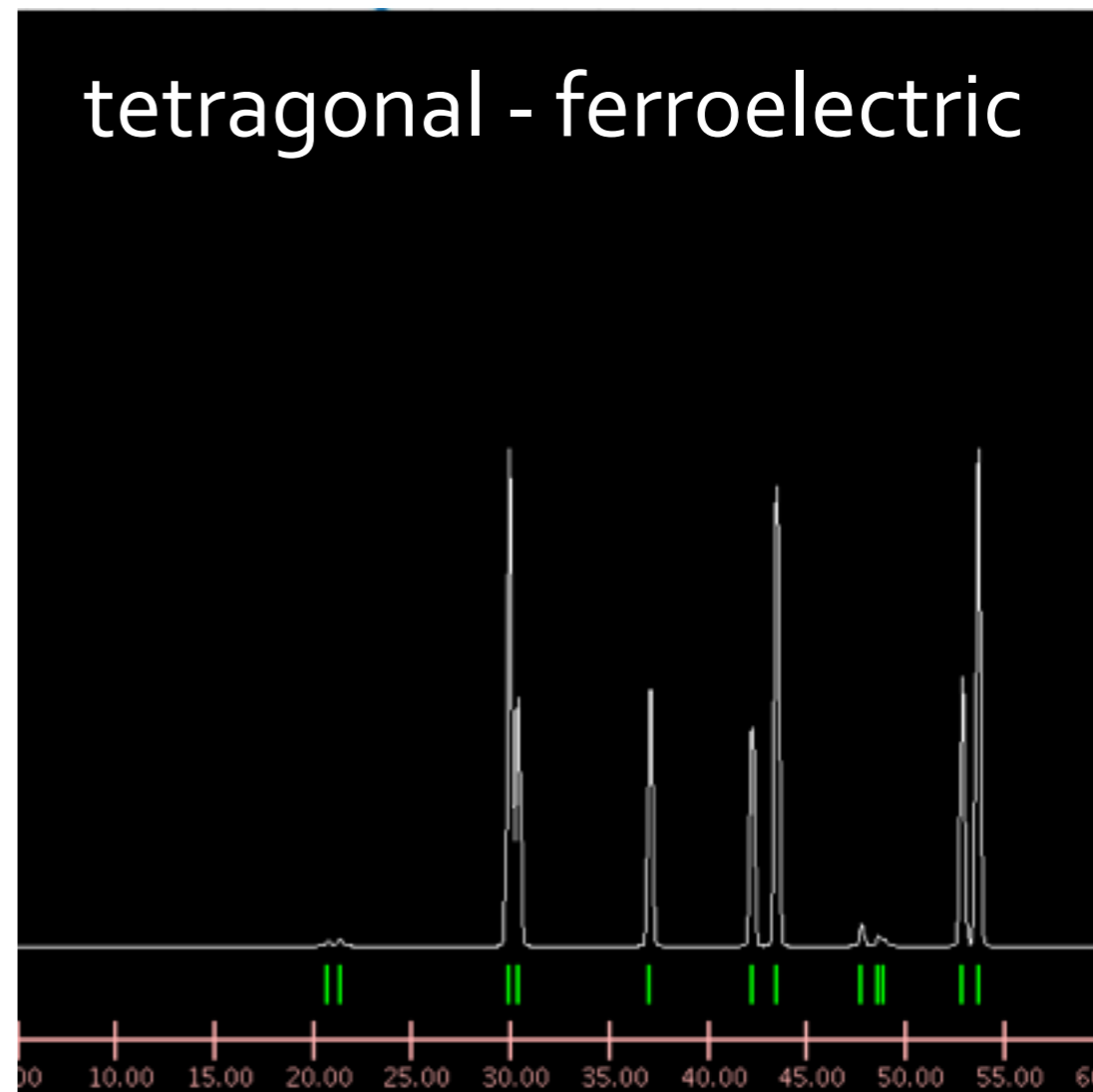
simulated with ISODISTORT and VESTA

BaTiO₃ powder x-ray diffraction

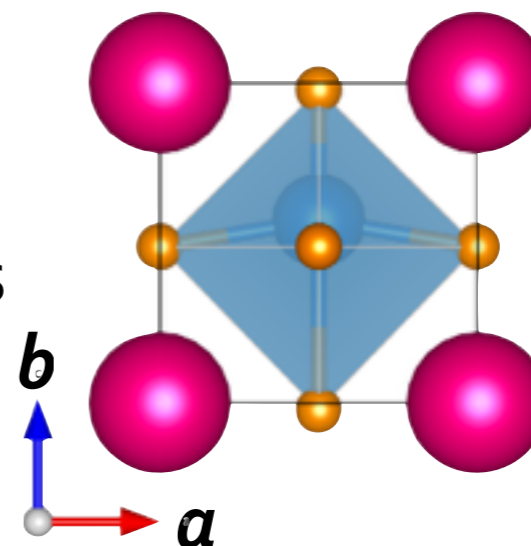
cubic - paraelectric



tetragonal - ferroelectric



distinct peak splitting
and intensity differences



simulated with ISODISTORT and VESTA

Why do crystallography in the TEM?

take advantage of the 'local' probe!

- Phase identification (lattice or point group)
- Orientation of crystals
- Finite sizes precludes macroscopic crystallography
 - *the nanostructure problem; challenges in crystal growth.*
- Supersymmetry: strong electron contrast
 - *modulated structures, charge density wave materials (e.g., TaS₂, LiFeBO₃)*
- Disorder: domains and local distortions
 - *e.g., metal-insulator transitions, ferroelectrics, superconductors*

Disorder: distortions and metal-insulator transitions

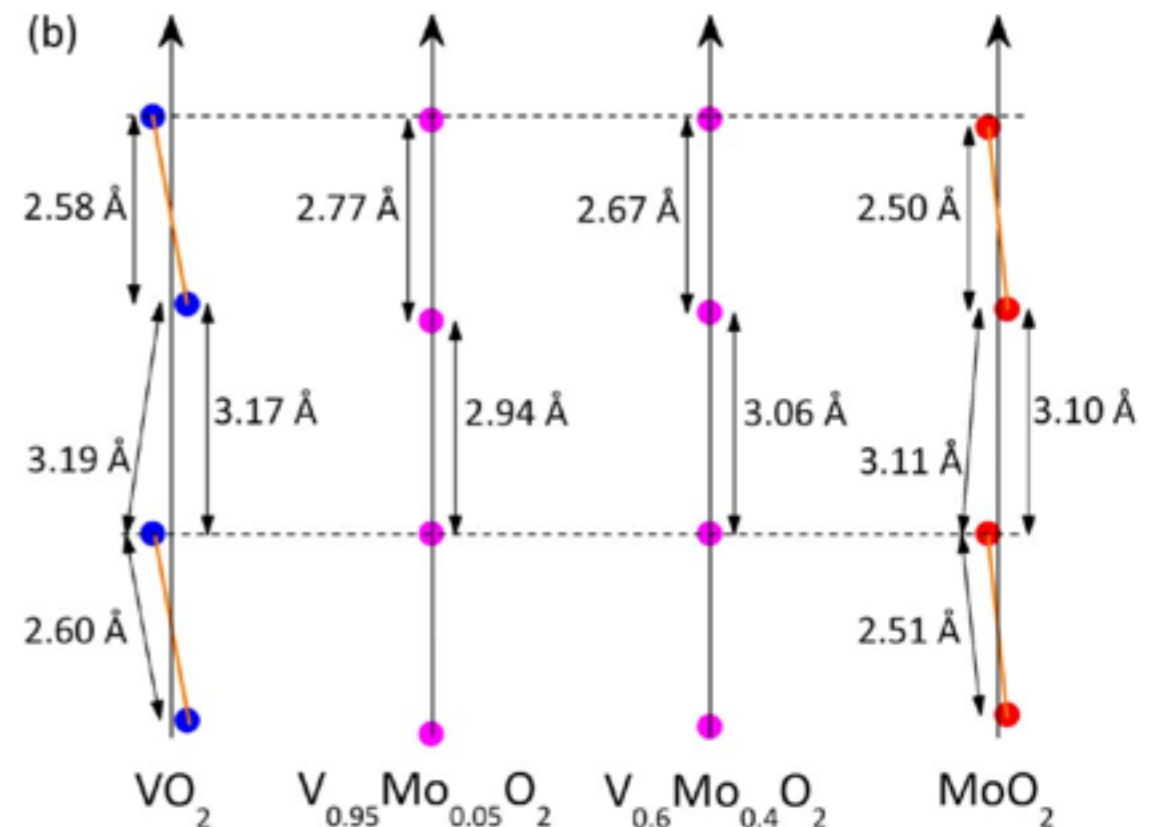
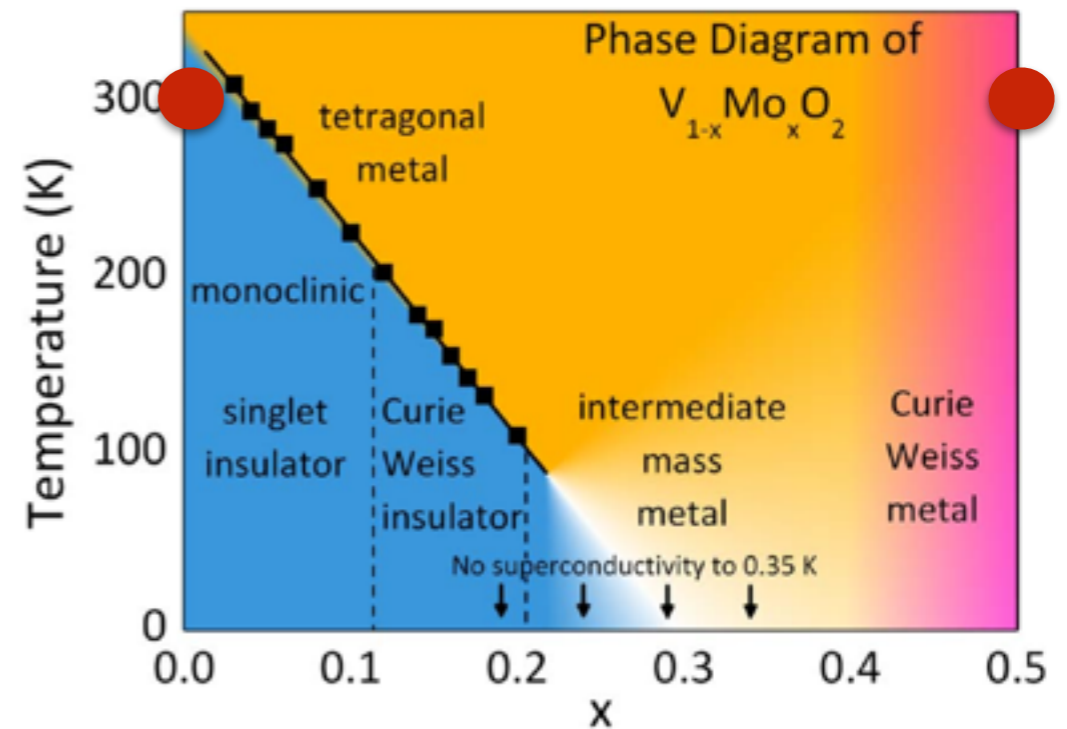


At $T = 300$ K:

Monoclinic for VO_2
– **insulator**

Tetragonal for $V_{0.5}Mo_{0.5}O_2$
– *weird* metal

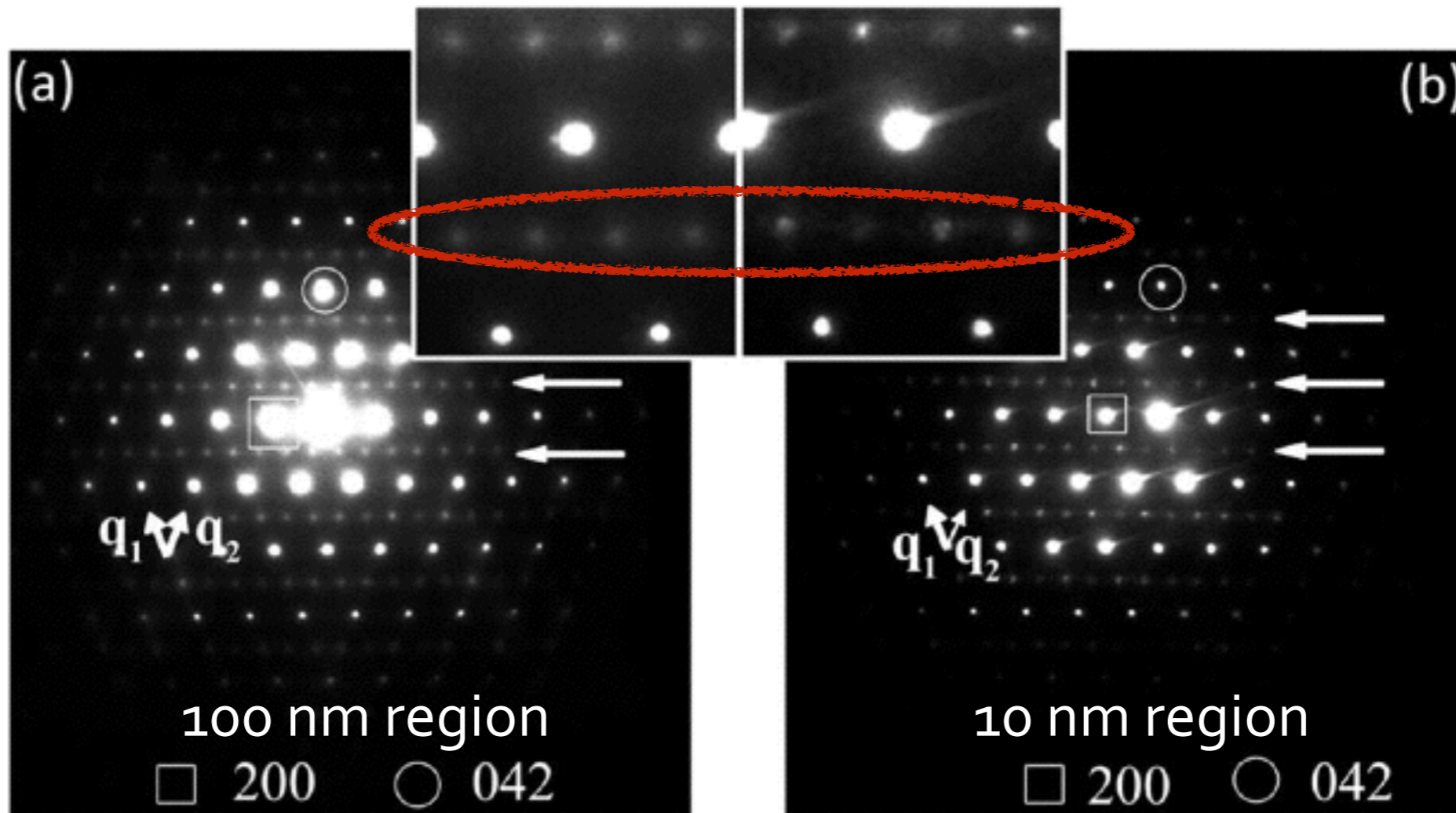
modulation of $M-M$ distances
within the material



Local distortions in $V_{0.5}Mo_{0.5}O_2$

X-ray crystallography: tetragonal *on average*

Zooming into a small region of the sample with TEM:



Extra reflections: broken tetragonal symmetry!

monoclinic on a short length scale

yet still metallic

Kinematic theory for diffraction

- Main assumptions:
 - single scattering (no multiple scattering)
 - elastic scattering
 - monochromatic incident beam
 - distortion free crystal
 - only a small fraction of e's scattered
 - no absorption

Kinematic theory for diffraction

- Main assumptions: (**non-ideal assumptions for the TEM**)
 - **single scattering (no multiple scattering)**
except maybe for very thin samples
 - elastic scattering
 - monochromatic incident beam
 - **distortion free crystal**
 - **only a small fraction of e's scattered**
 - no absorption

intensities: structure factor (kinematic scattering)

- F_{hkl} = structure factor

$$F_{hkl} = \sum_n f_n \exp[2\pi i(hx_n + ky_n + lz_n)]$$

- atom positions are at $x_n, y_n,$ and z_n . f_n is the atomic scattering factor. The summation is taken over all atoms in the unit cell.
- We actually measure intensity that is proportional to $|F_{hkl}|^2$
 - Since we measure the square, we lose the relative phases of the periodic components: “the crystallographic phase problem”

systematic absences from the structure factor

Consider a body-centered cell, with identical atoms at $(0,0,0)$ and $(1/2, 1/2, 1/2)$

$$F_{hkl} = \sum_n f_n \exp[2\pi i(hx_n + ky_n + lz_n)]$$

$$F_{hkl}/f_n = \exp[2\pi i(h \cdot 0 + k \cdot 0 + l \cdot 0)] + \exp[2\pi i(1/2h + 1/2k + 1/2l)]$$

$$F_{hkl}/f_n = 1 + \exp[\pi i(h + k + l)]$$

if $h + k + l = \text{even}$:

$$F_{hkl}/f_n = 1 + 1 + 0 = 2$$

intensity is observed

if $h + k + l = \text{odd}$:

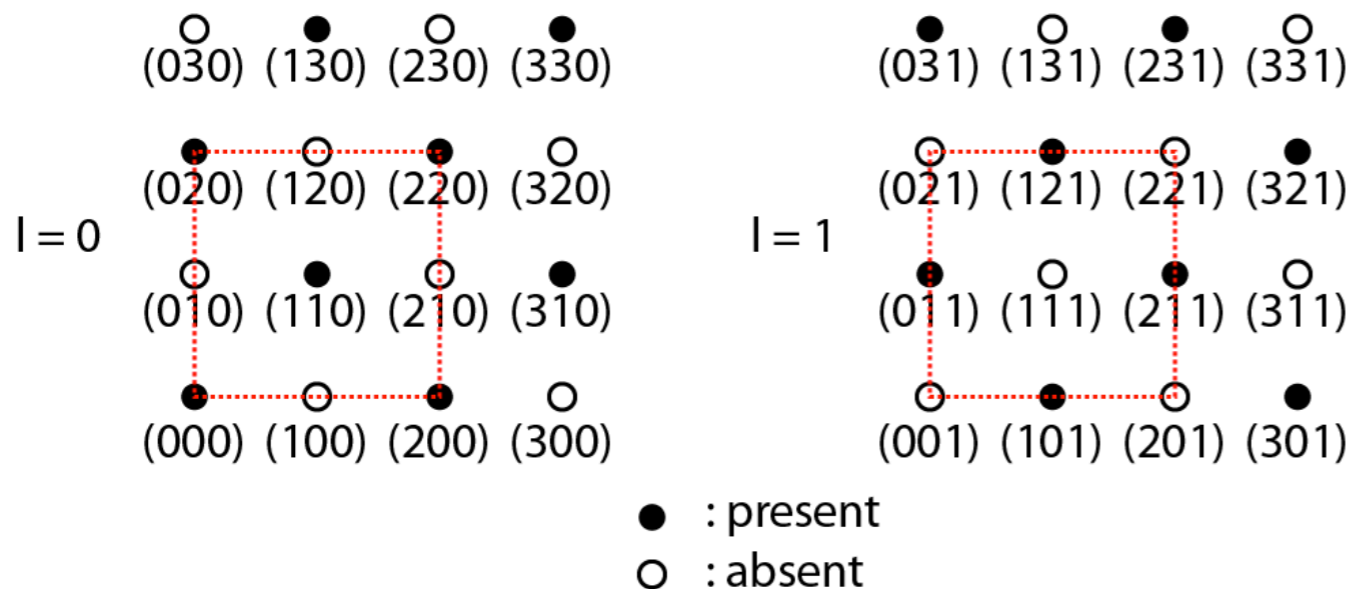
$$F_{hkl}/f_n = 1 - 1 + 0 = 0$$

these reflections are absent

systematic absences from the structure factor

A consequence of this is that a BCC crystal has a FCC reciprocal lattice.

Other translation operations create different systematic absences.



Translational symmetry elements and their systematic absences

Symmetry element	Affected reflection	Condition for systematic absence ←
<i>Centred-cells:</i>		
Body-centred (I)	hkl	$(h+k+l)$ odd
Face-centred (A)	hkl	$(k+l)$ odd
Face-centred (B)	hkl	$(h+l)$ odd
Face-centred (C)	hkl	$(h+k)$ odd
Face-centred (F)	hkl	$(h+k)$ odd $(h+l)$ odd $(k+l)$ odd
<i>Screw axes:</i>		
Two-fold screw (2_1) along a	$h00$	h odd
2_1 along b	$0k0$	k odd
2_1 along c	$00l$	l odd
<i>Glide planes $\perp b$:</i>		
Translation ($a/2$) (a -glide)	$h0l$	h odd
Translation ($c/2$)	$h0l$	l odd
Translation ($a/2 + c/2$) (n -glide)	$h0l$	$(h+l)$ odd
Translation ($a/4 + c/4$) (d -glide)	$h0l$	$(h+l) = 4n + 1, 2, 3$

other examples

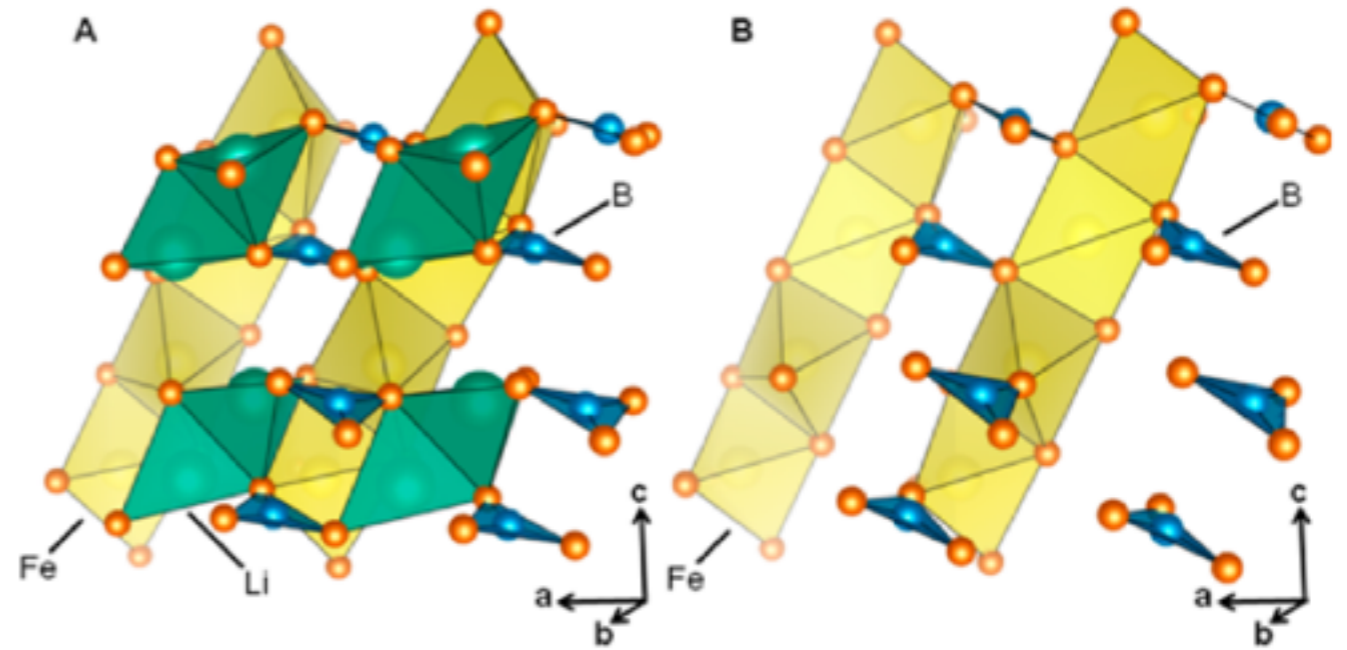
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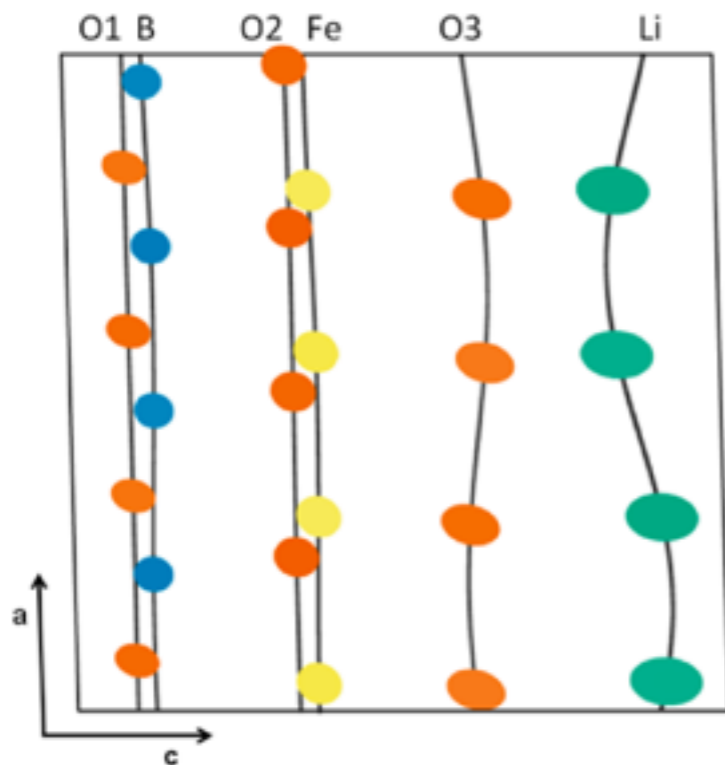
what is a modulated, supersymmetric structure?

- example: LiFeBO_3
- Li-ion battery cathode material

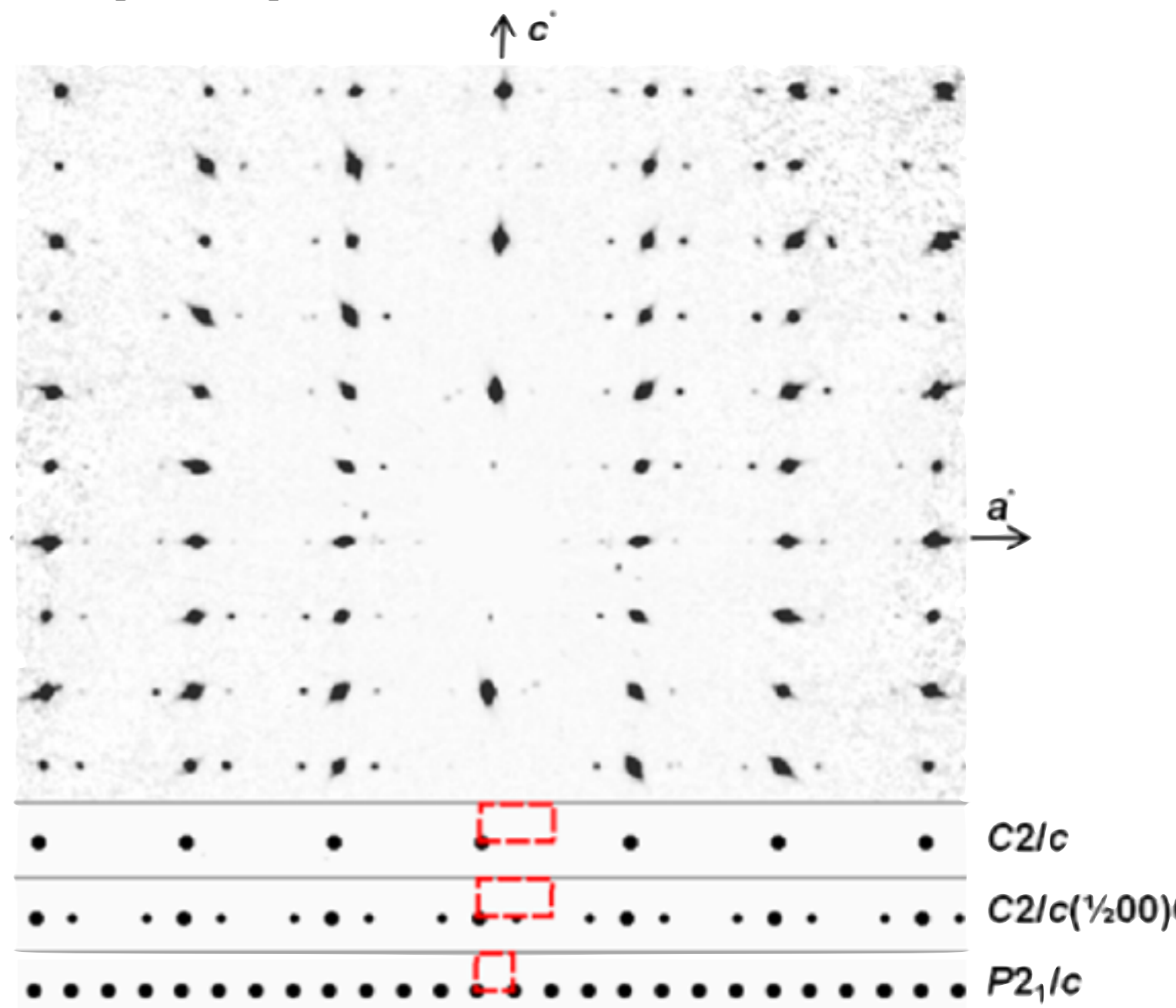


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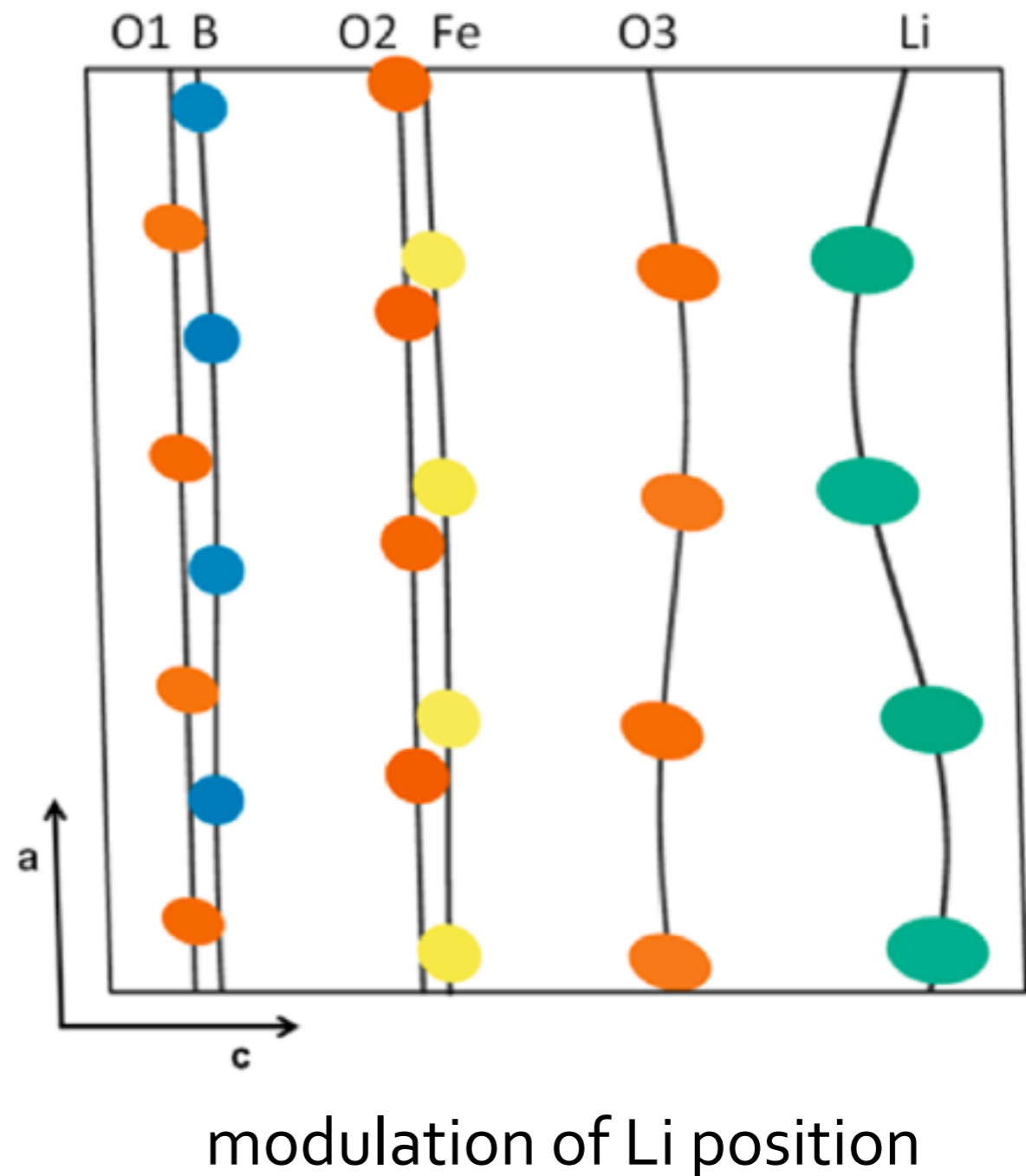
modulation of Li position



extra, weak reflections observed
at $1/2 a^*$, or $2a$!!!

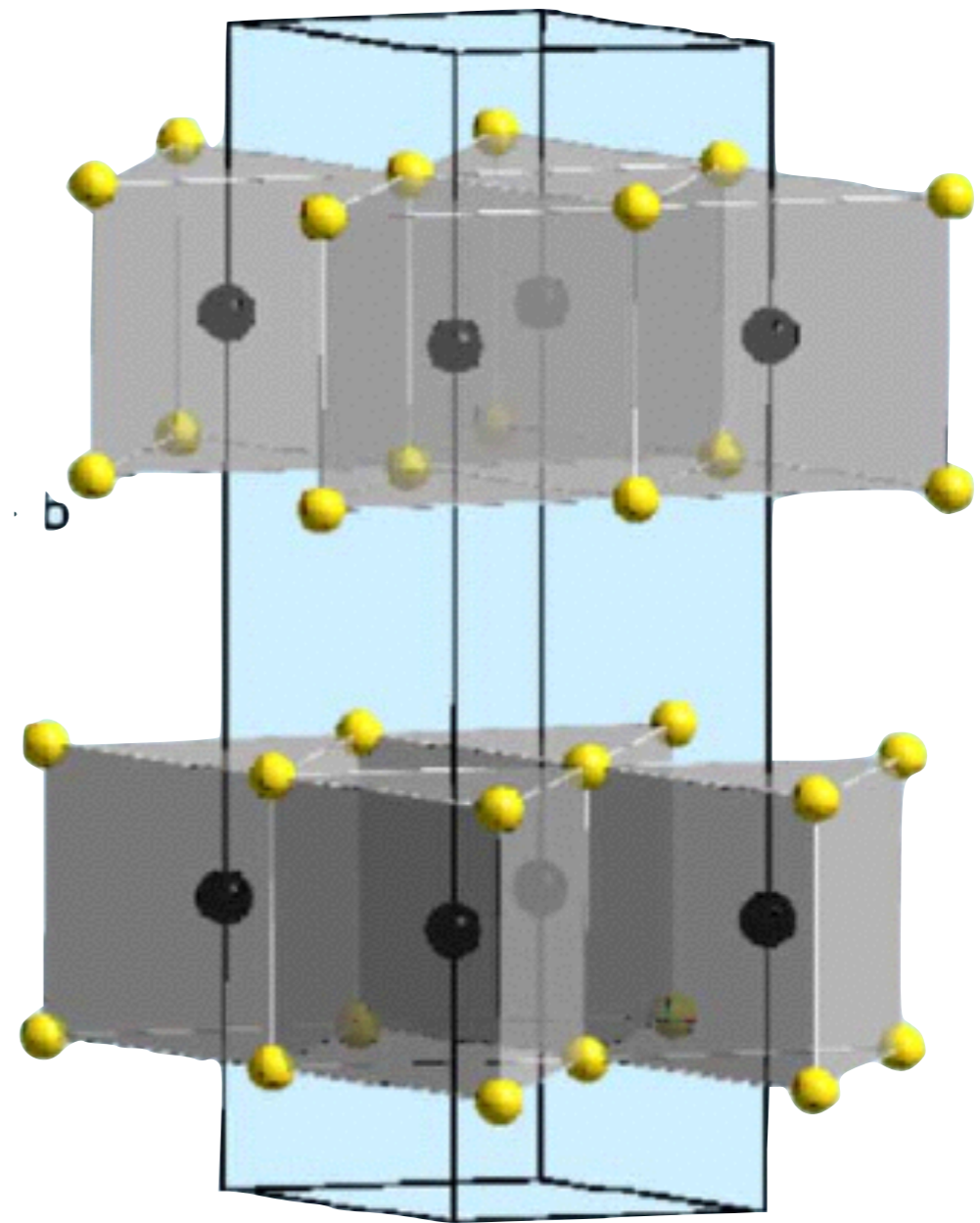
what is a modulated, supersymmetric structure?

- example: LiFeBO_3
- Li-ion battery cathode material
- Periodicity of the Li positions is **commensurate** with the underlying reciprocal lattice but has longer-range order.



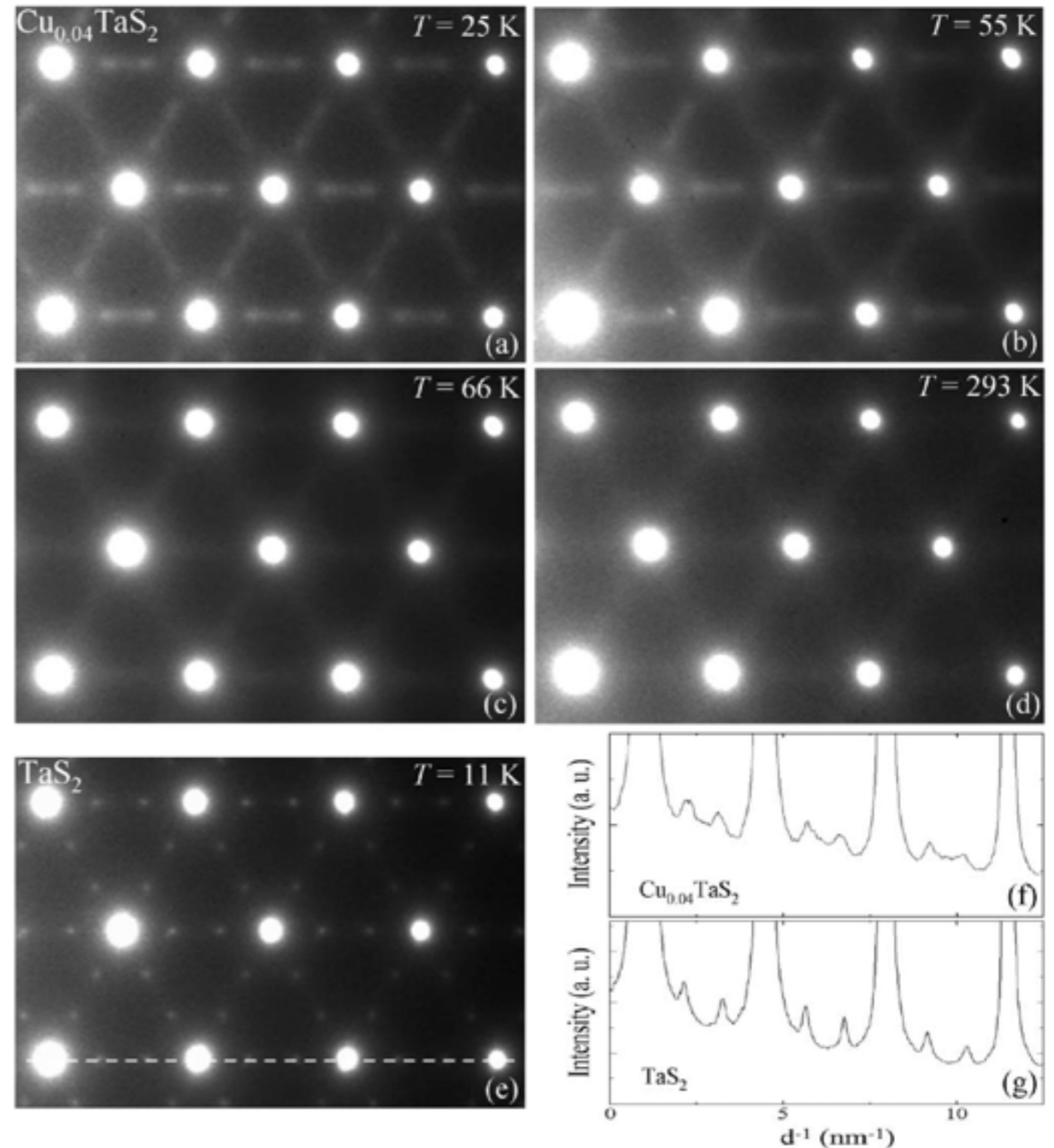
supersymmetry: higher dimensional crystallography

- example: charge-density wave order in TaS_2
 - $q = [0.341, 0, 0]$
- that is, the actual unit cell is $(0.341)^{-1}$ times longer along \mathbf{a} !



supersymmetry: higher dimensional crystallography

- example: charge-density wave order in TaS_2
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additional crystallography resources:

http://www.aps.anl.gov/Xray_Science_Division/Powder_Diffraction_Crystallography/Introduction_to_Crystallography/

Web-based video lectures and other resources

<http://www.doitpoms.ac.uk/tlplib/crystallography3/index.php>

great animations and illustrations

<http://www.iucr.org/education/pamphlets>

handouts from the International Union on Crystallography

<http://jp-minerals.org/vesta/en/>

software for visualizing crystal structures and lattice planes (VESTA)

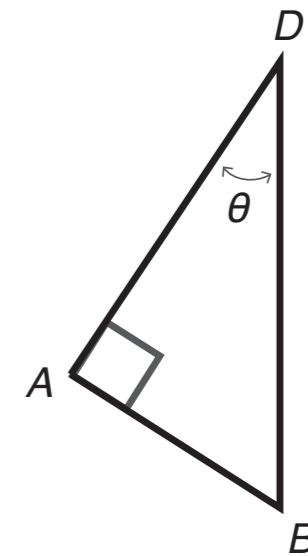
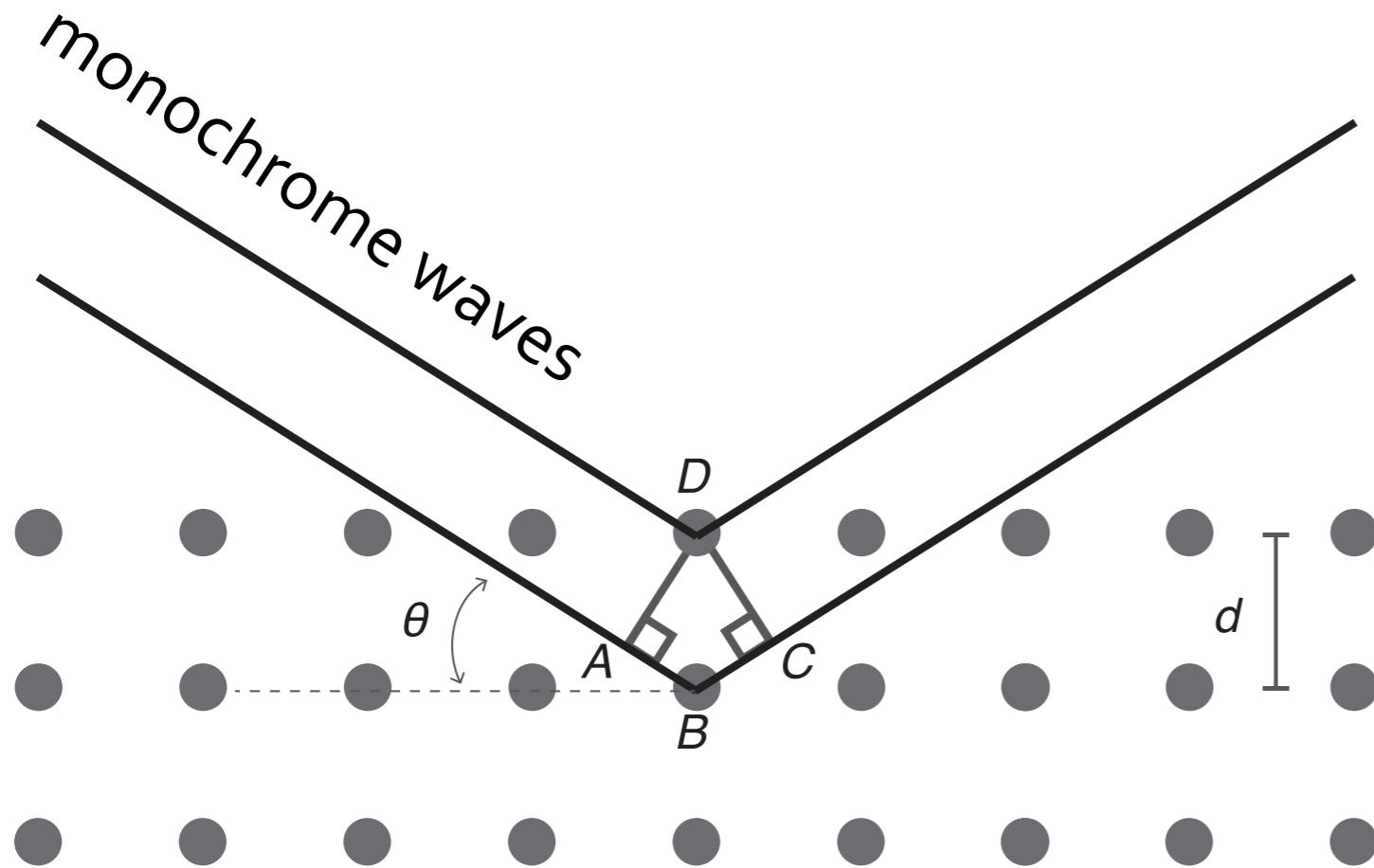
<http://stokes.byu.edu/iso/isodistortform.php>

ISODISTORT: online tool for simulating crystals, their distortions, and their diffraction patterns

more on diffraction

Diffraction: Bragg's Law

$$2d \sin \theta = n\lambda$$



$$\sin \theta = AB/BD = AB/d$$

$$AB = BC$$

if the path-length difference of the two waves is equal to an integer multiple of the wavelength, then constructive interference

constructive if $2AB = n\lambda$

$$2d \sin \theta = n\lambda$$

Diffraction: the Laue Equations

Instead, we can use vector notation

$$|\mathbf{s} - \mathbf{s}_0| = 2 \sin \vartheta$$

and

$$|\mathbf{d}^*_{hkl}| = 1/d_{hkl}$$

Their ratio is λ :

$$(\mathbf{s} - \mathbf{s}_0) / \lambda = \mathbf{d}^*_{hkl}$$

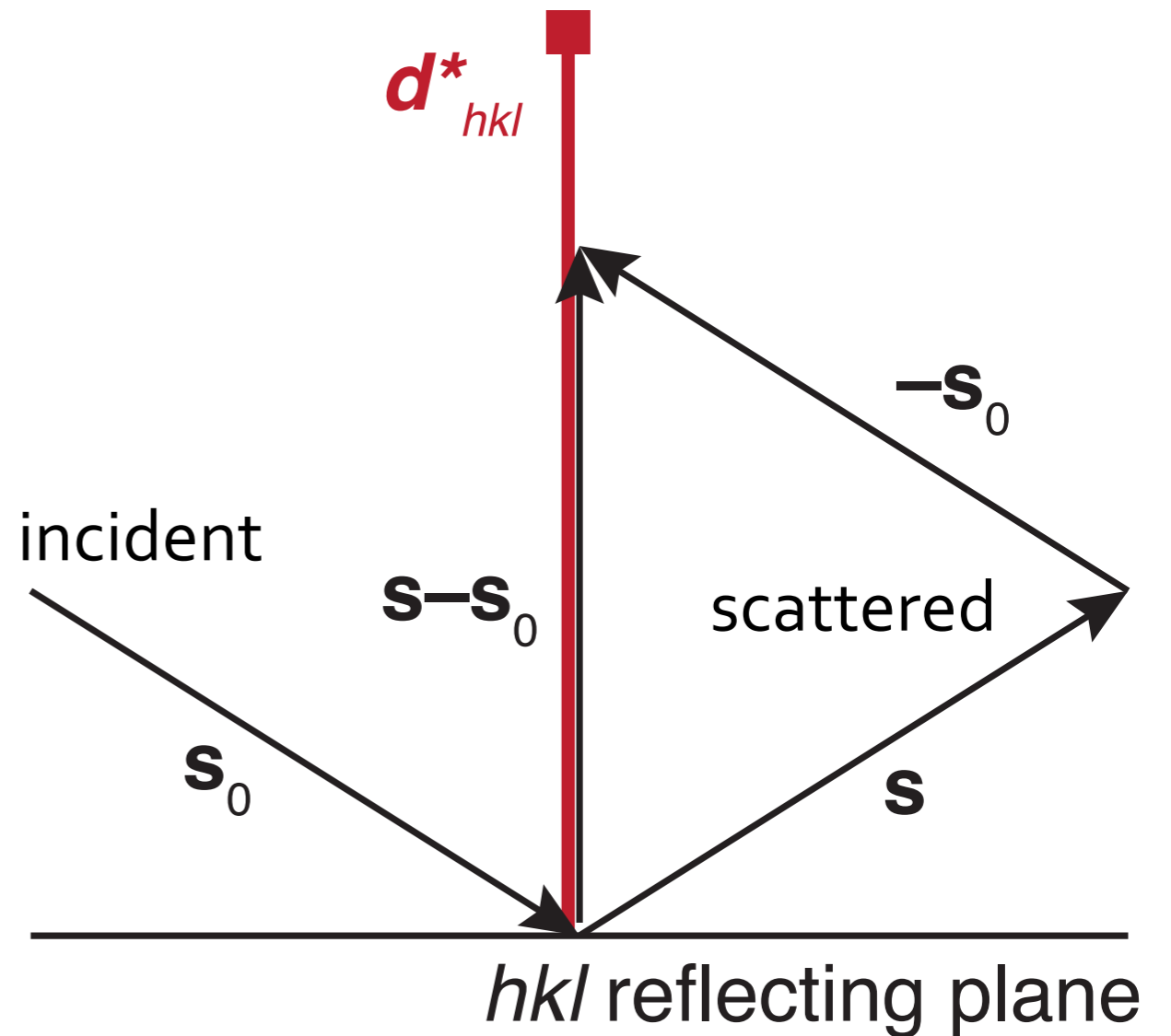
Because:

$$\mathbf{d}^*_{hkl} = h\mathbf{a}^* + k\mathbf{b}^* + l\mathbf{c}^*$$

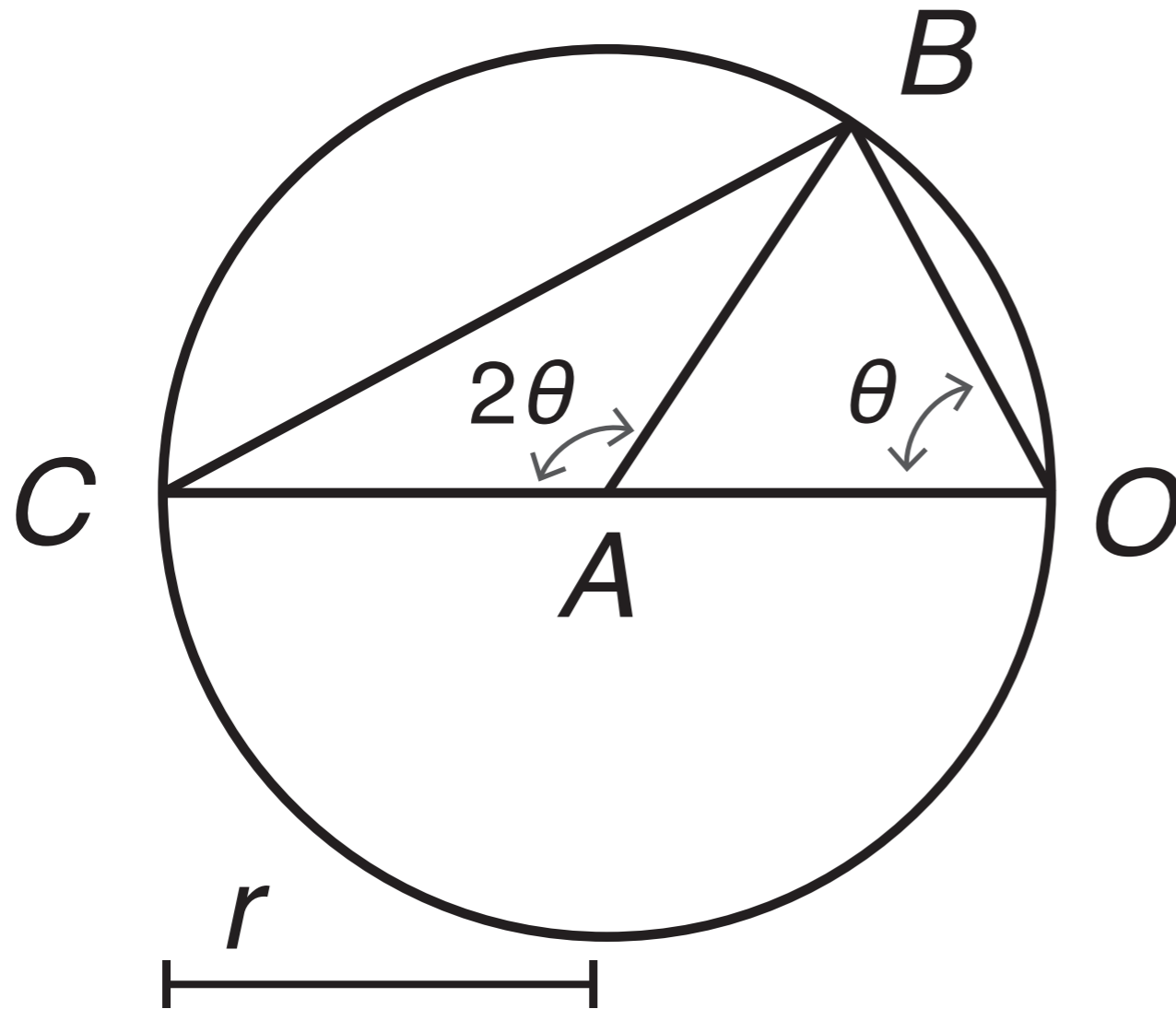
The first Laue equation is:

$$\mathbf{a}(\mathbf{s} - \mathbf{s}_0) = n_x \lambda$$

then repeat for \mathbf{b} , n_y ; \mathbf{c} , n_z .



Diffraction: The Ewald Sphere



From trigonometry:

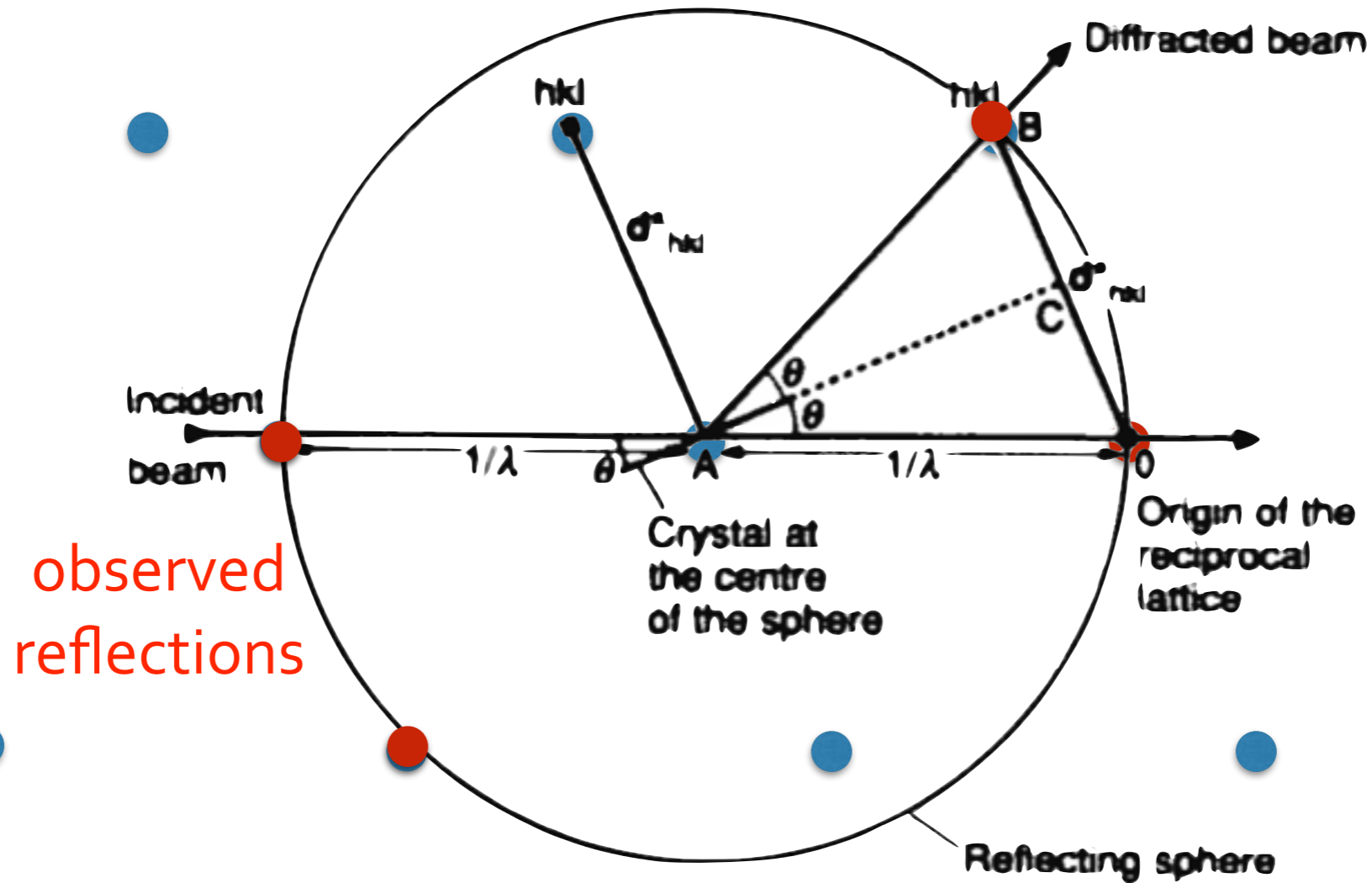
$$\sin \theta = \frac{OB}{2r}$$

In reciprocal space, we can set $r = 1/\lambda$. If O is the origin in reciprocal space, and B falls on a reciprocal lattice point, then:

$$\sin \theta = \frac{1}{\frac{d_{hkl}}{2\lambda}}$$

$$2d_{hkl} \sin \theta = \lambda$$

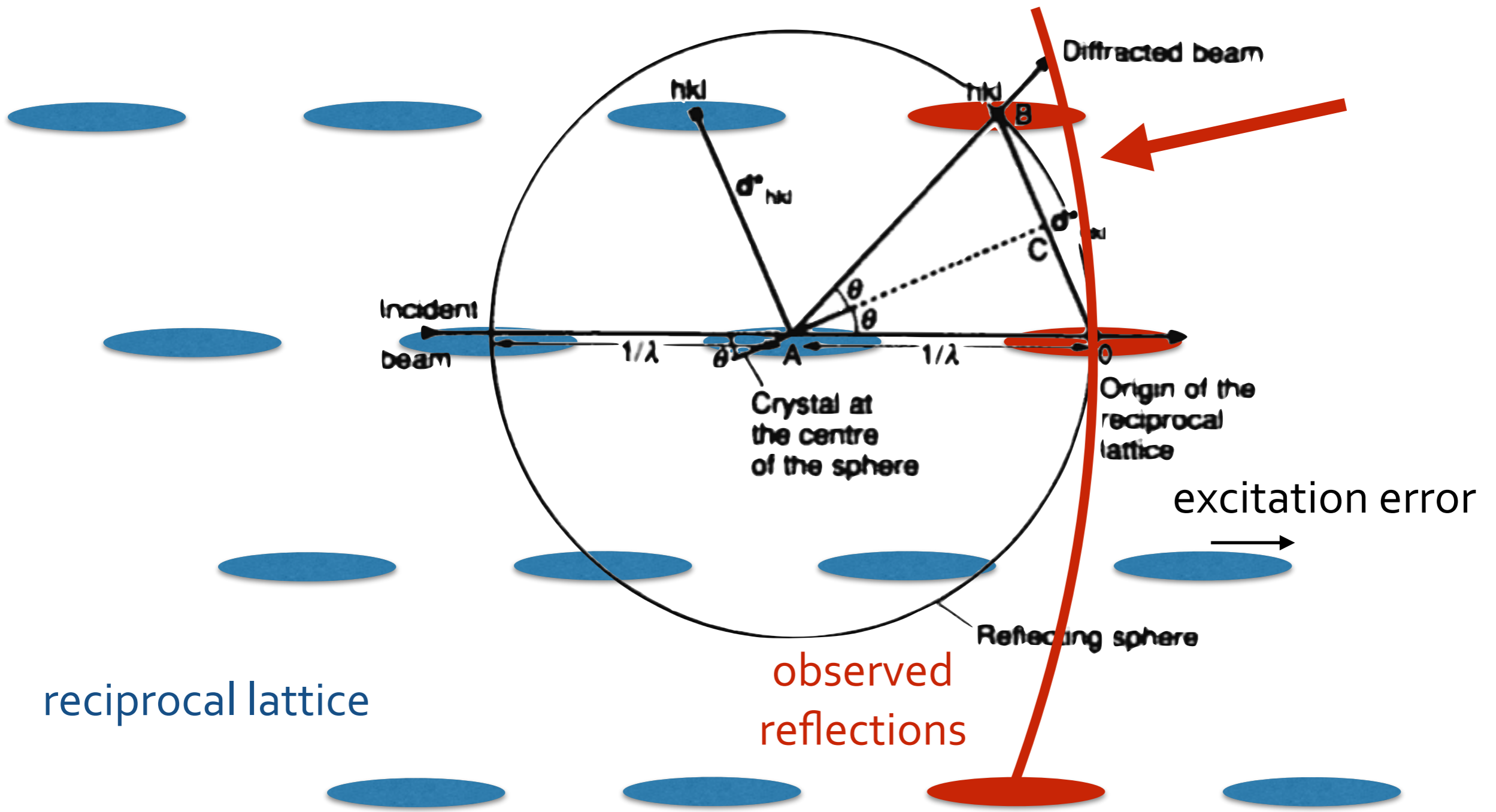
Diffraction: The Ewald Sphere



reciprocal lattice

the Ewald sphere in the TEM

relrods: reciprocal lattice rods *finite size effect of a thin sample*
the Ewald sphere appears flat! ($\lambda \sim 0.05 \text{ \AA}$ vs. 1.54 \AA for Cu $K\alpha$ X-rays)



deviations from kinematic theory in electron diffraction

- **multiple scattering - dynamical diffraction**
 - need to use the “Howie-Whelan equations”
 - simultaneous equations that describe the direct beam and scattered beam
- **“relrods”**
 - thin samples cause reciprocal lattice points to be streaked (think nanoparticle broadening)
 - “excitation error”

Diffraction in the Electron Microscope

V (kV)

λ (Å)

20

0.086

100

0.037

1000

0.009

structure factor in the TEM

- In the TEM, because we have *multiple scattering*, we often **see reflections that should be systematically absent**.
 - *This is a breakdown of the kinematic approximation.*
 - Electrons strongly interact with matter because they are charged particles; therefore, they often scatter many times.
 - Electrons and X-rays scatter $\propto (\text{atomic number})^2$; however, X-rays interact much less.
- **Conclusion: intensities in electron diffraction patterns require non-trivial analysis to solve atomic structures**

Diffraction in the Electron Microscope

- observable patterns:
 - amorphous ring pattern
 - single crystal SAED pattern
 - polycrystalline SAED pattern
 - CBED pattern
 - Kikuchi patterns